

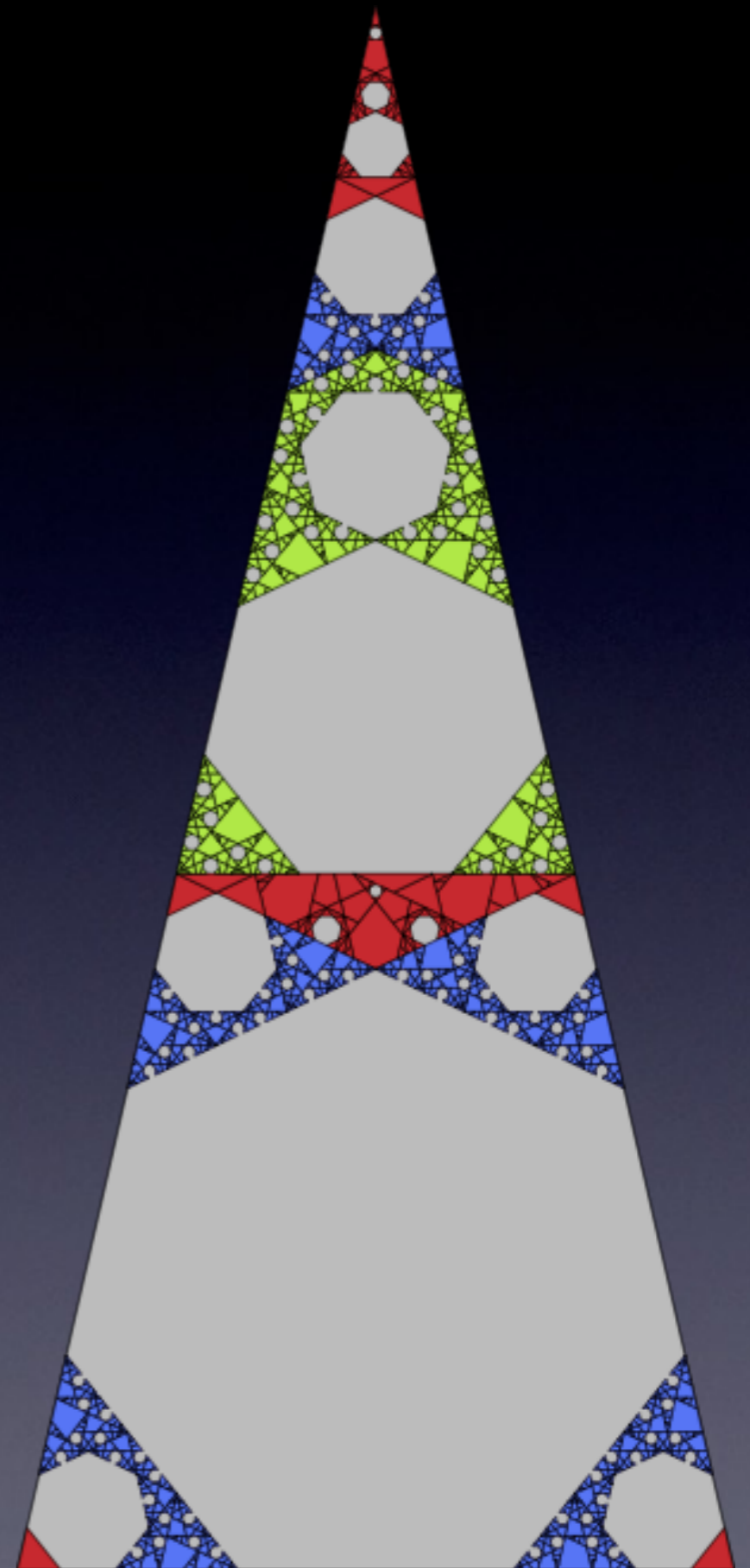
# Renormalization

in parametrised families  
of polygon-exchange transformations

**Franco Vivaldi**

Queen Mary, University of London

with J H Lowenstein (New York)



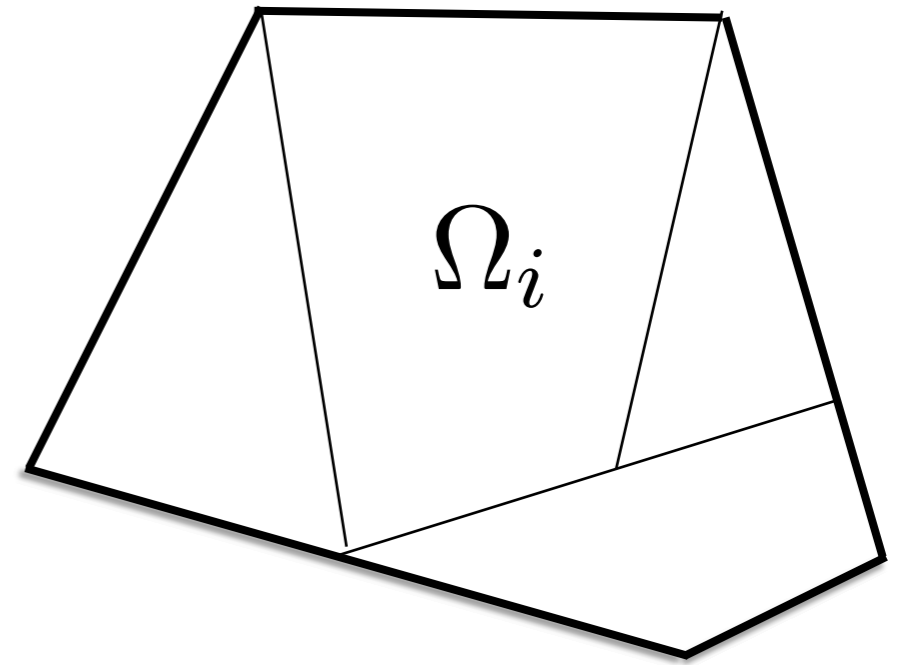
# Piecewise isometries

the space:

$$\Omega \subset \mathbb{R}^n$$

$$\Omega = \overline{\bigcup \Omega_i}$$

a finite collection of pairwise disjoint open polytopes (intersection of open half-spaces), called the **atoms**.



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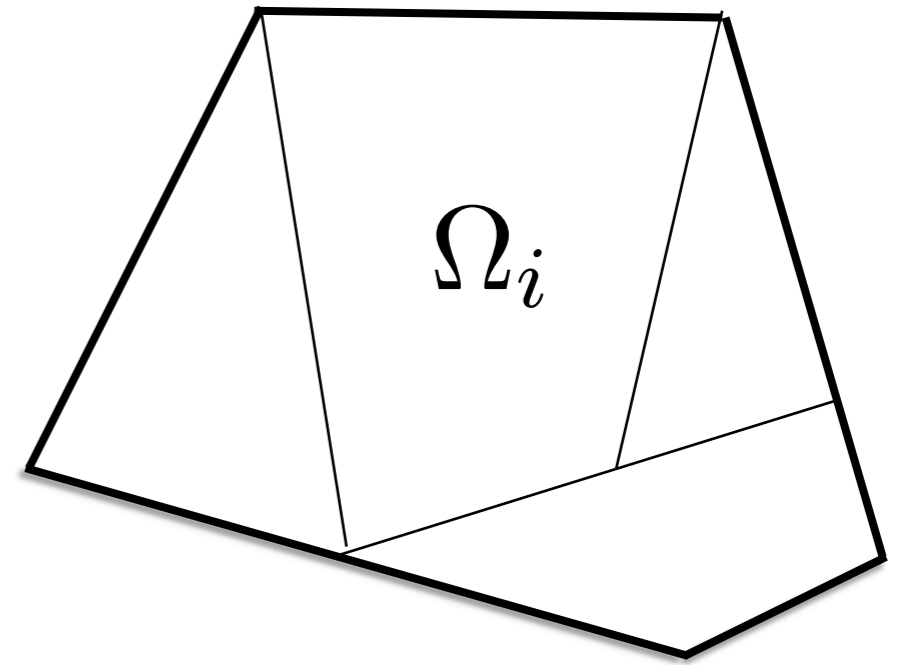
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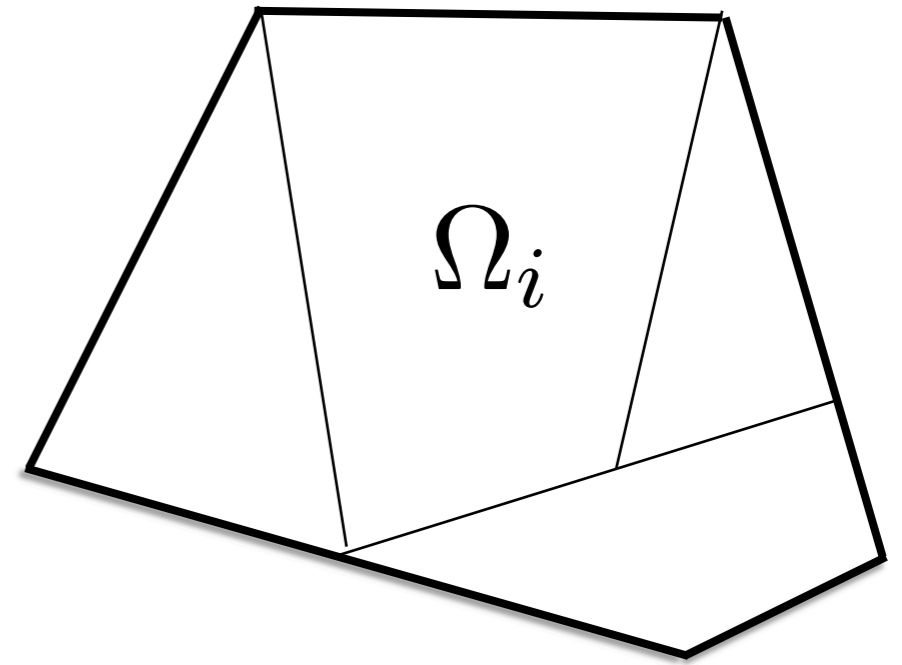
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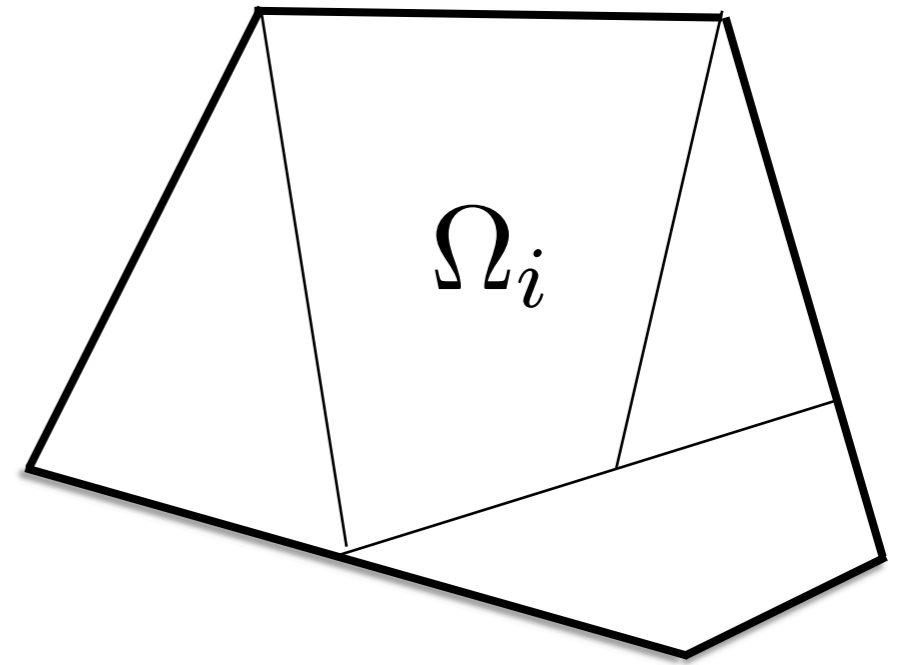
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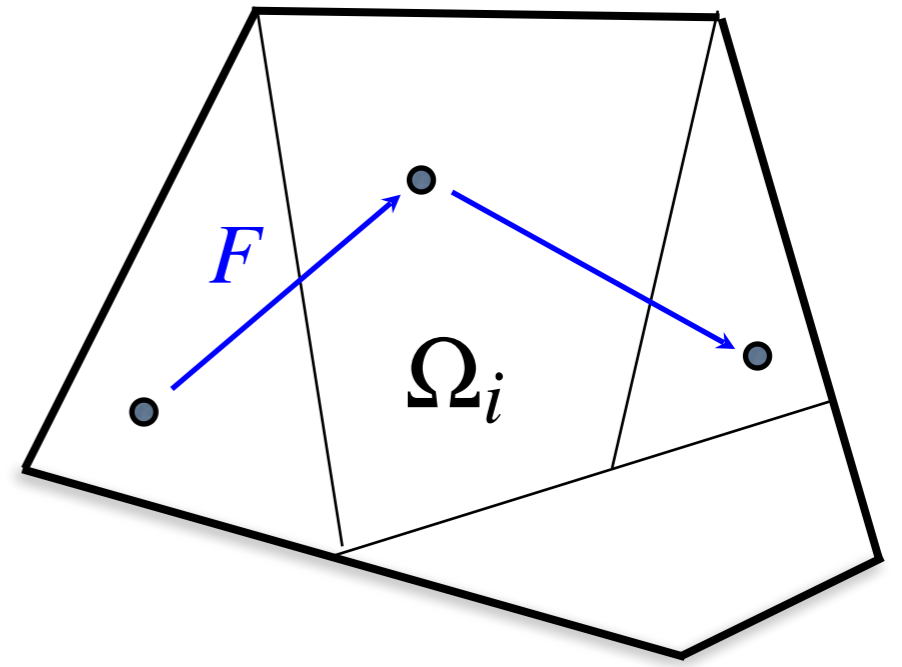
$$F : \Omega \rightarrow \Omega \quad F|_{\Omega_i} \text{ is an isometry}$$

If  $F$  is invertible, then  $F$  is volume-preserving.

**Theorem** (Gutkin & Haydin 1997, Buzzi 2001)

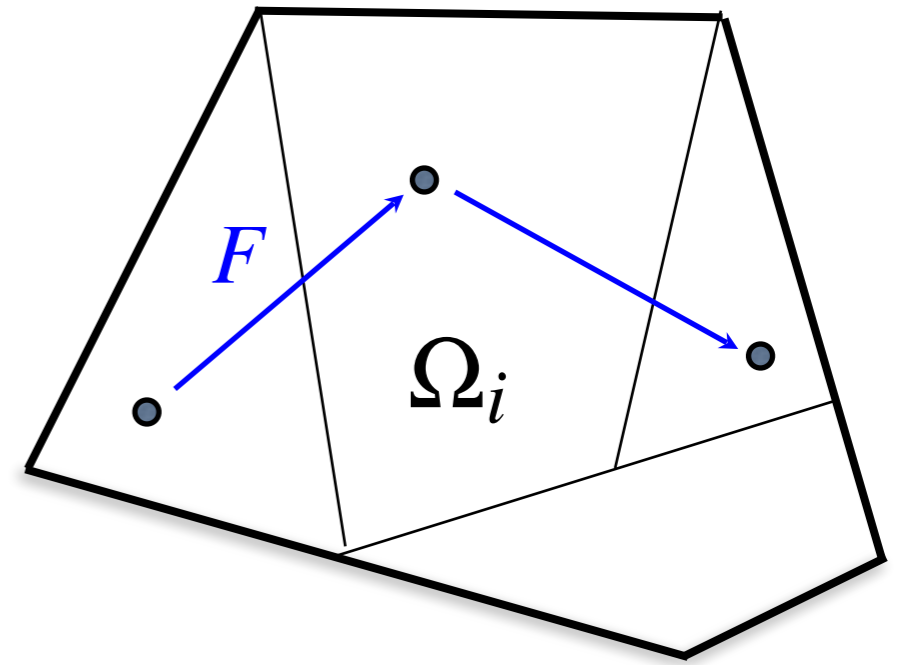
*The topological entropy of a piecewise isometry is zero.*

# Cells



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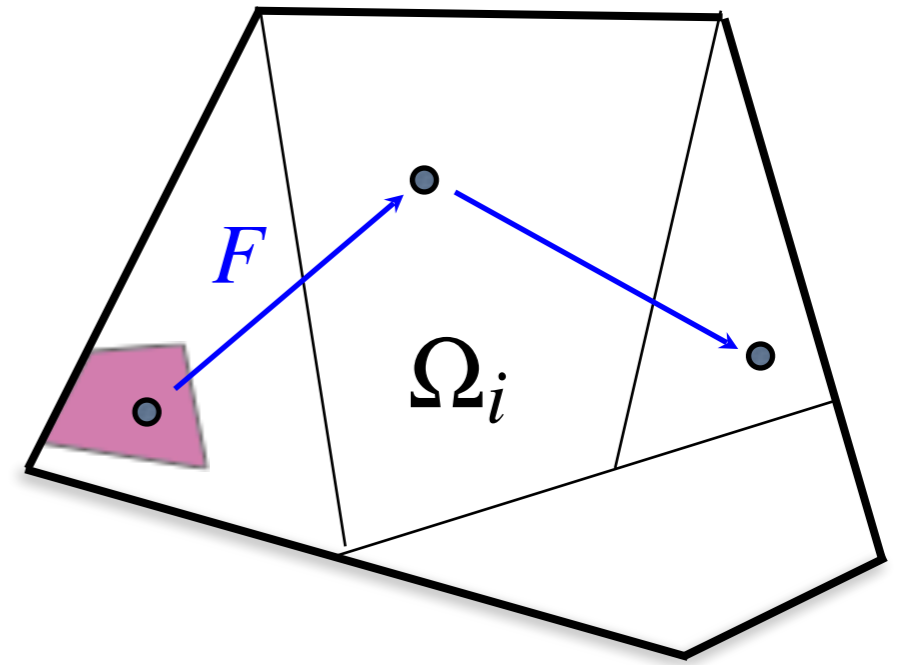
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A **cell** is a set of points with the same symbolic dynamics; cells are convex sets.

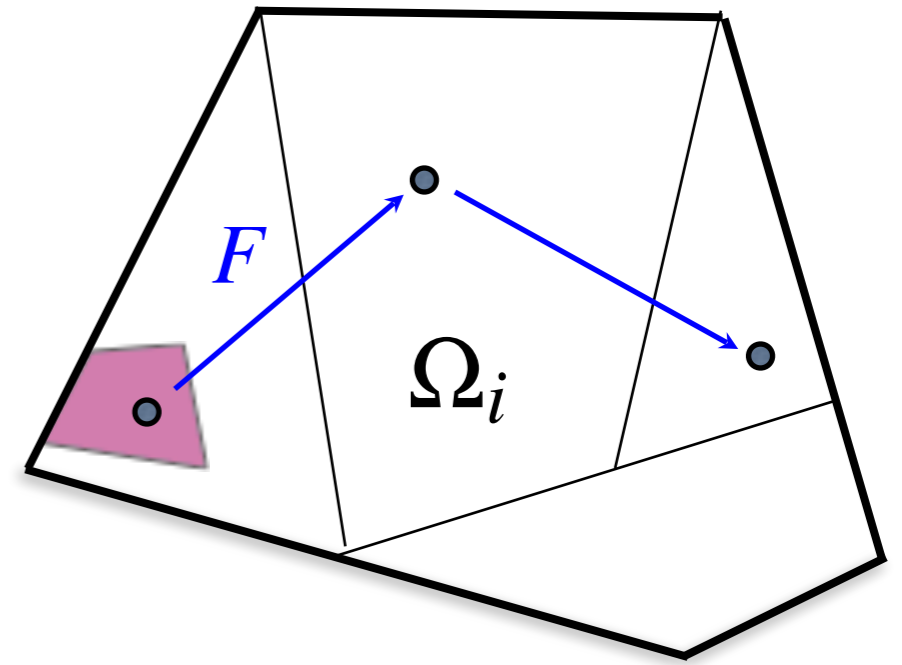




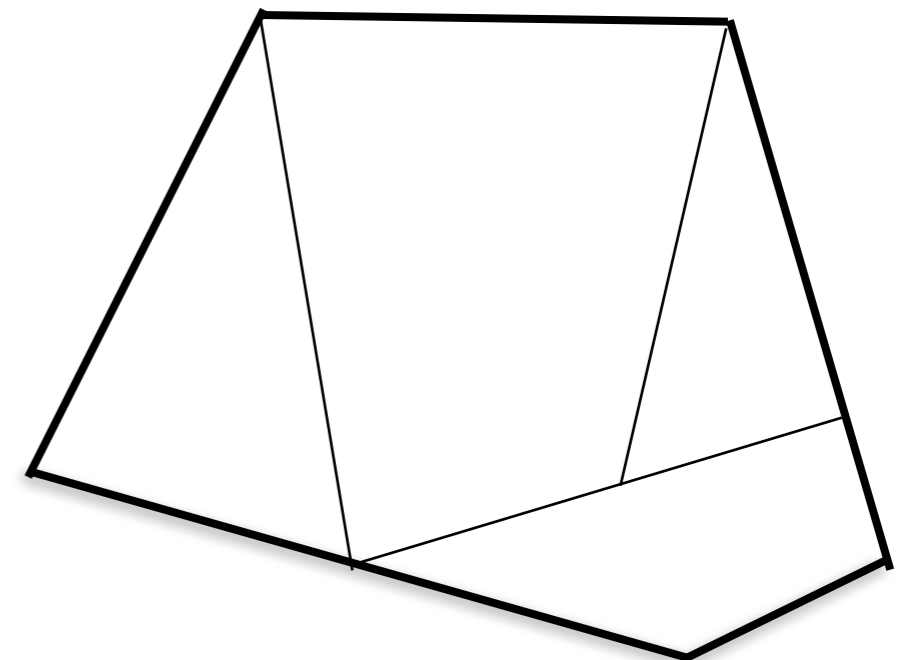
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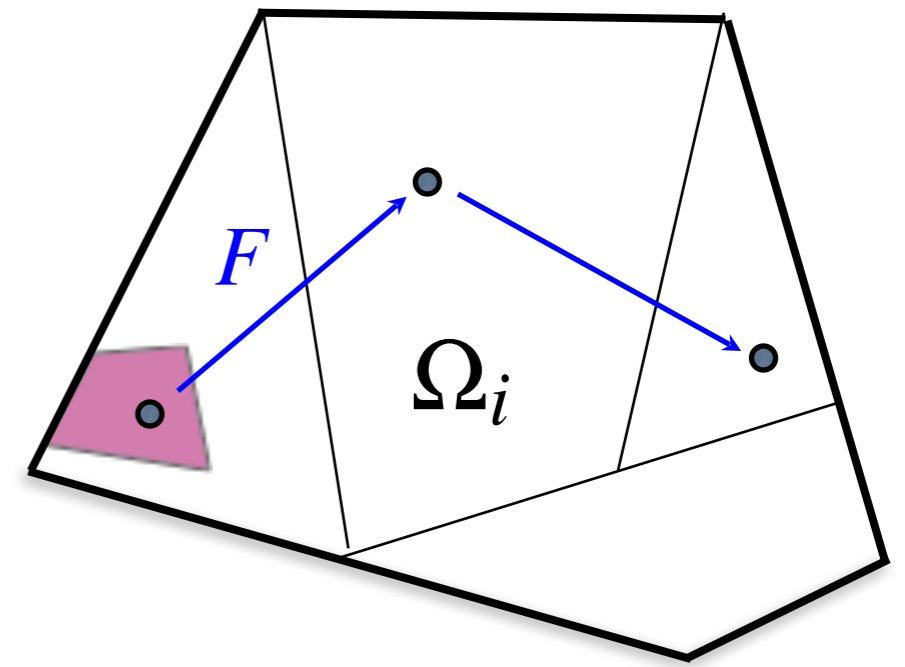
## Induced maps: the basis of renormalization



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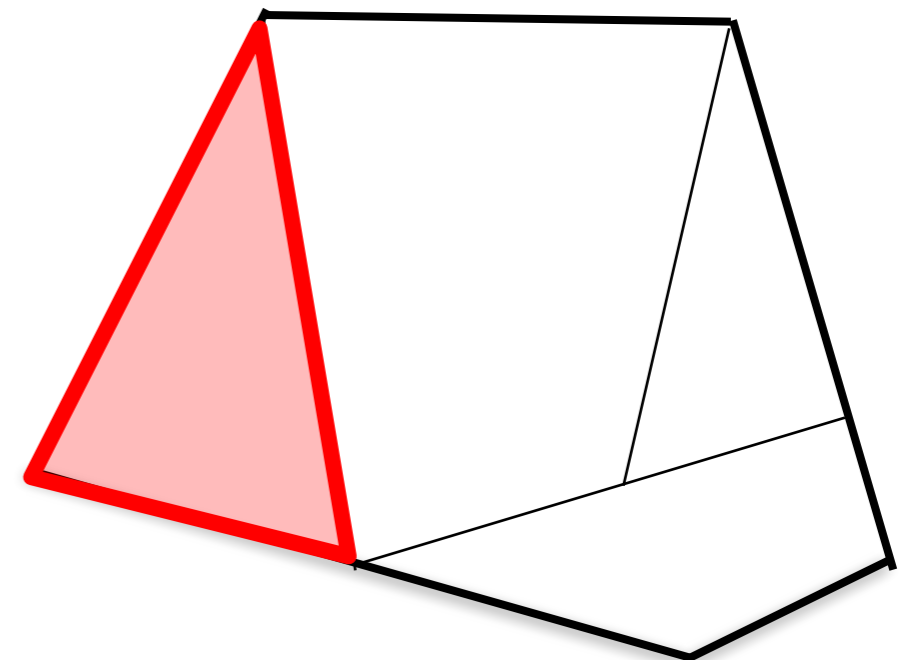
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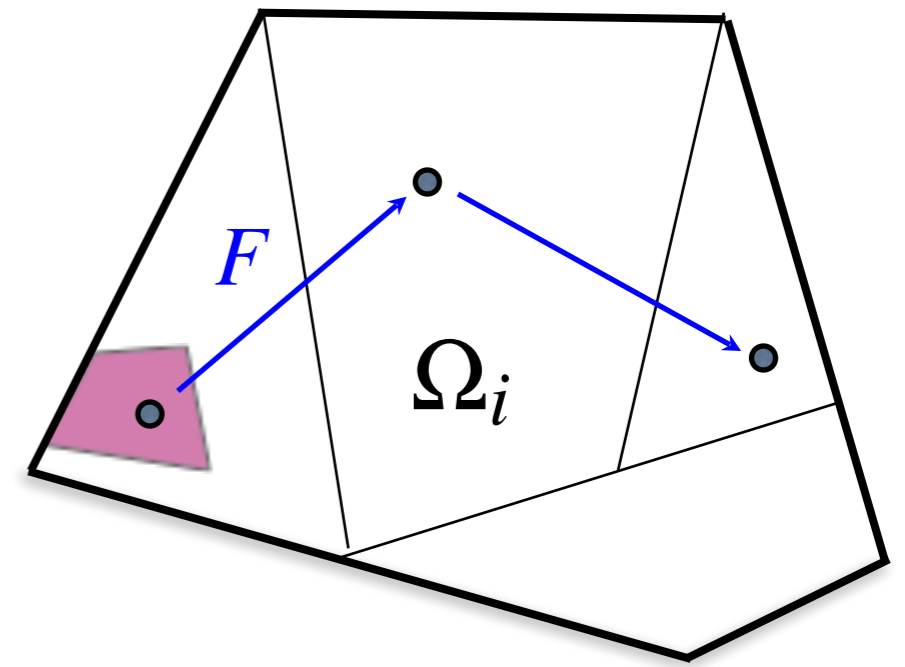
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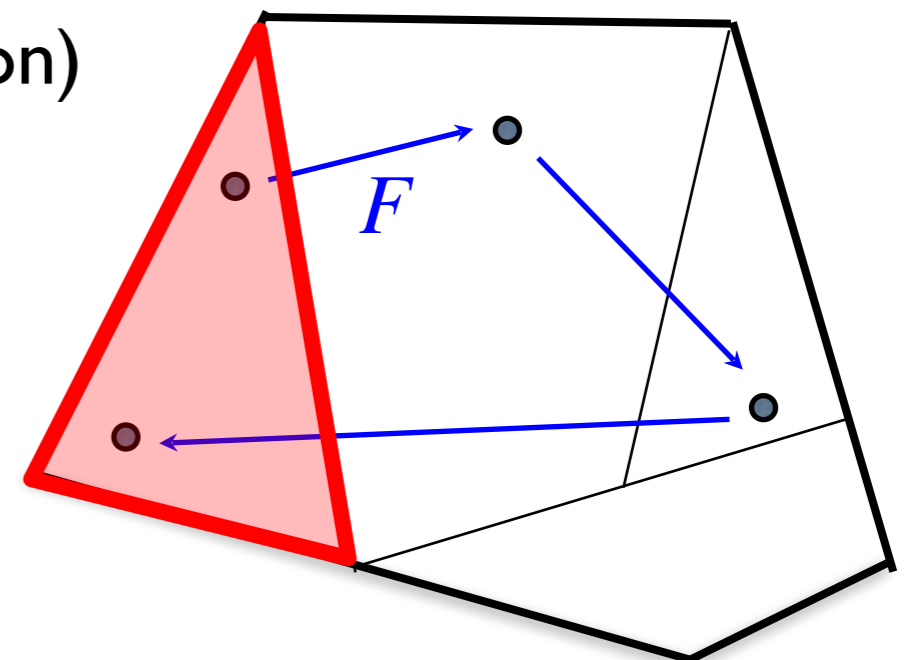
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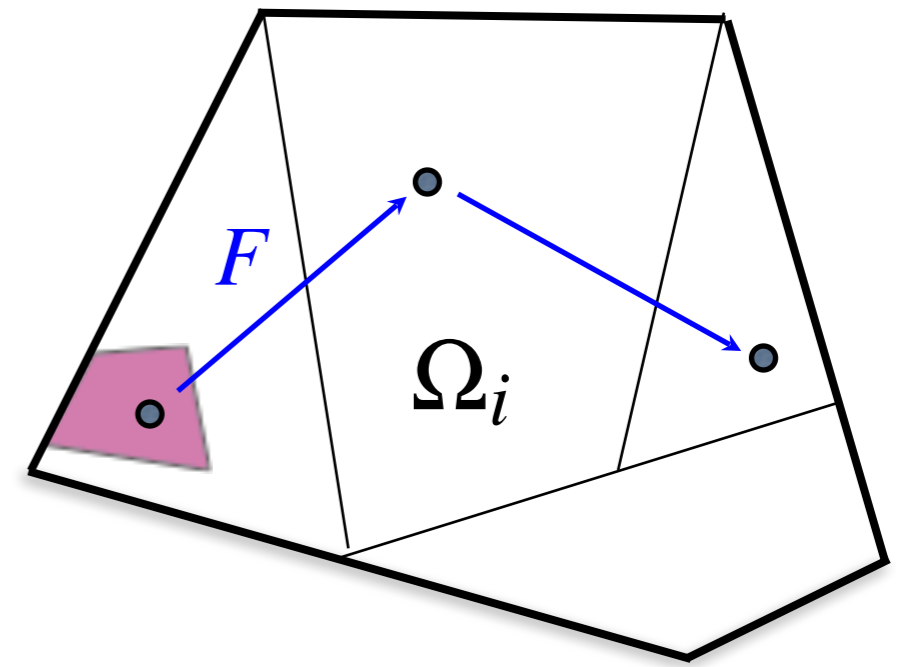
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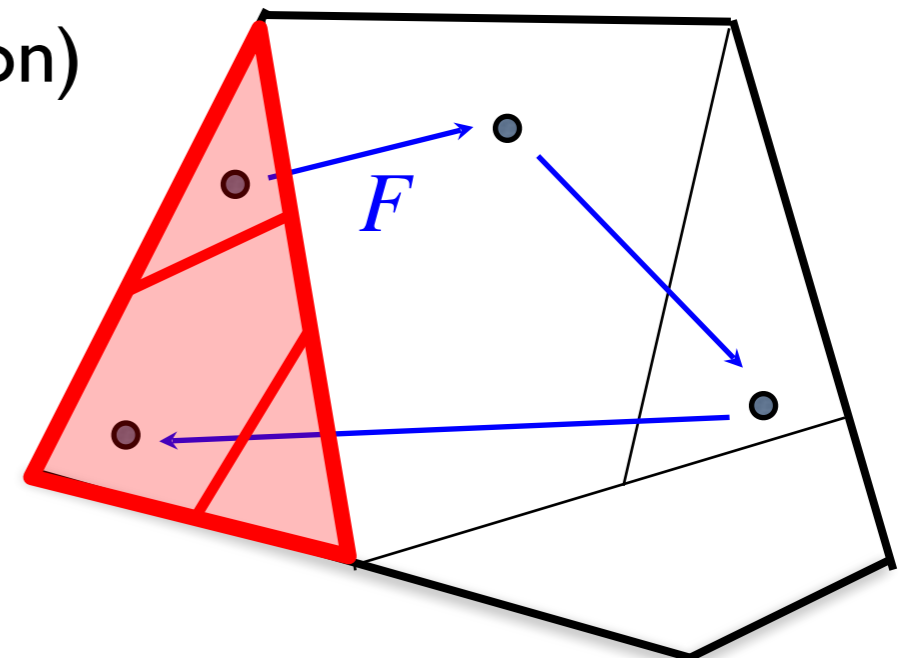
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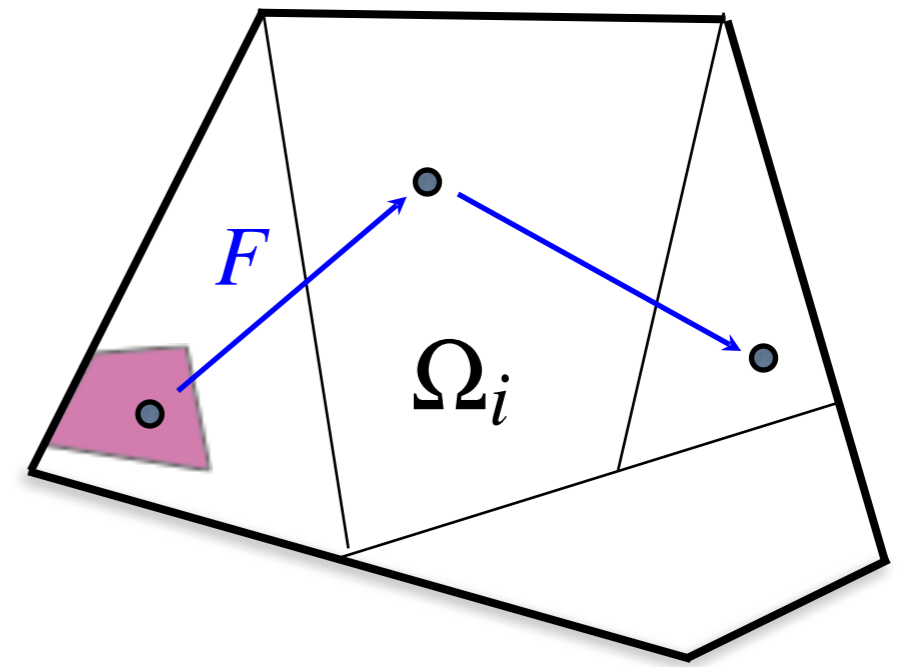
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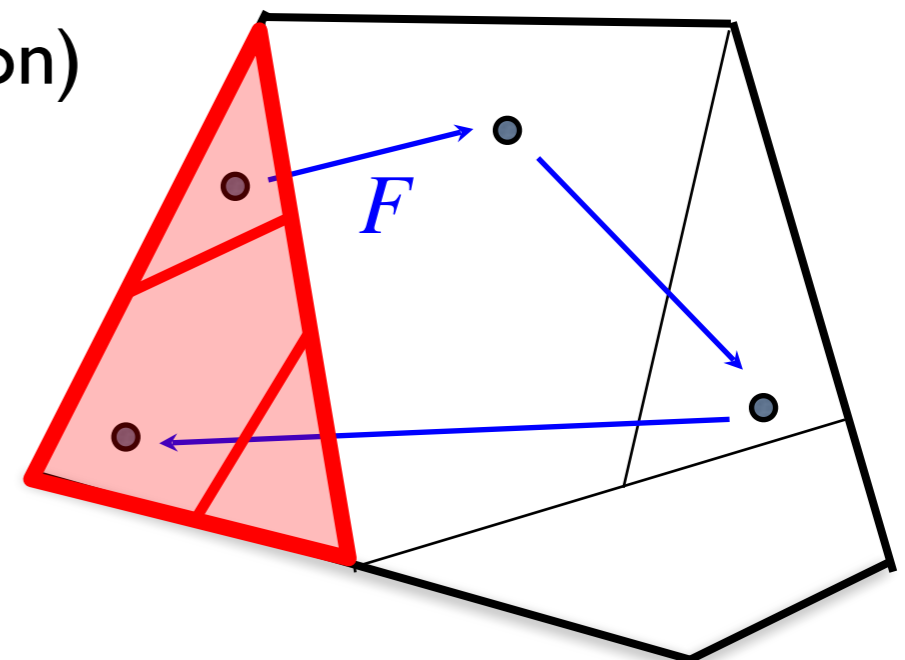
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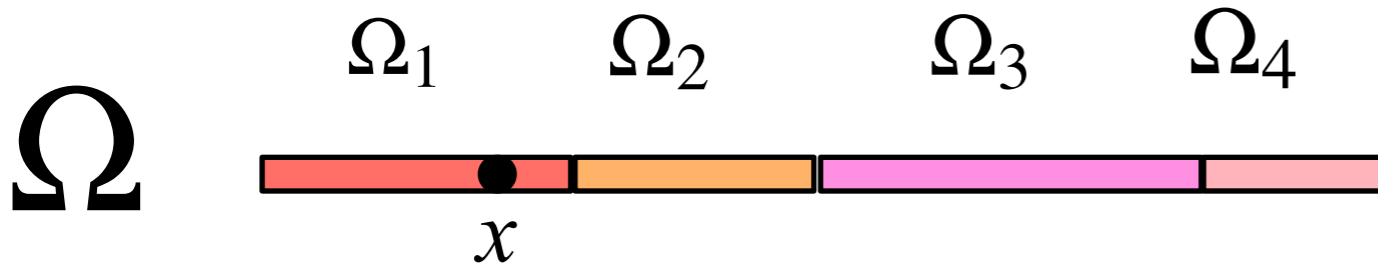


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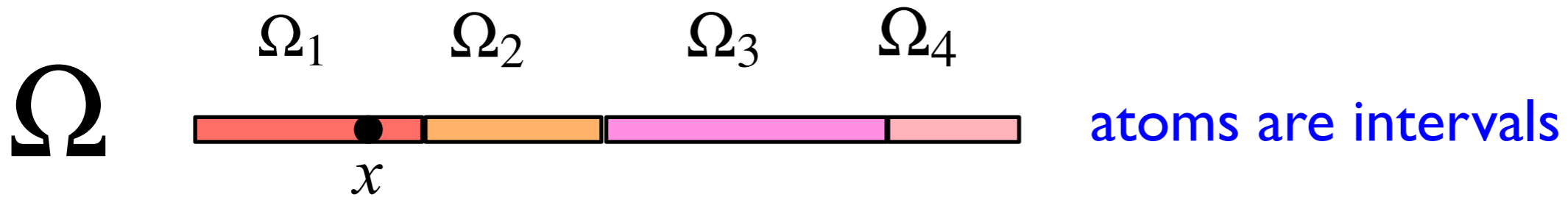
- Choose a domain (typically, an atom or a union of atoms).
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- A new PWI on a smaller domain.
- Different return times in different atoms.



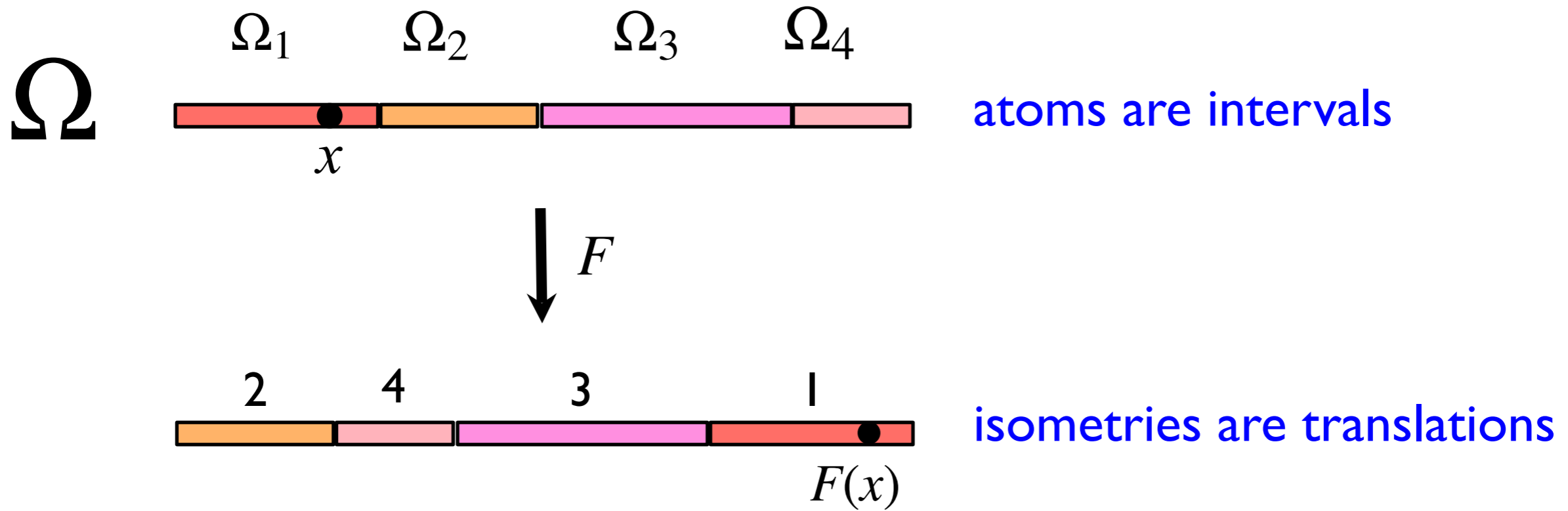
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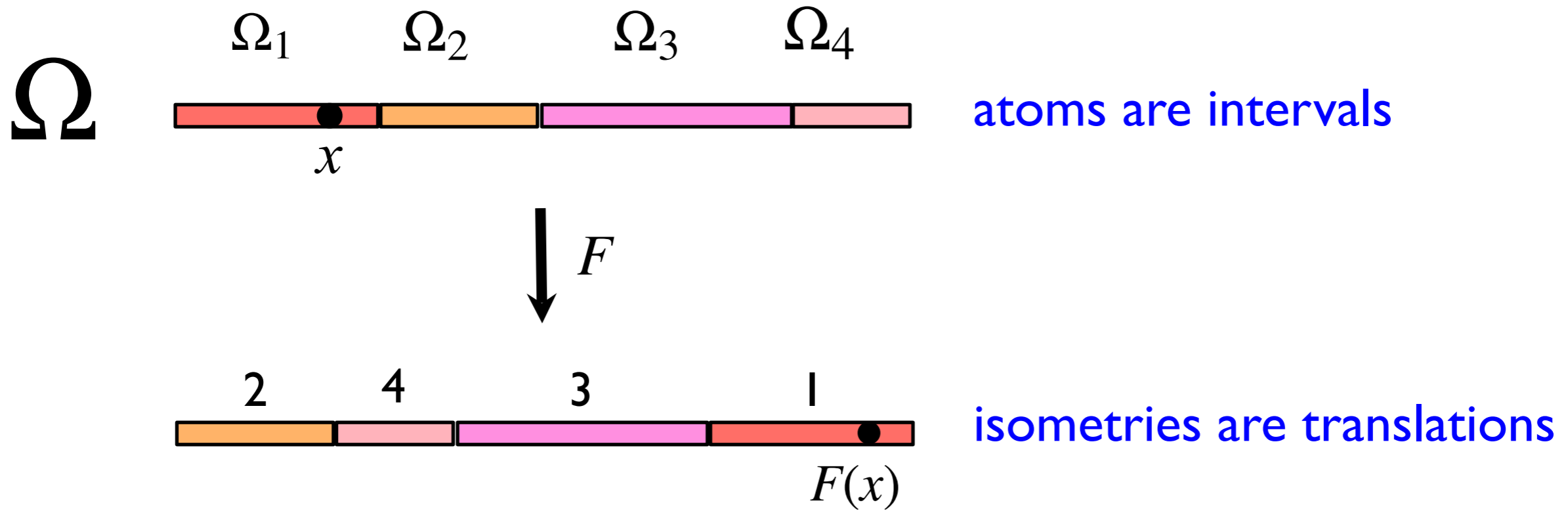


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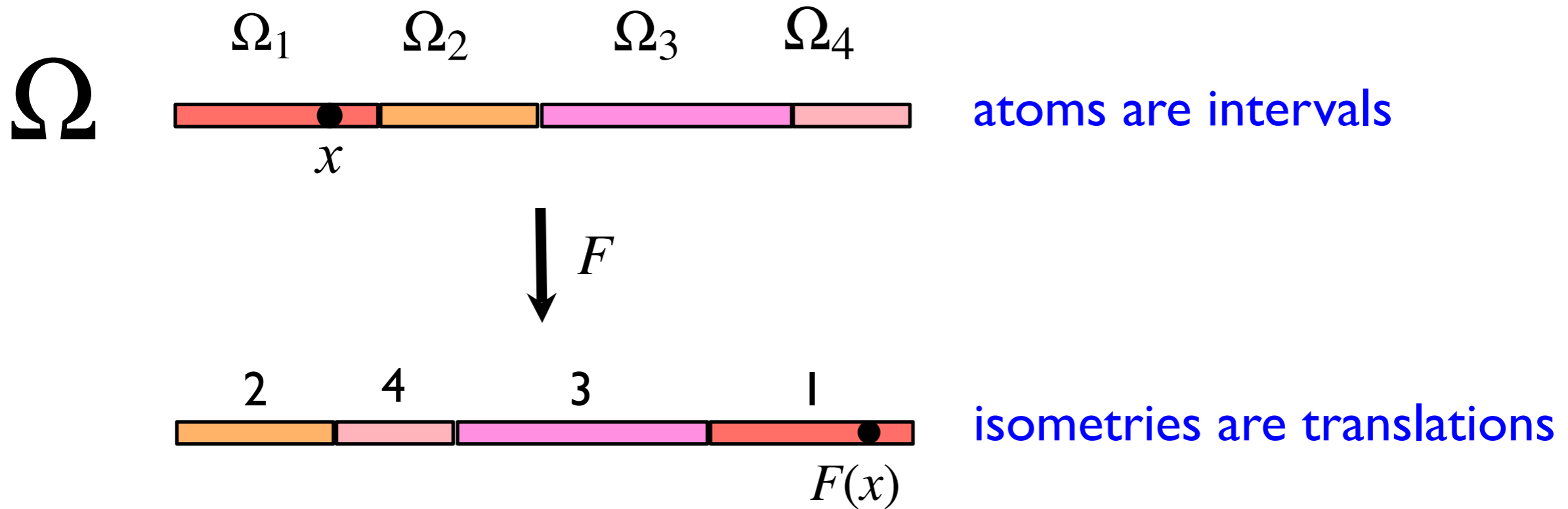
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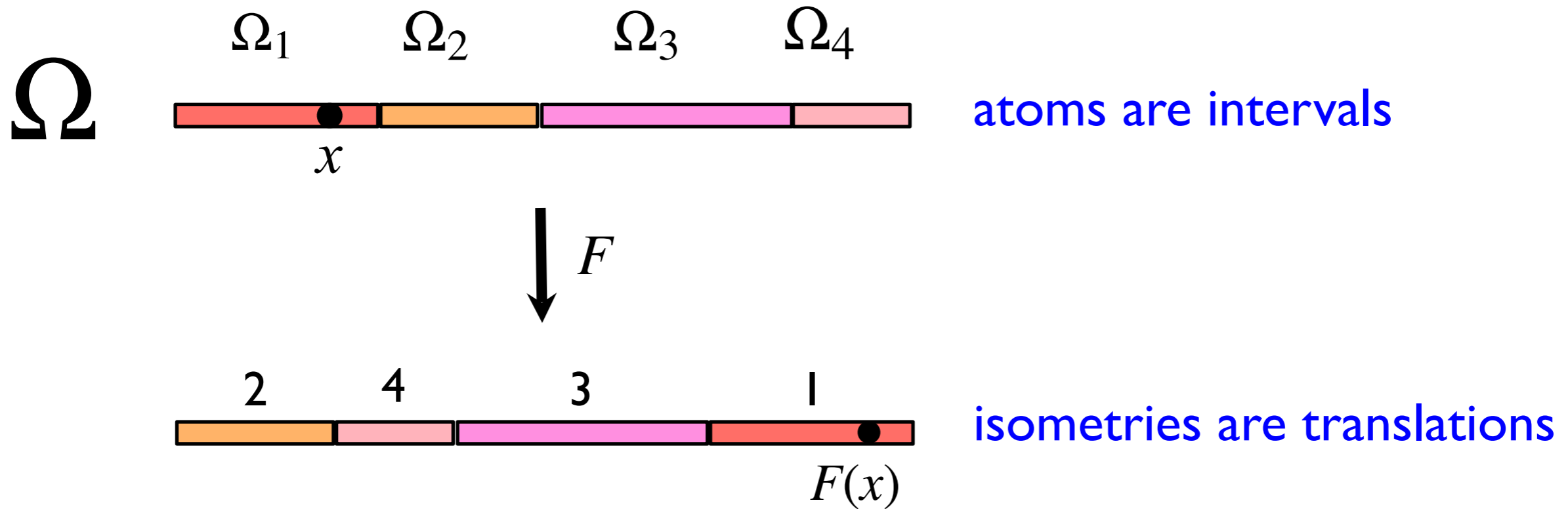
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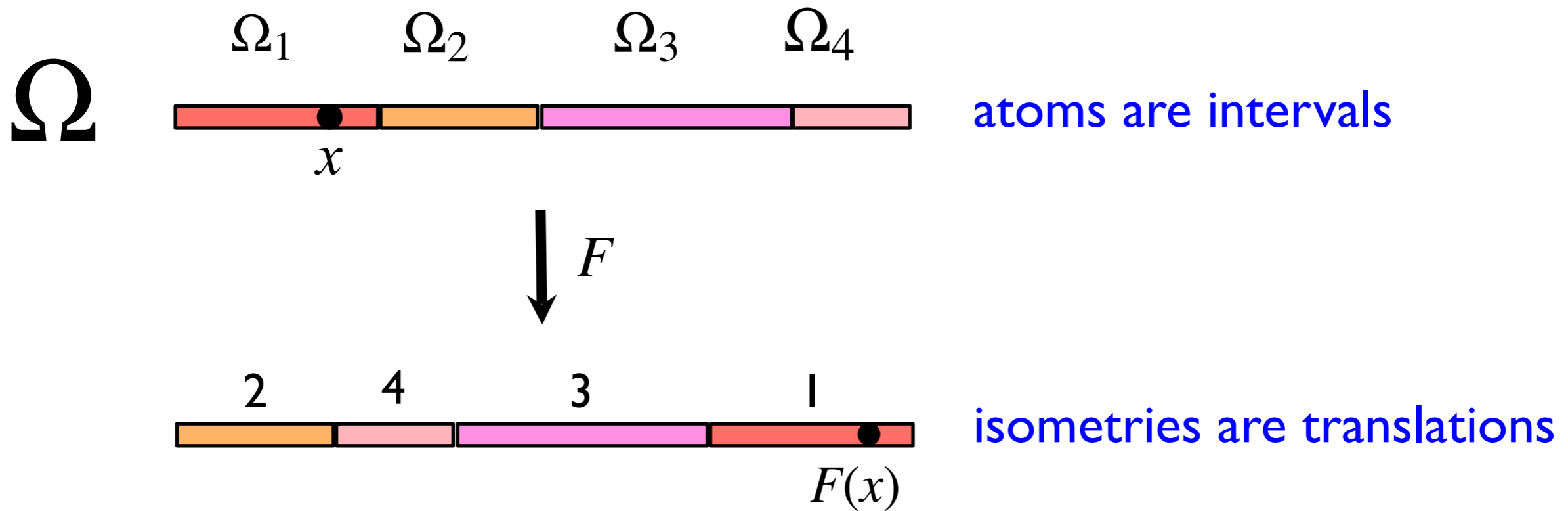


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$$K = \mathbb{Q}(|\Omega_1|, \dots, |\Omega_n|)$$

**Theorem** (Boshernitzan & Carrol, 1997)

If an IET is defined over a quadratic field, then, up to scaling, the number of induced maps over sub-intervals is finite.

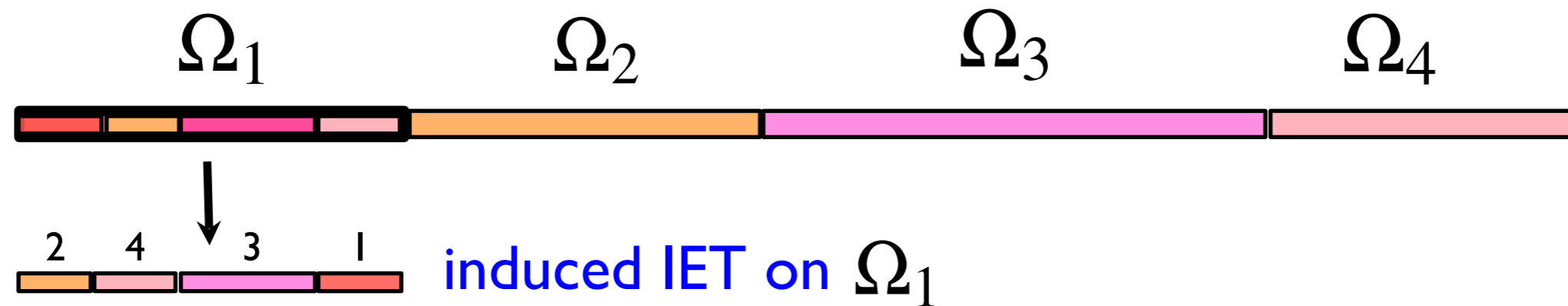
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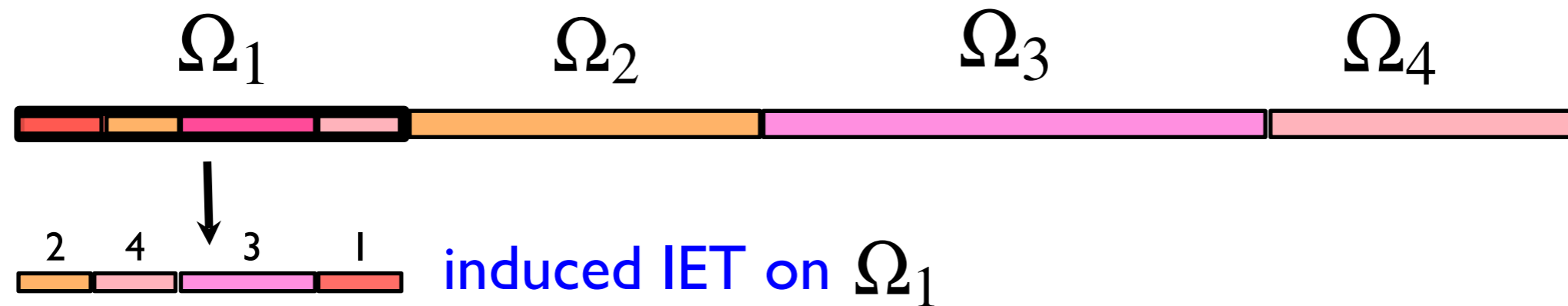
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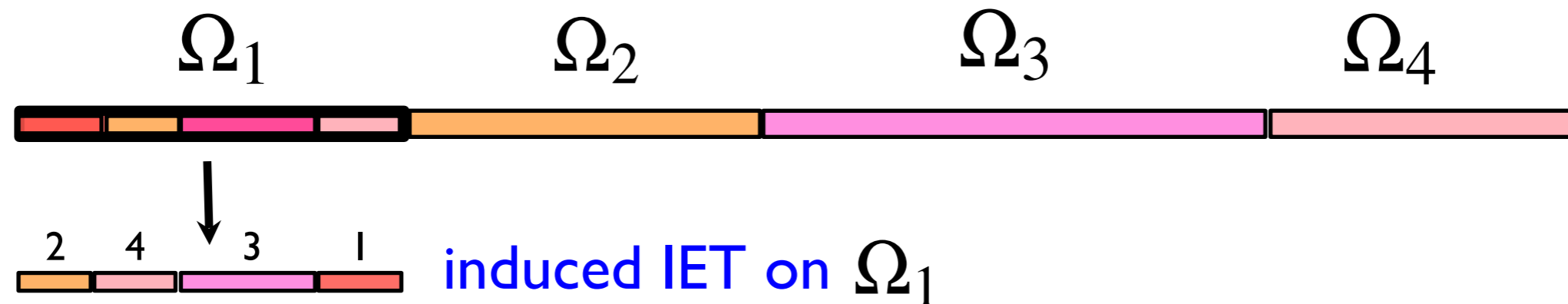


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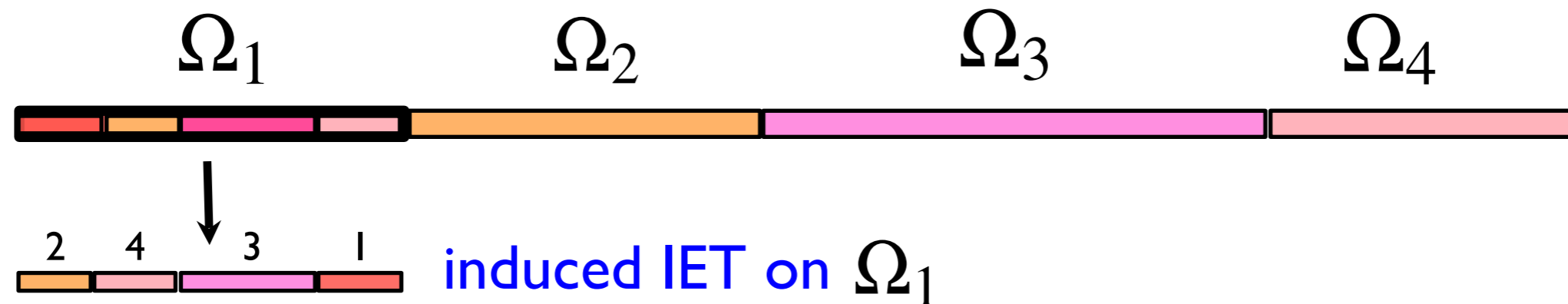


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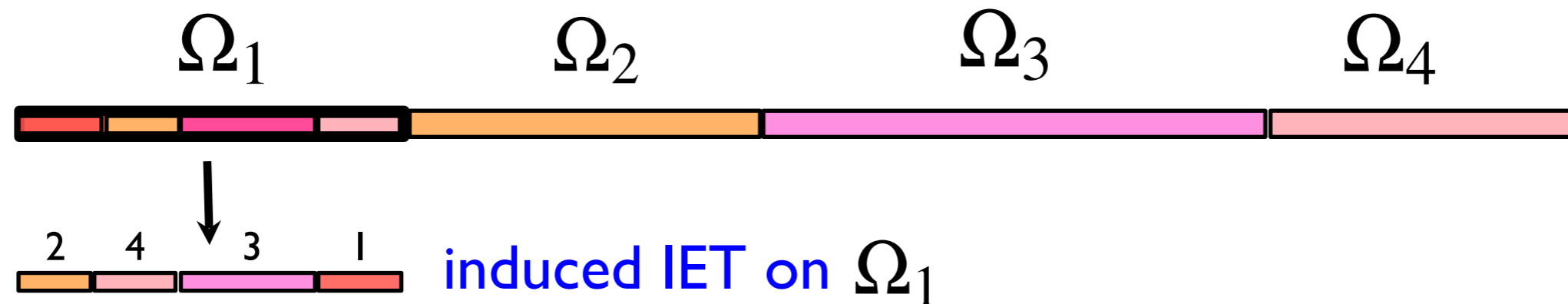
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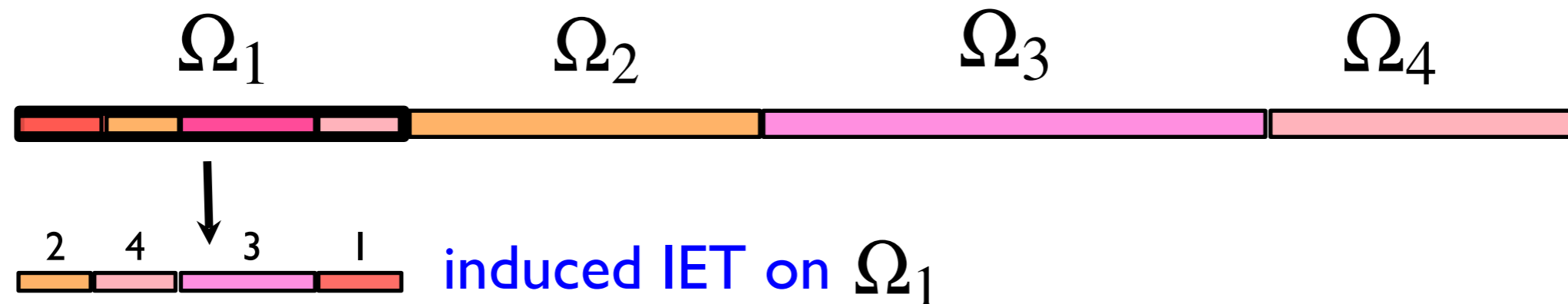
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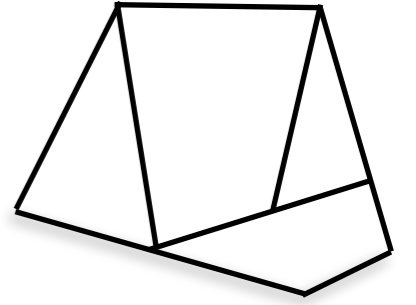
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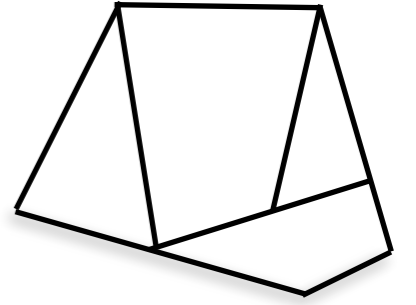
- We construct two one-parameter families of maps which are renormalizable iff the parameter belongs to a distinguished quadratic field.
- With two parameters, we only find a degenerate form of renormalizability (one parameter is free).

# Higher dimensions: topology

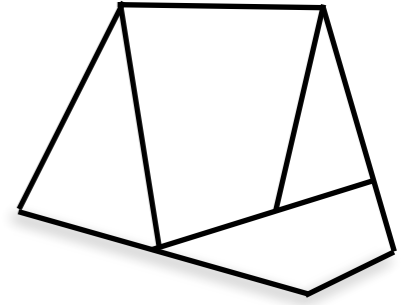


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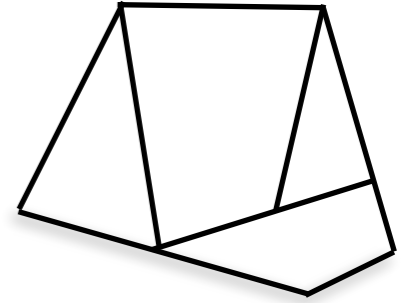


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discontinuity set

$$\mathcal{D} = \bigcup_{t \in \mathbb{Z}} F^t(\partial\Omega)$$

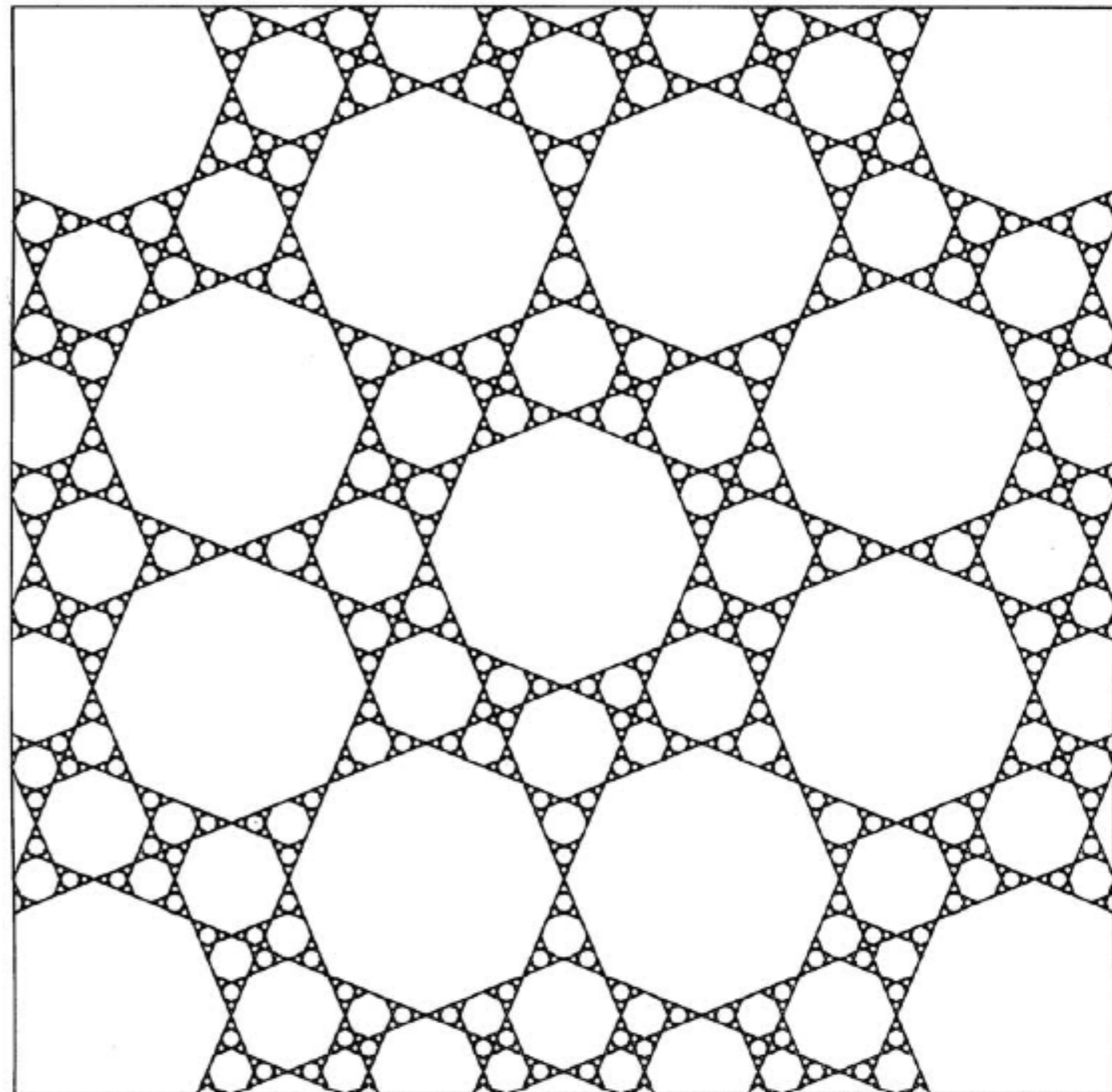
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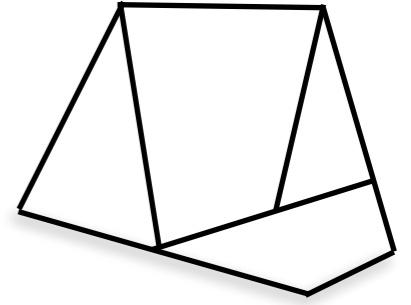
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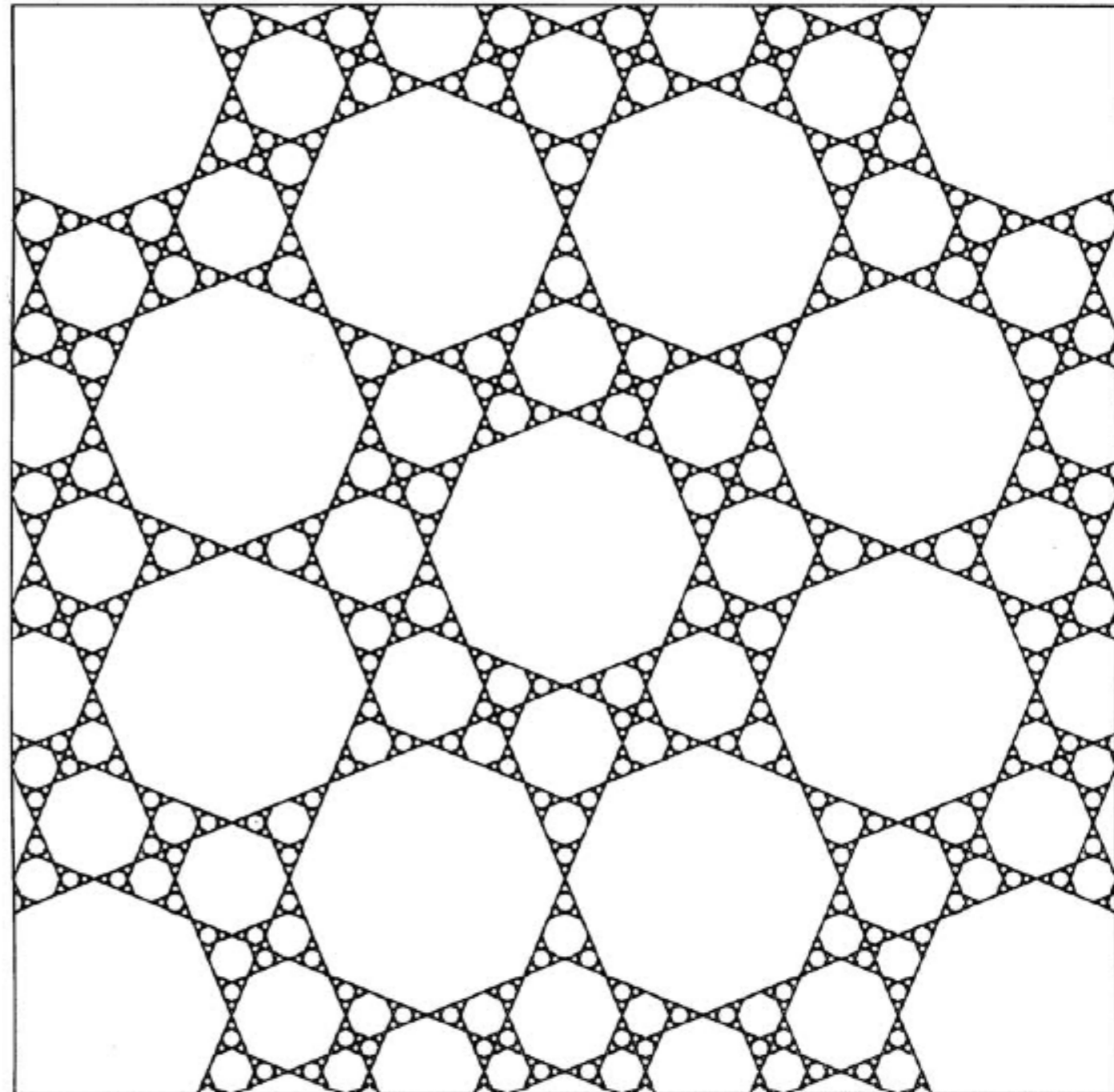


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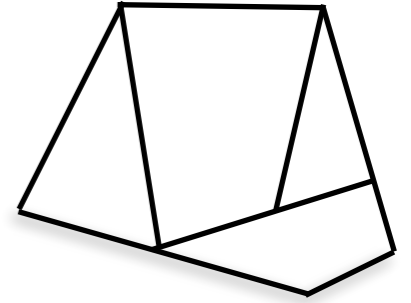
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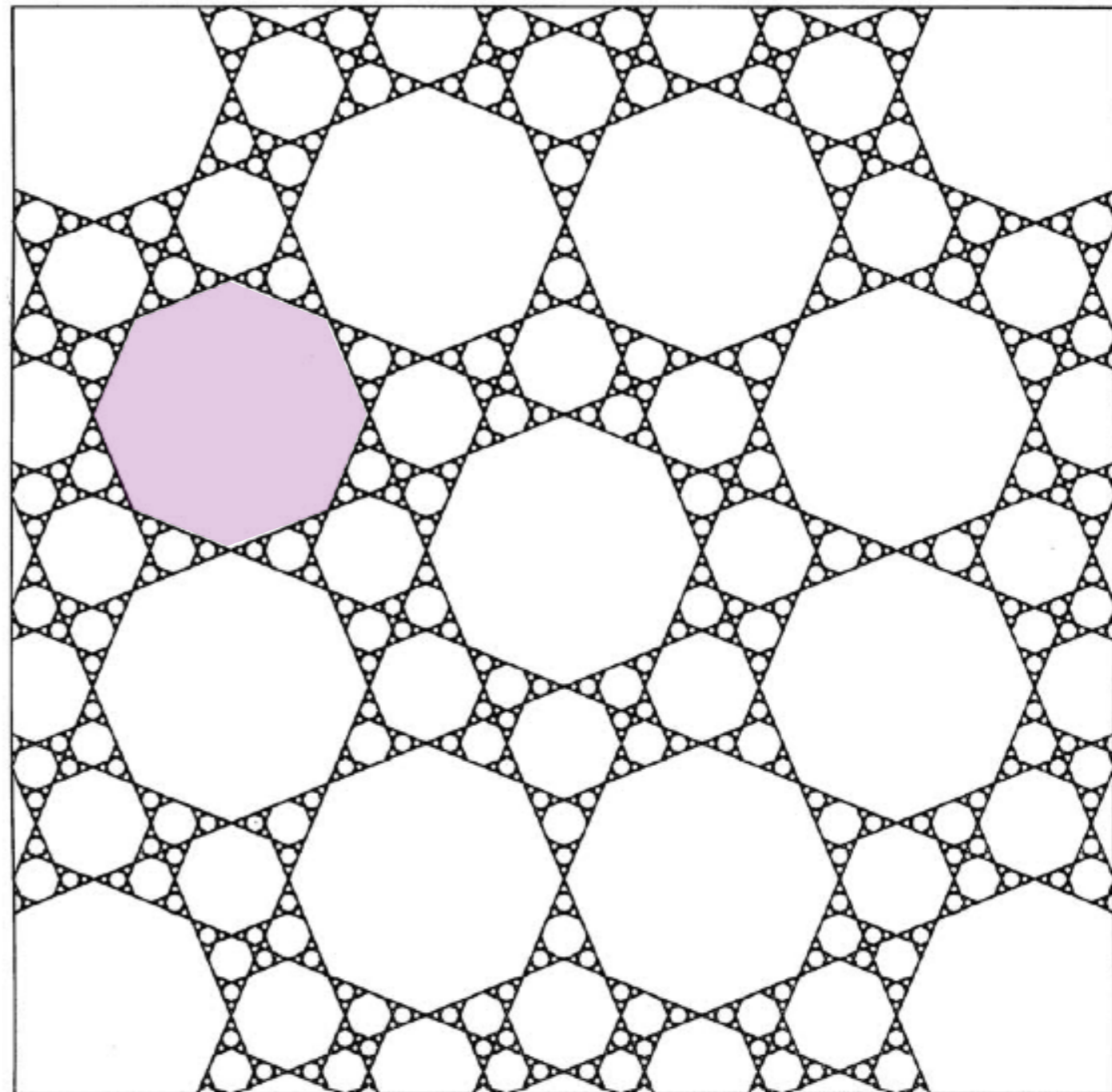
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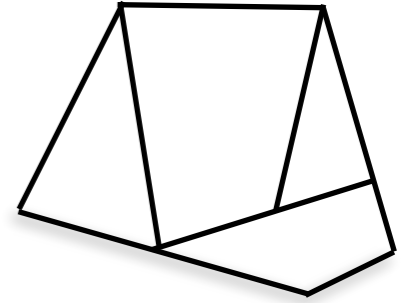
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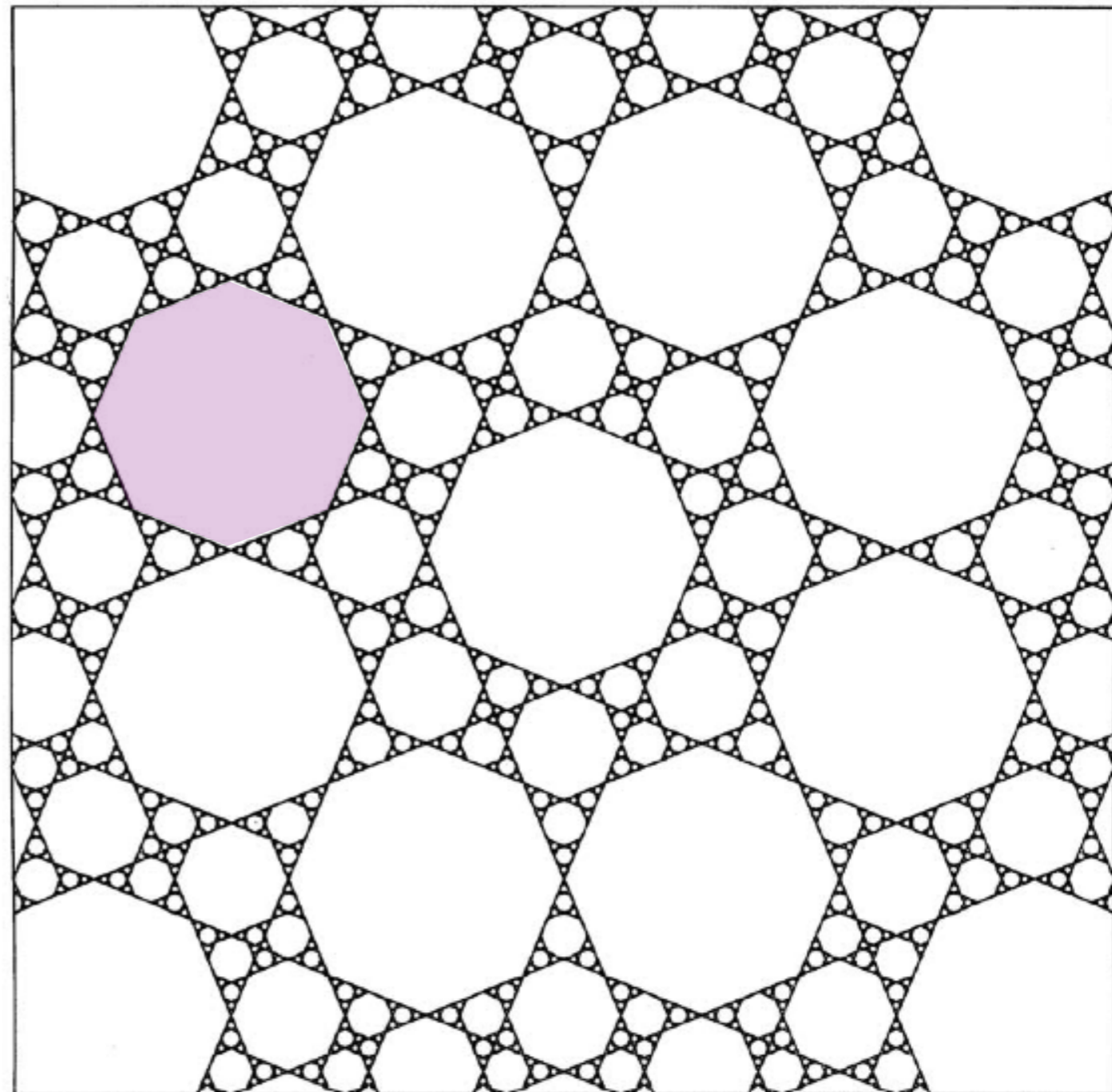
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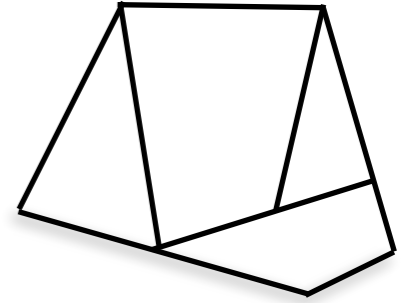
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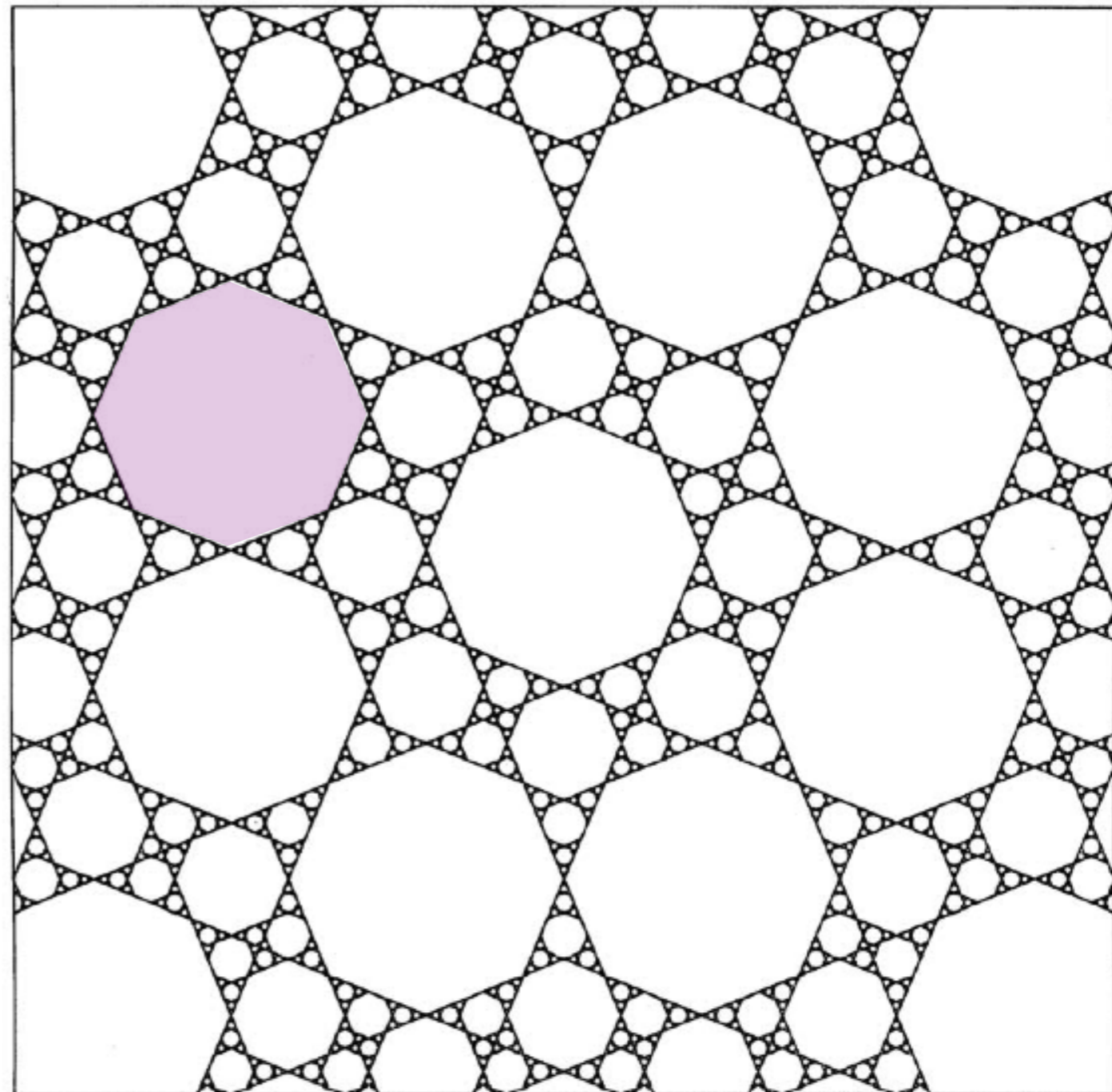
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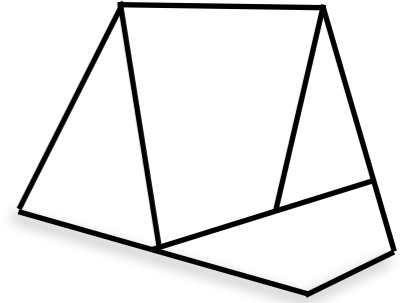
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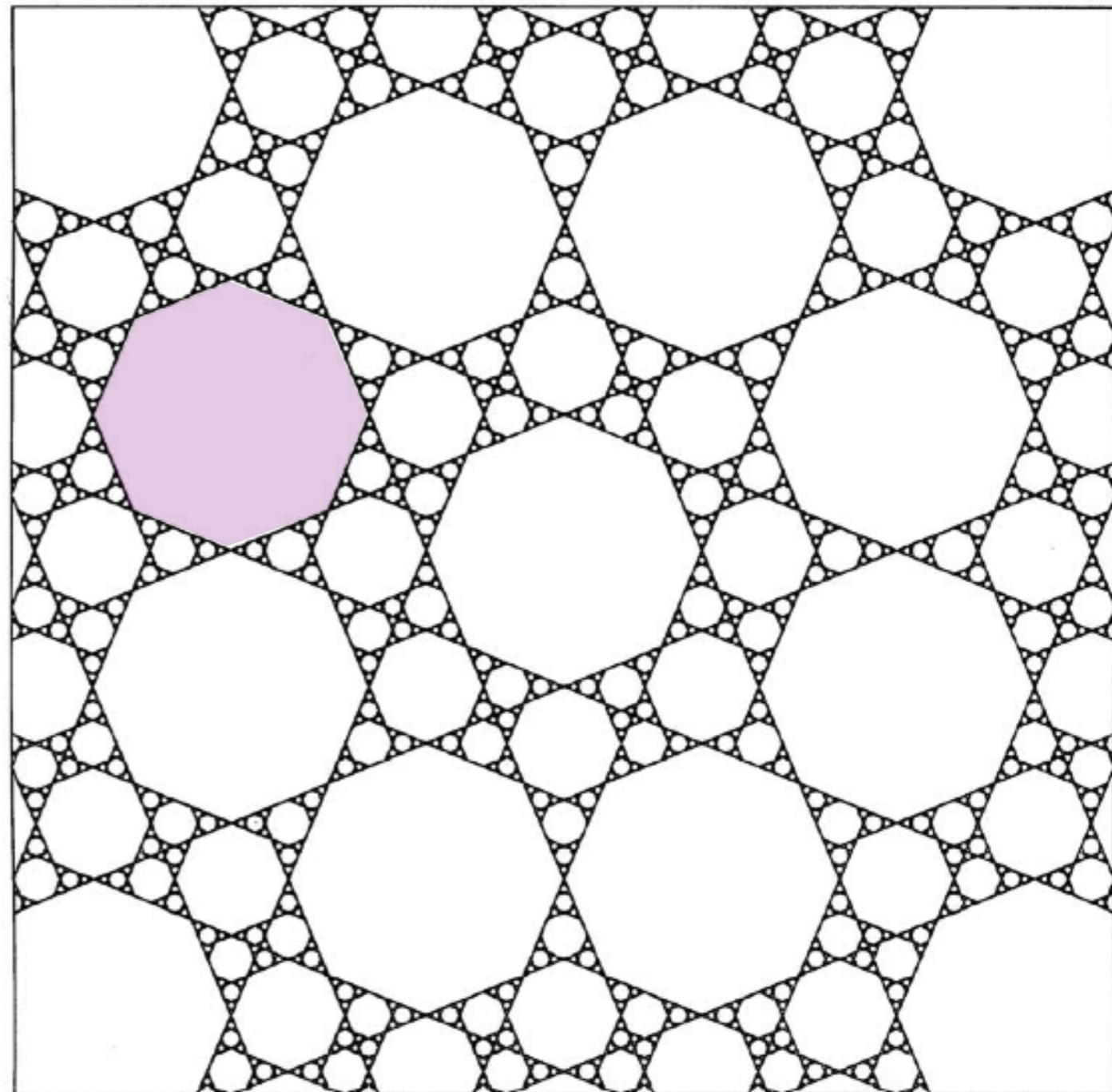
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(asymptotic phenomena)



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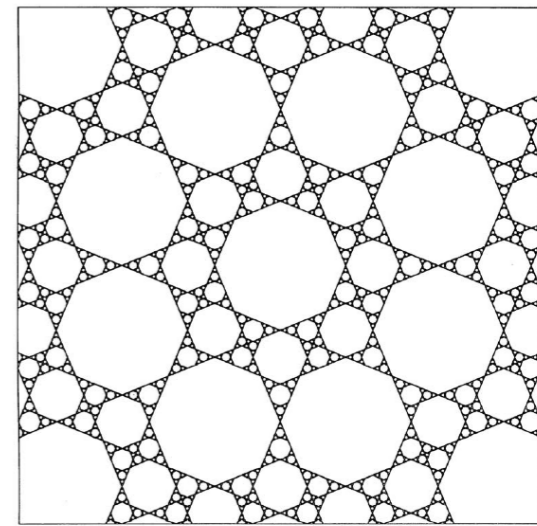
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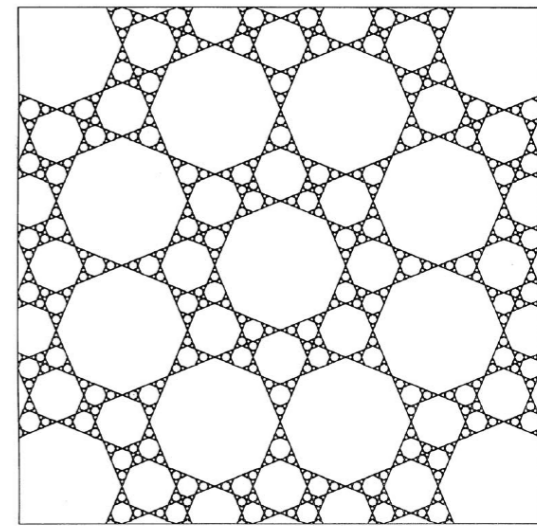
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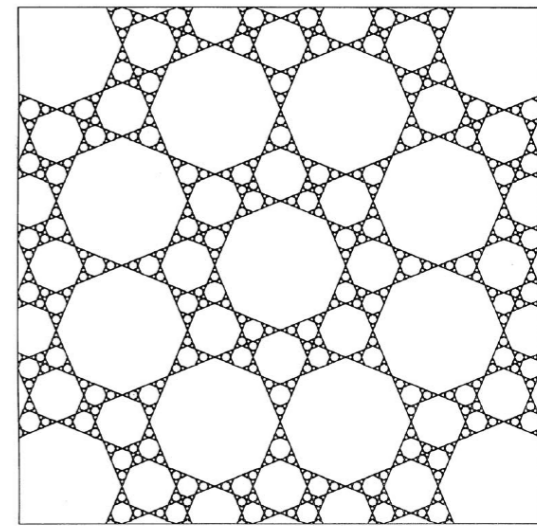
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No general theory, but many examples of quadratic PWIs, where:

- the number of induced maps is finite, up to scaling;
- the periodic set has full measure: tiling by regular polygons;





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The rotational component of a PWI is conjugate to  $\begin{pmatrix} \lambda & -1 \\ 1 & 0 \end{pmatrix}$

The rotation field is  $\mathbb{Q}(\lambda)$

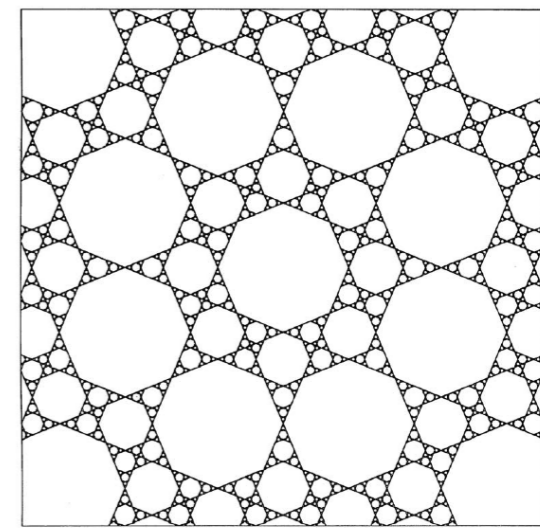
$$\lambda = 2 \cos(2\pi \nu) \leftarrow \text{rotation number}$$

Simplest case: rational rotation number.

$$\lambda = 2 \cos(2\pi p/q) \quad \begin{array}{ll} q = 1,2,3,4,6: & \mathbb{Q}(\lambda) \quad \text{rational} \\ q = 5,8,10,12: & \mathbb{Q}(\lambda) \quad \text{quadratic} \end{array}$$

No general theory, but many examples of quadratic PWIs, where:

- the number of induced maps is finite, up to scaling;
- the periodic set has full measure: tiling by regular polygons;
- scaling constants are units in the ring of integers of the relevant field.

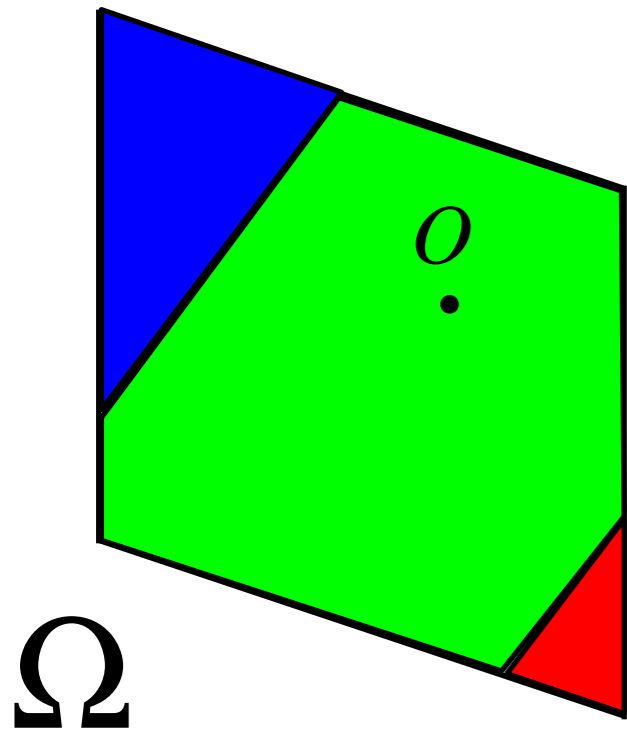


# One-parameter families of polygon-exchange transformations

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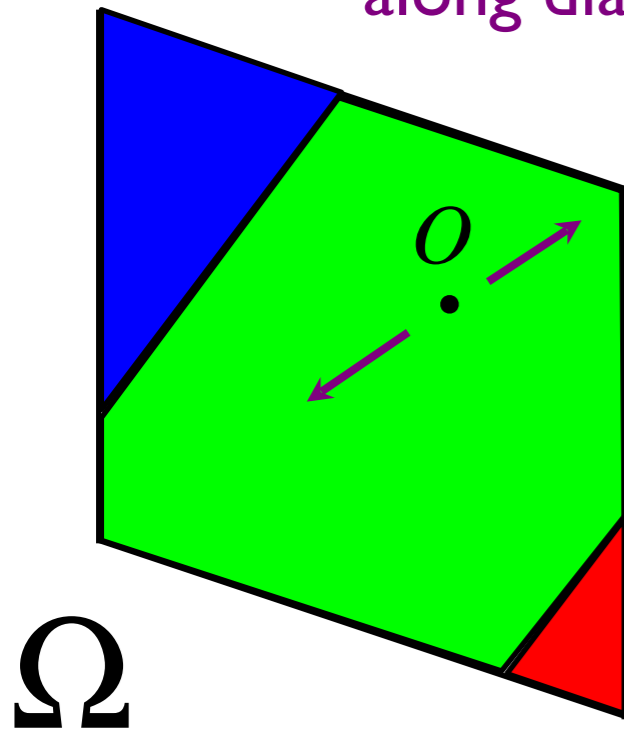
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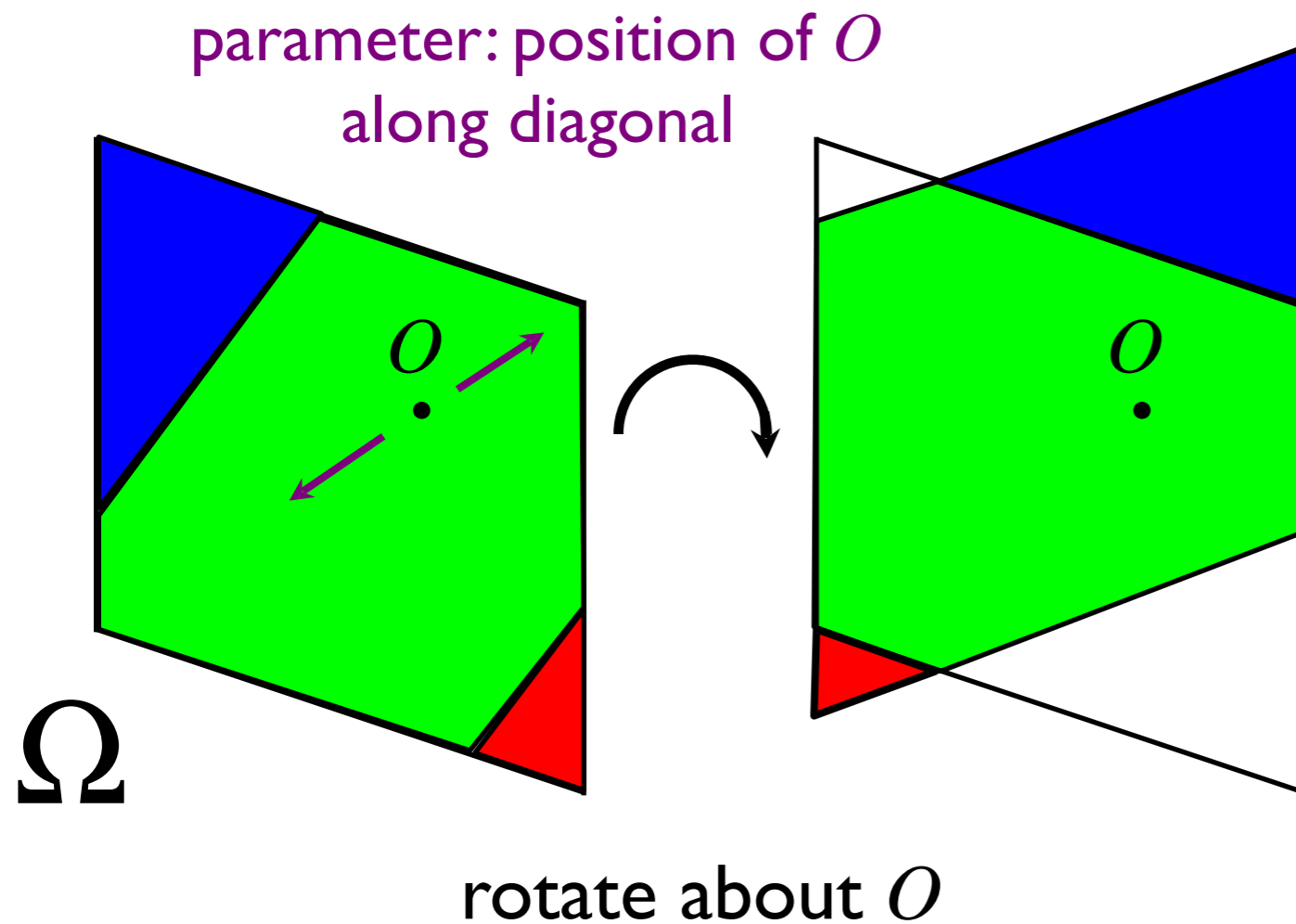
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parameter: position of  $O$   
along diagonal



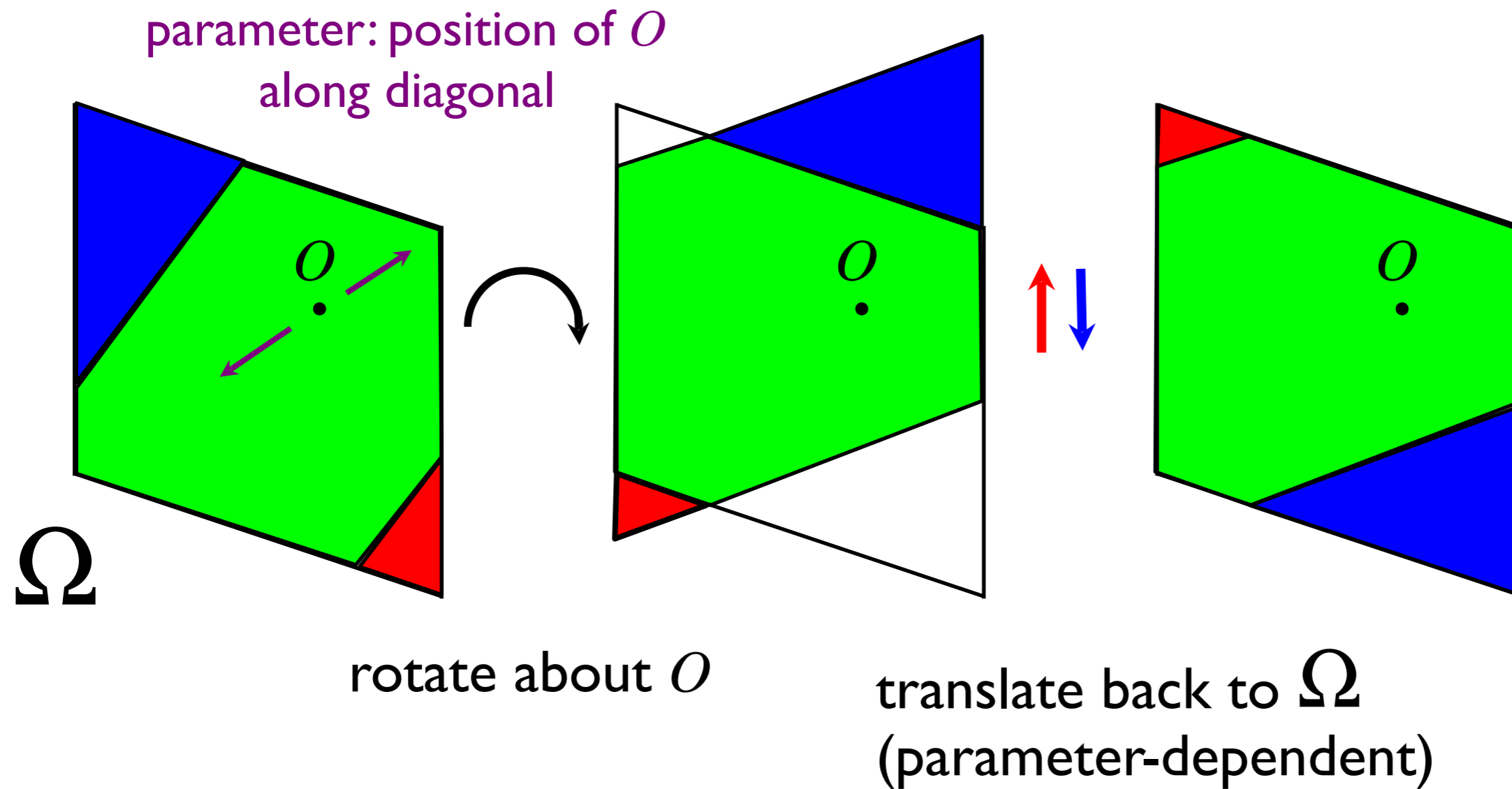
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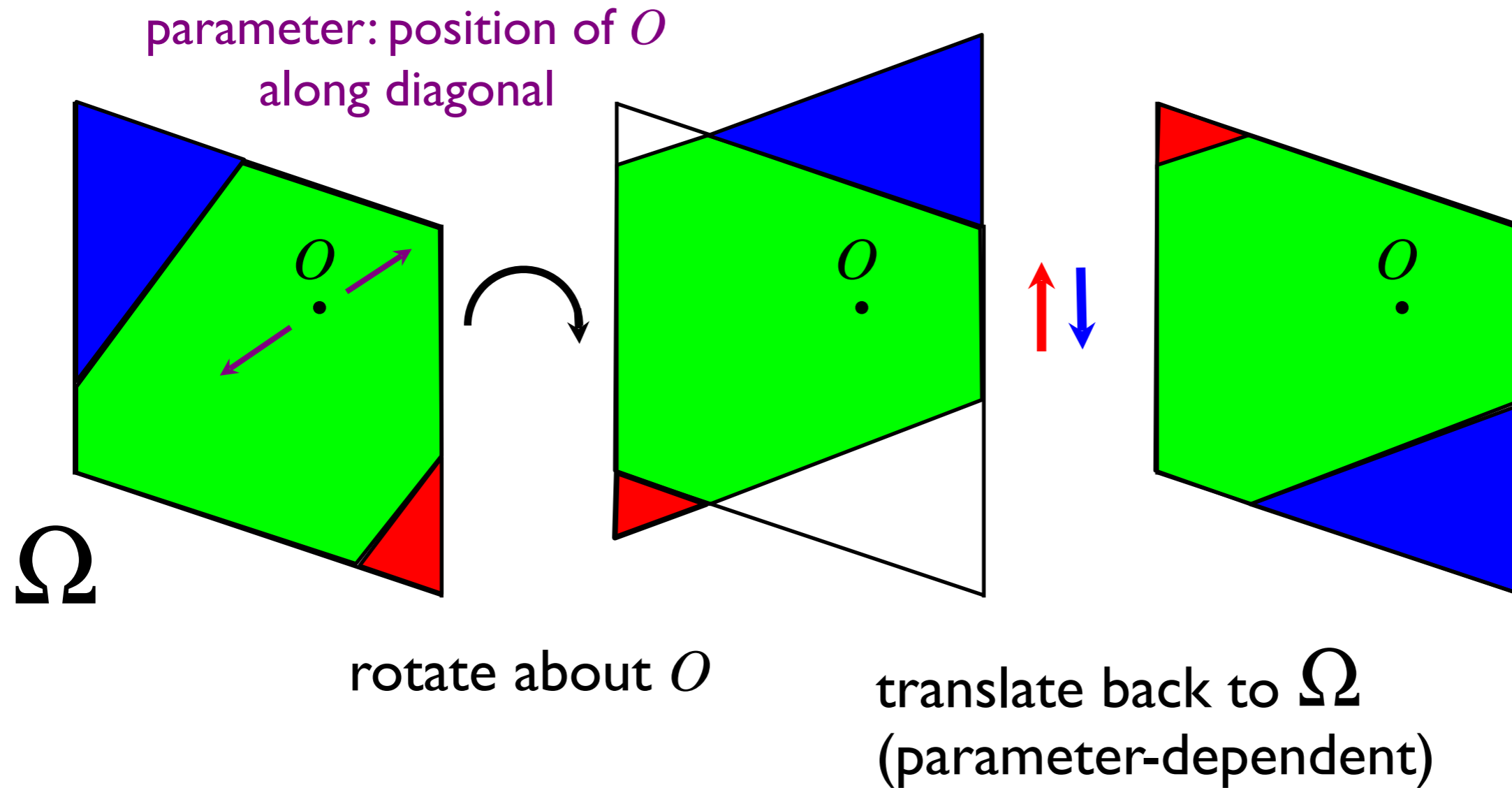
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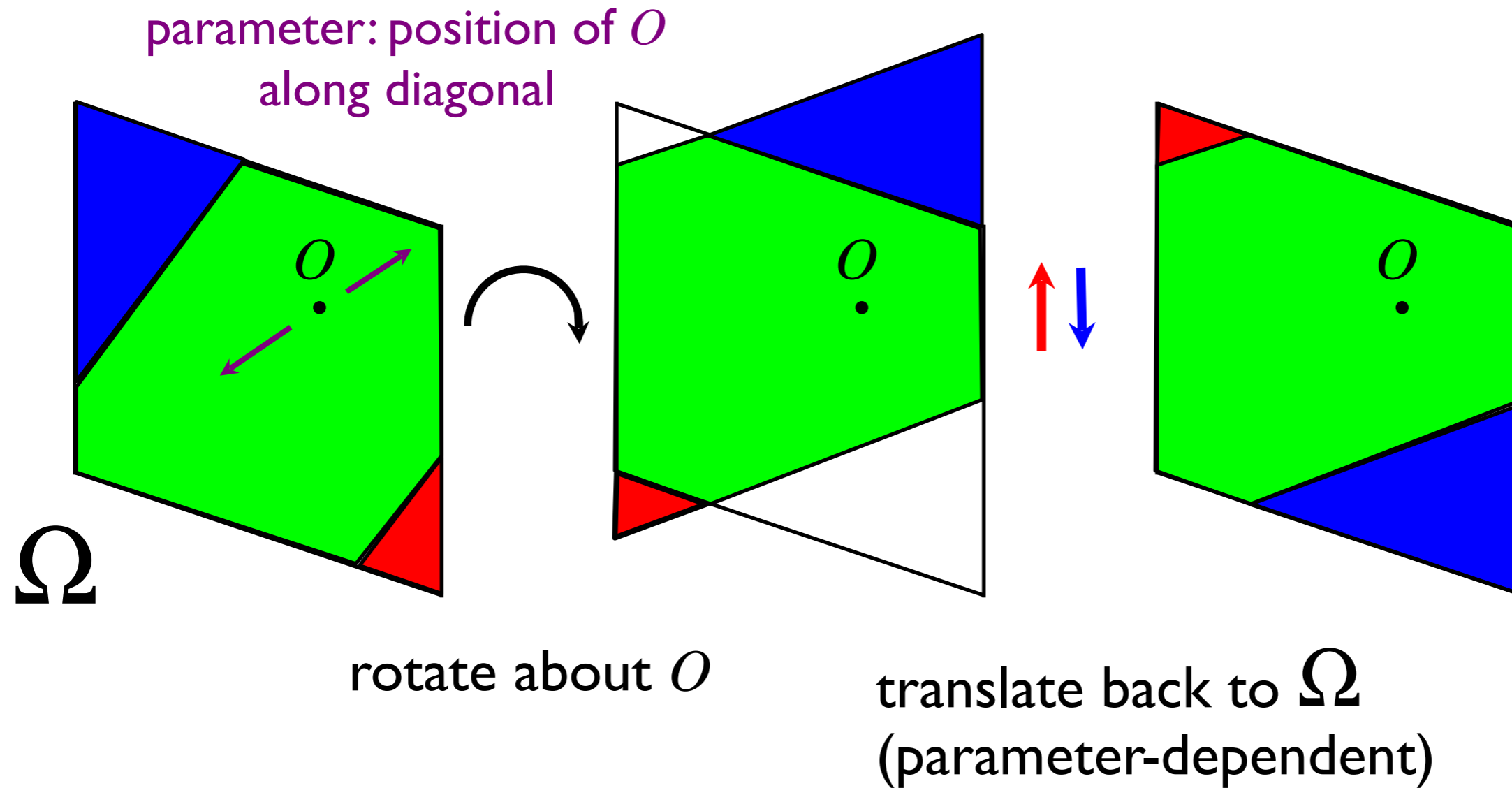
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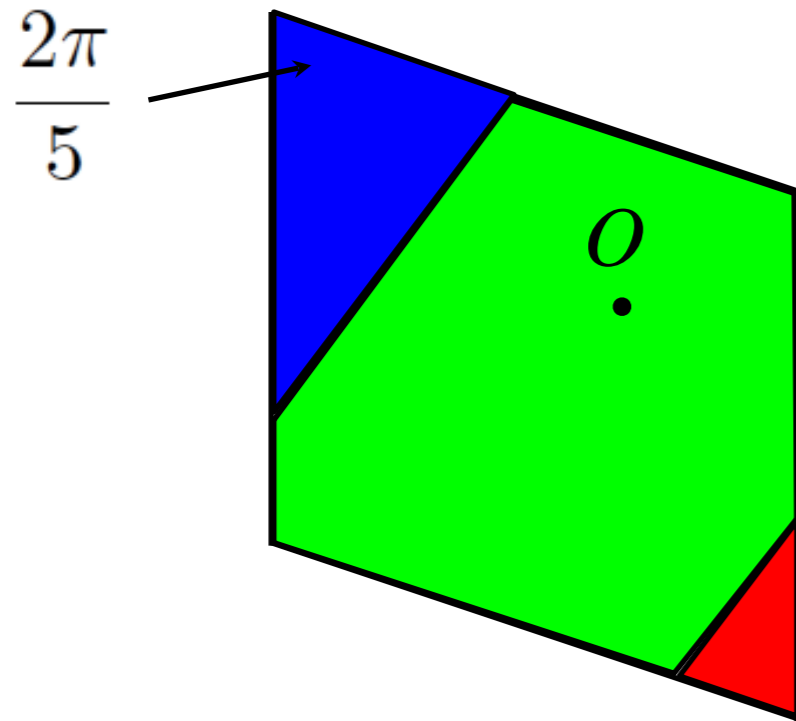


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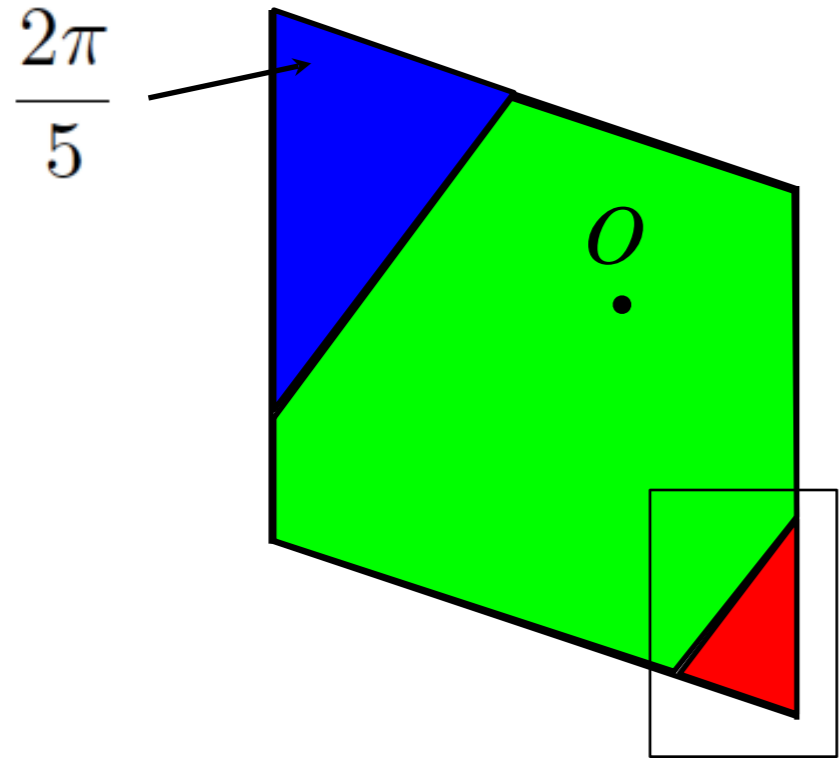
translation module:  $\mathbb{Q}(\lambda) + s\mathbb{Q}(\lambda)$   $s$ : parameter



# The pentagonal model (field $\mathbb{Q}(\sqrt{5})$ )

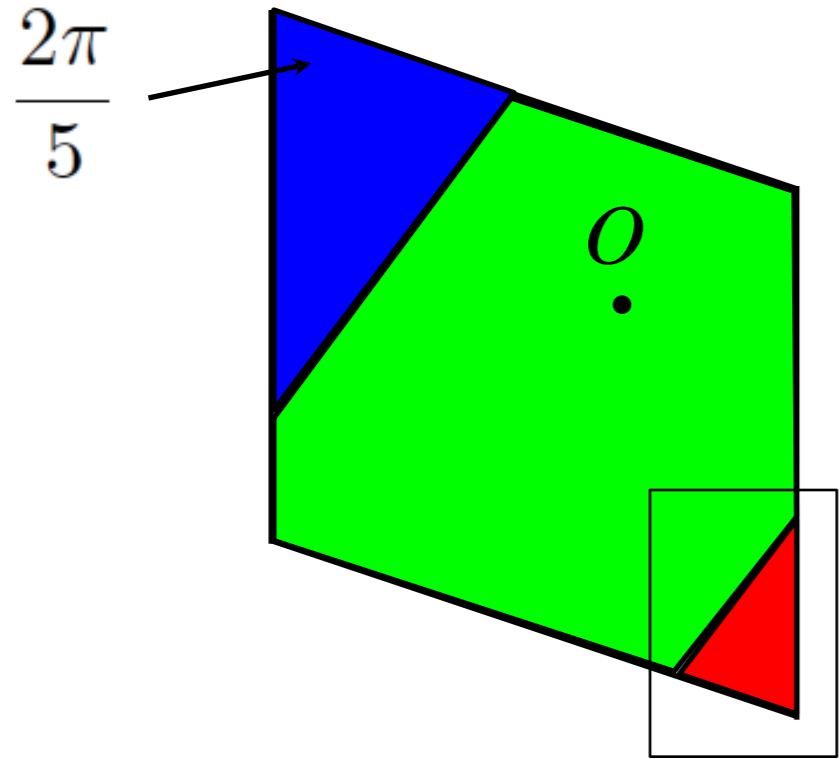


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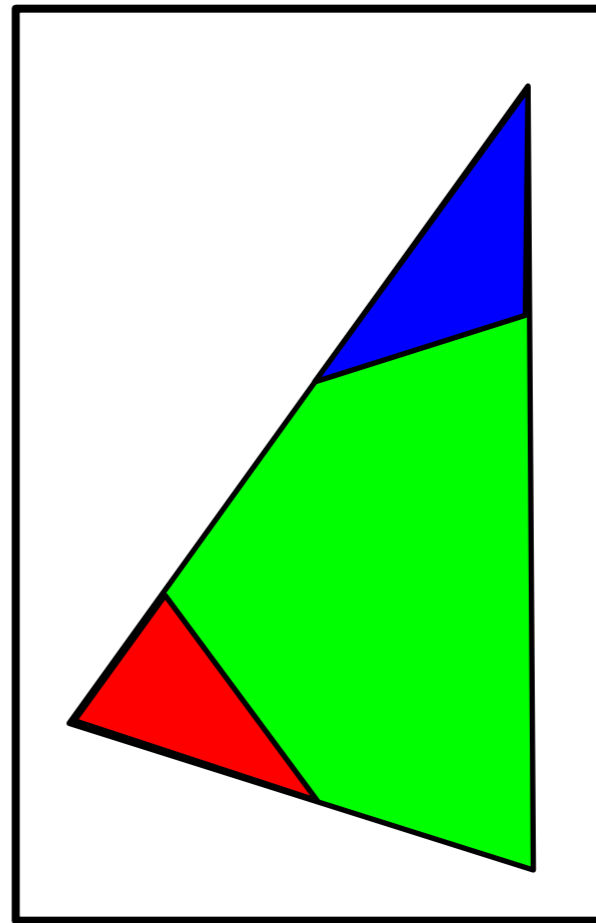
induction sequence

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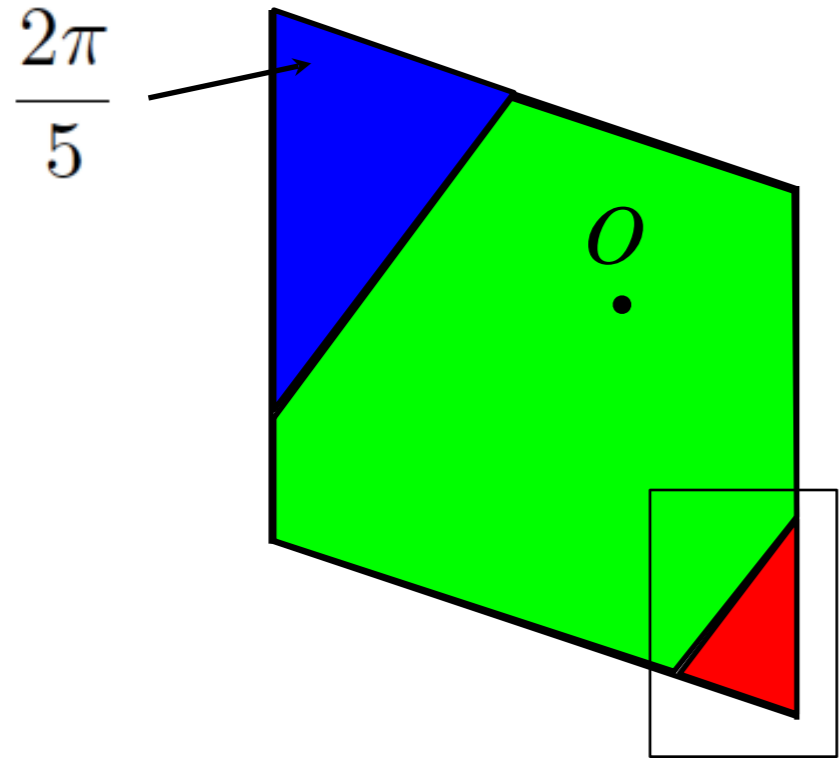


induction sequence

zoom in

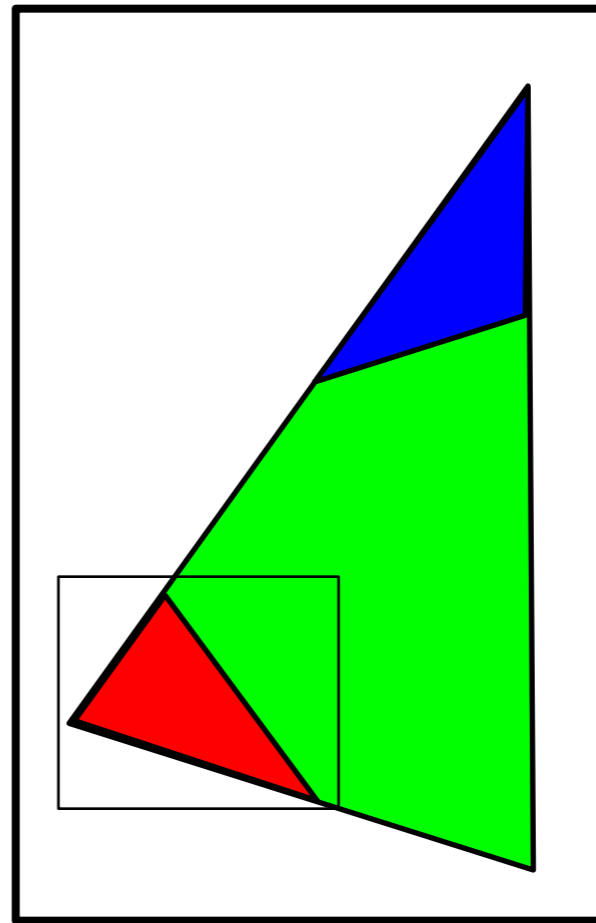


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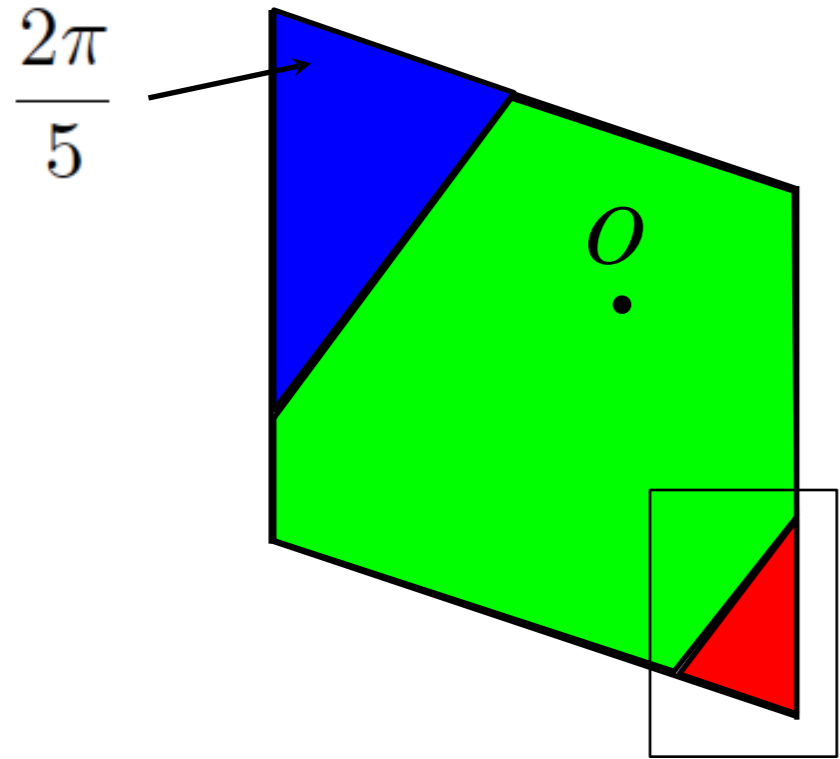


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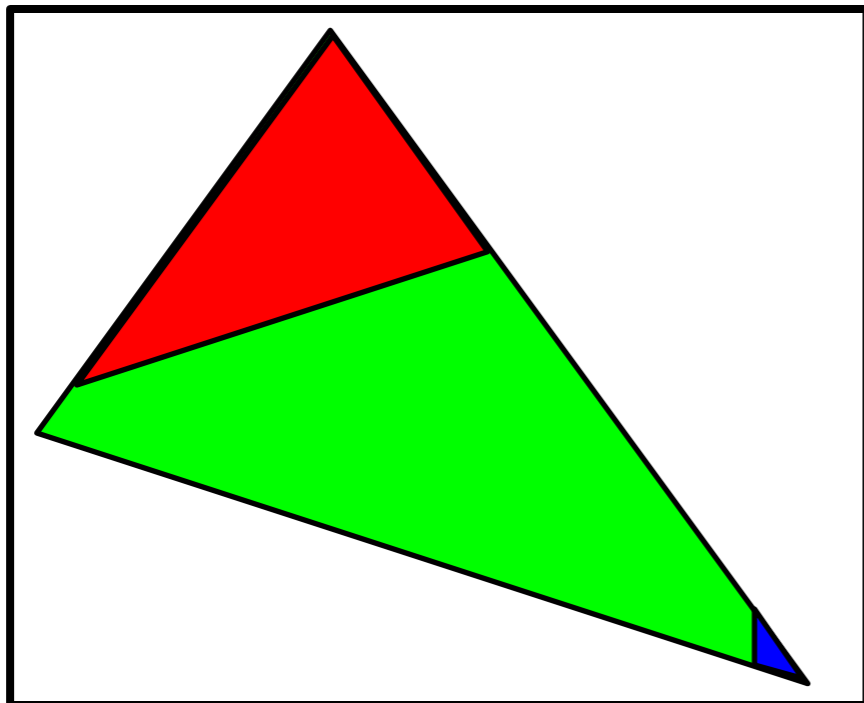
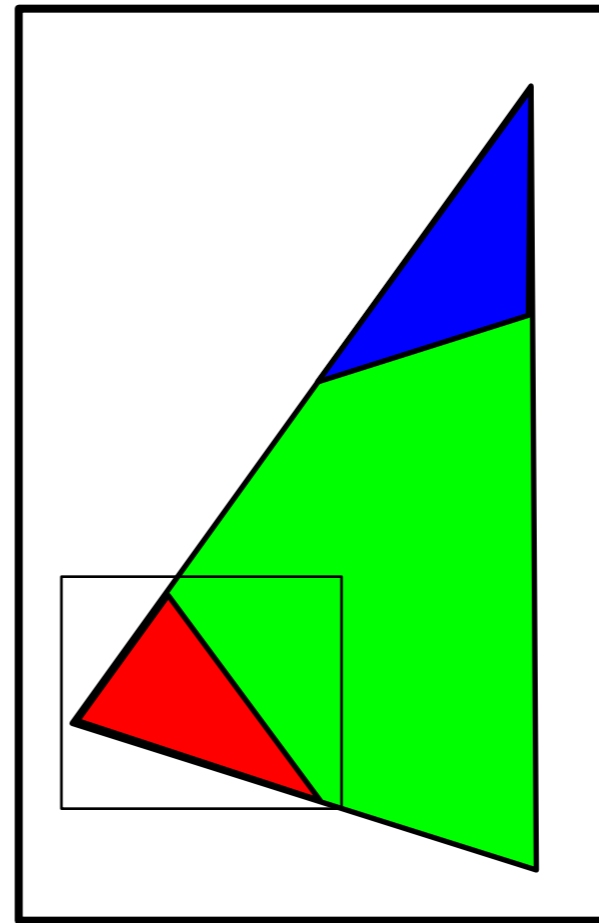


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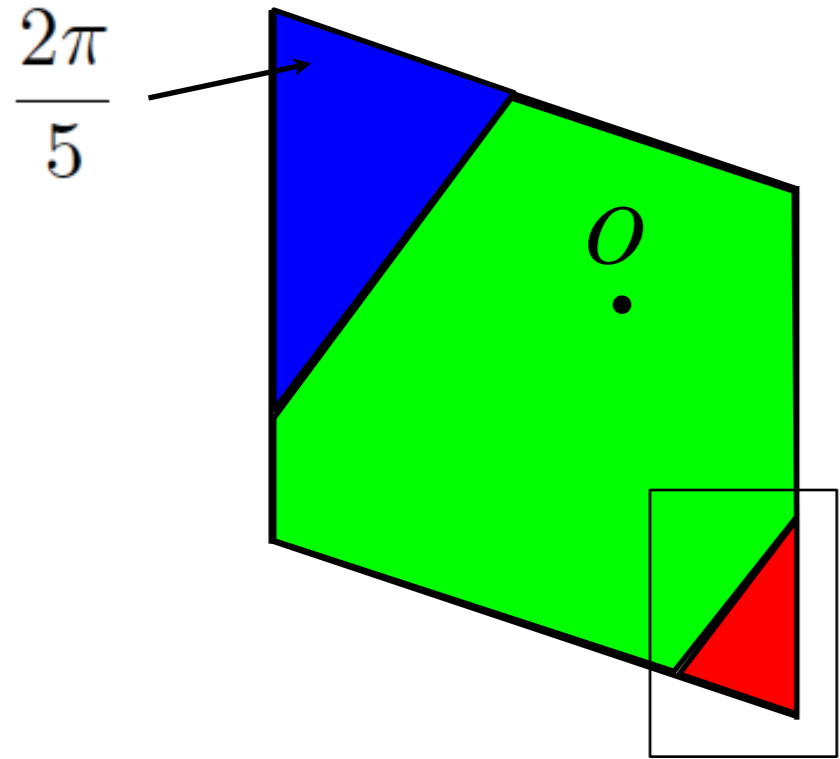


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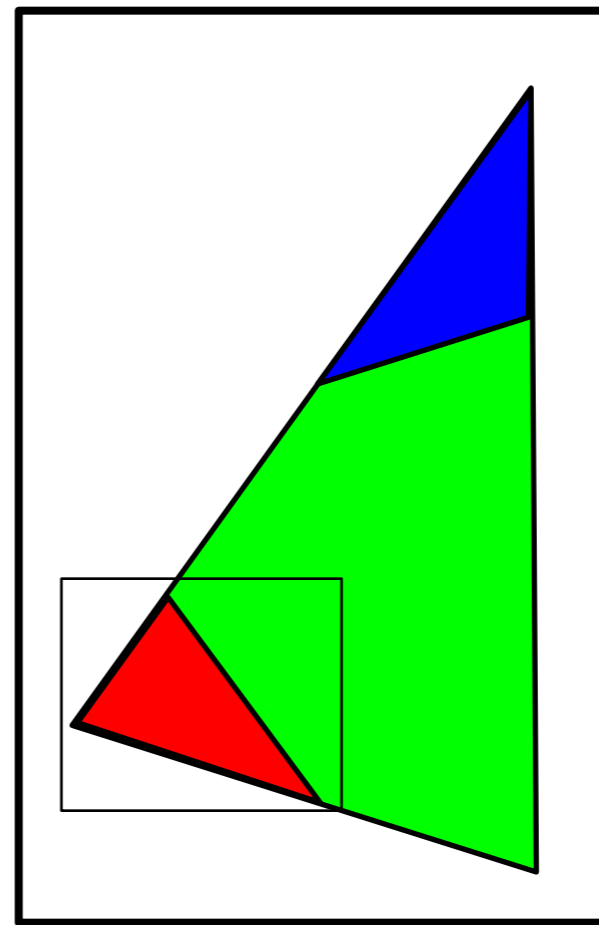


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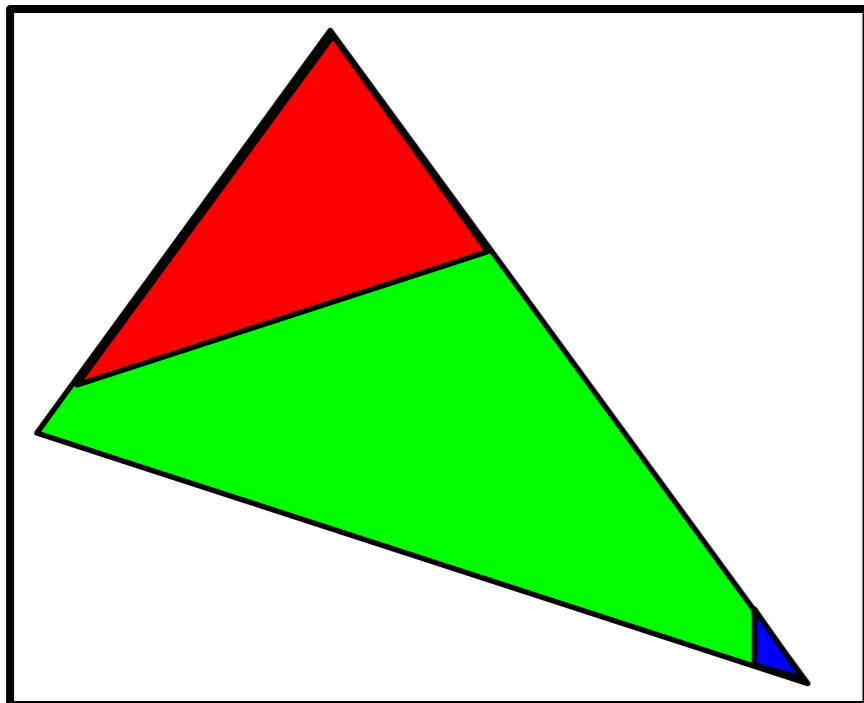


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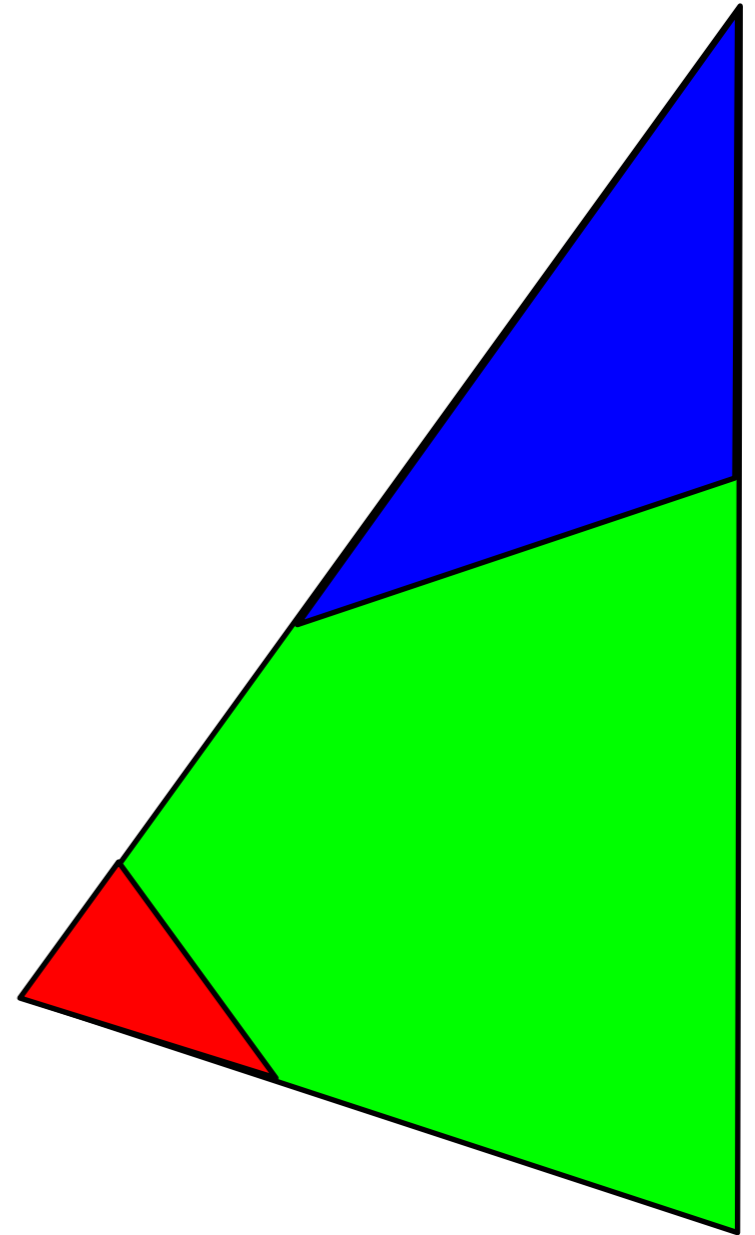
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base triangle

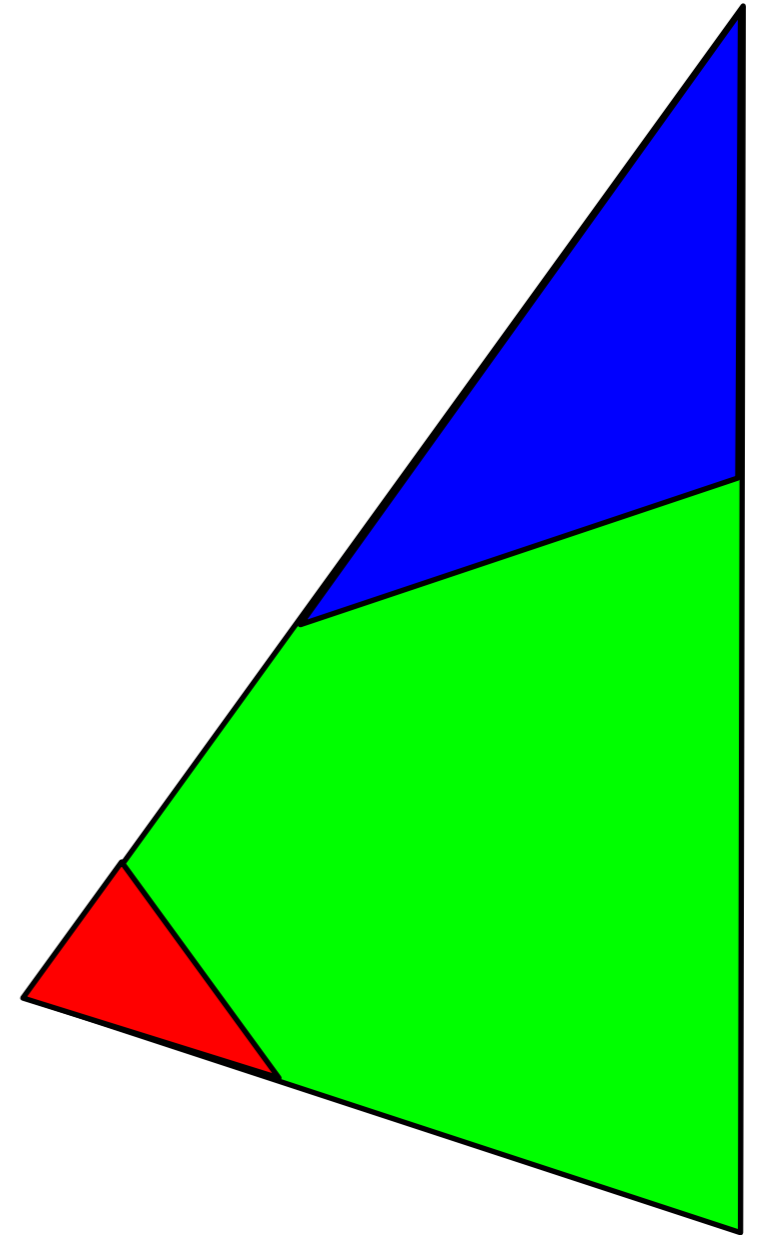


Base triangle: time-reversal symmetry



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The map is the composition of two involutions:  $F=GH$



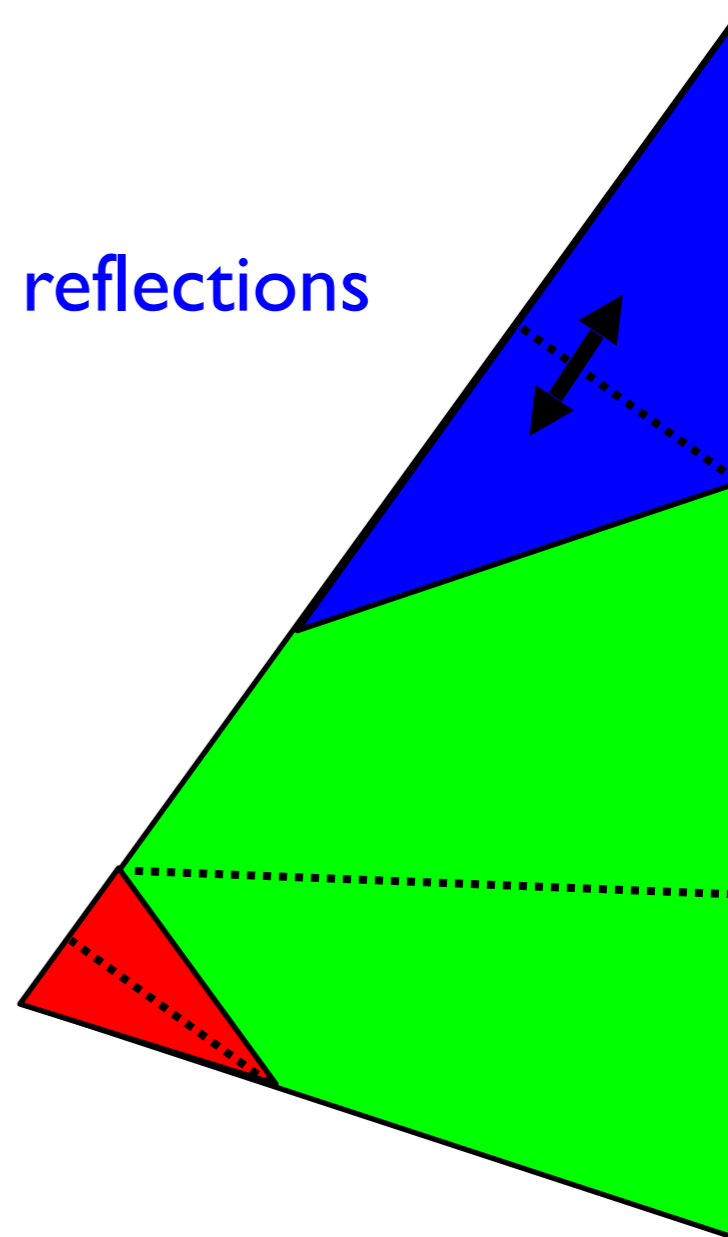


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local reflections

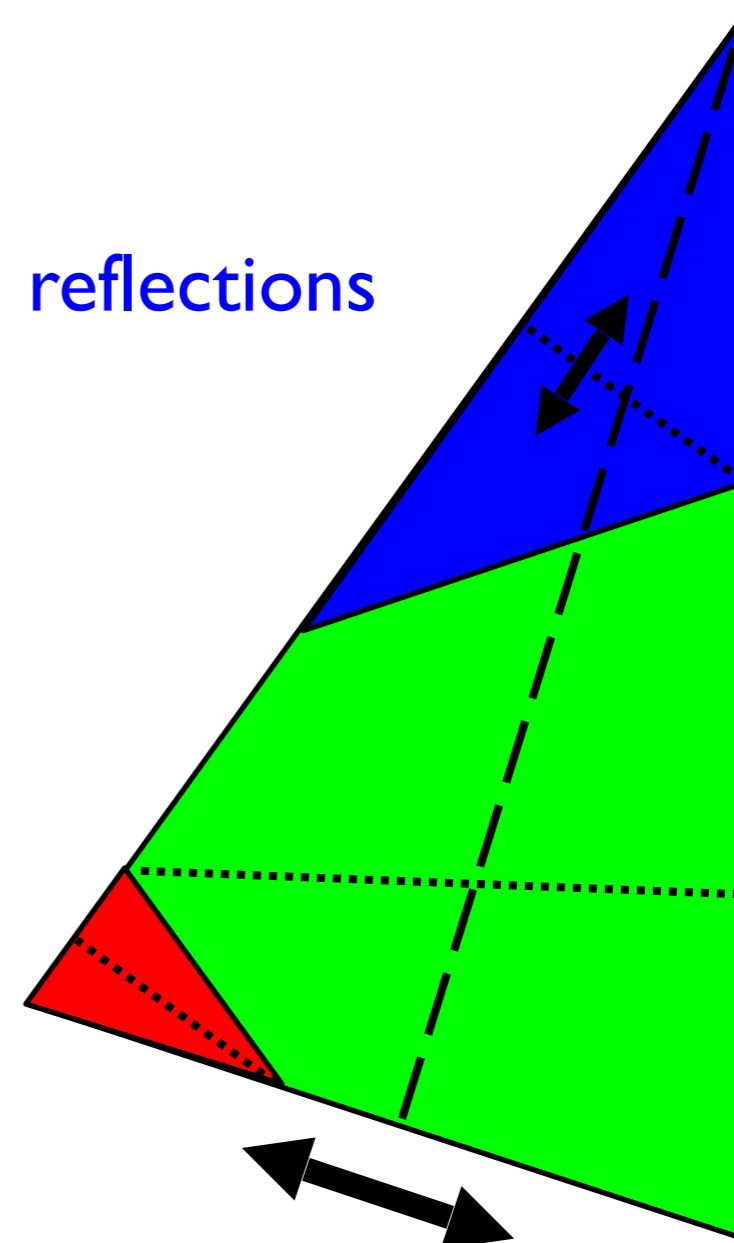


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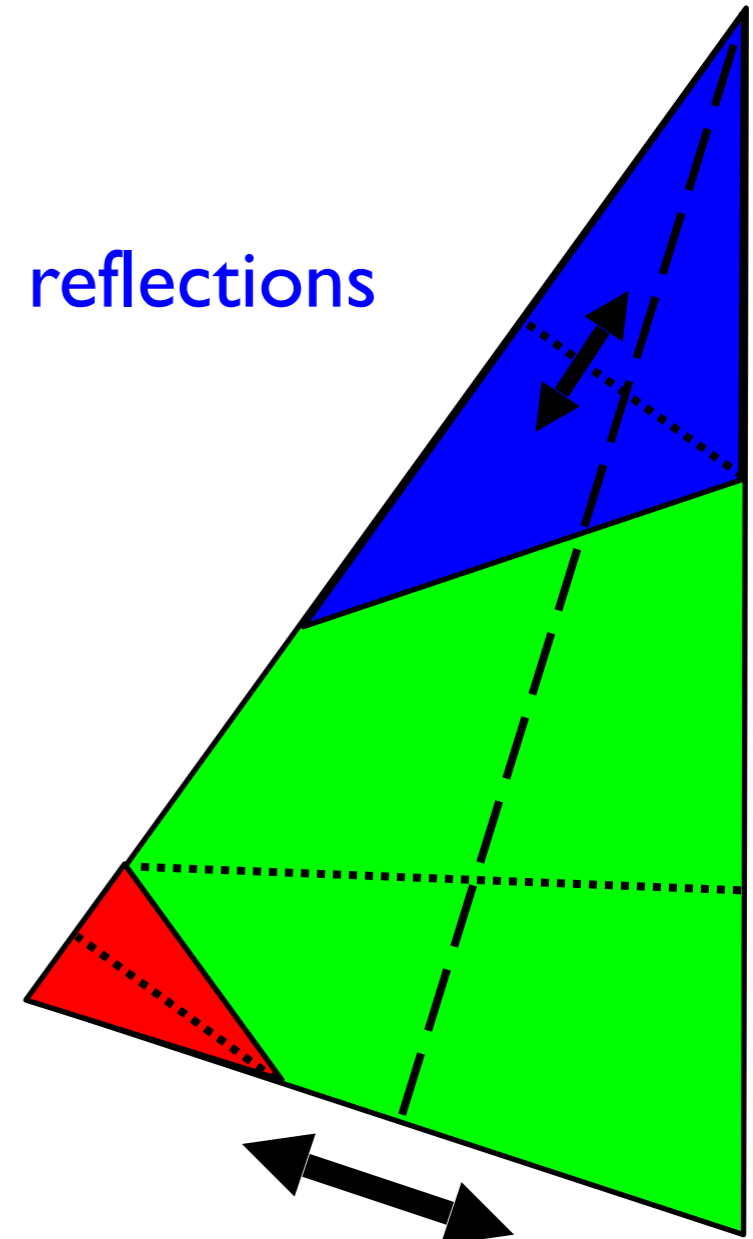
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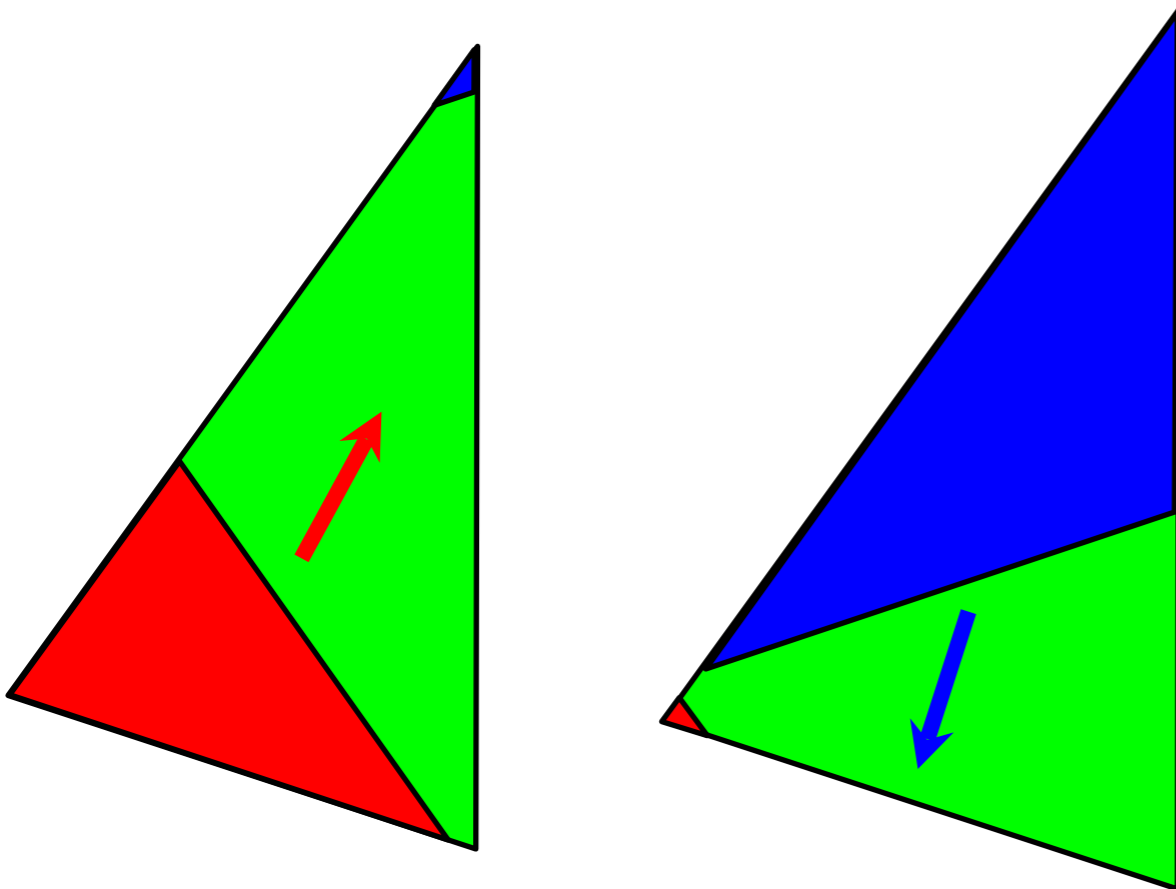
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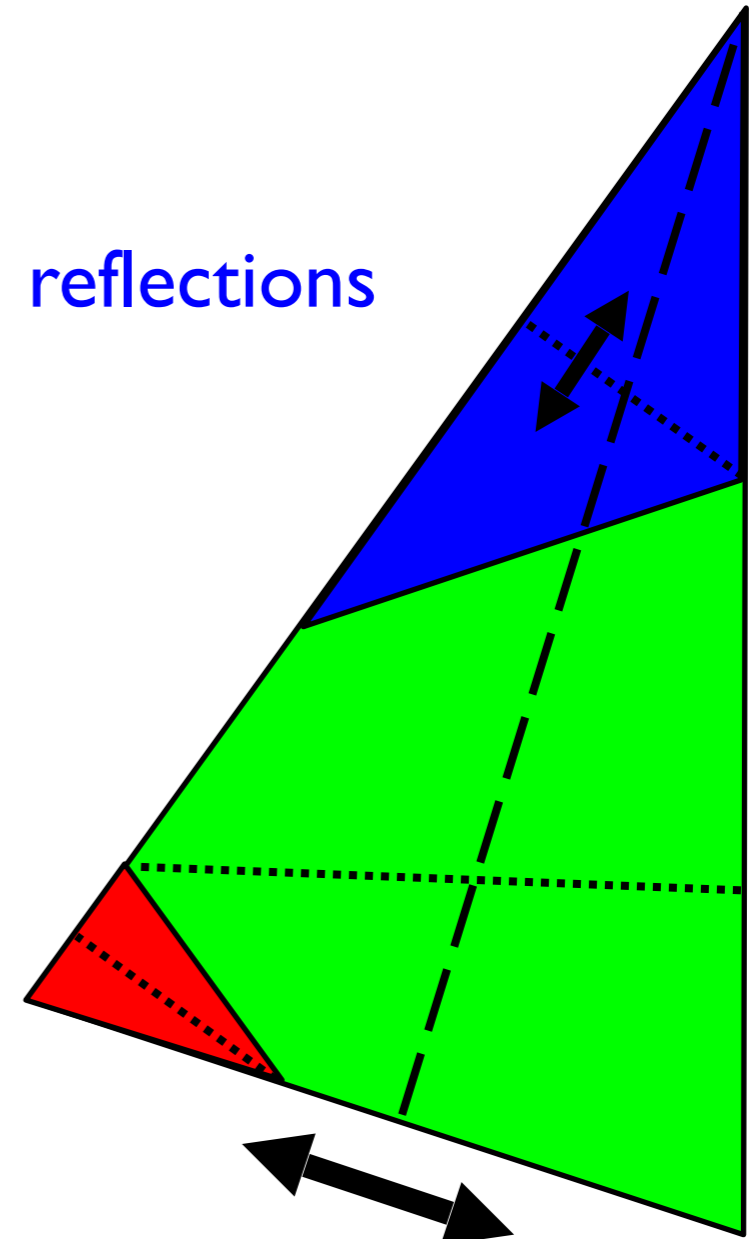
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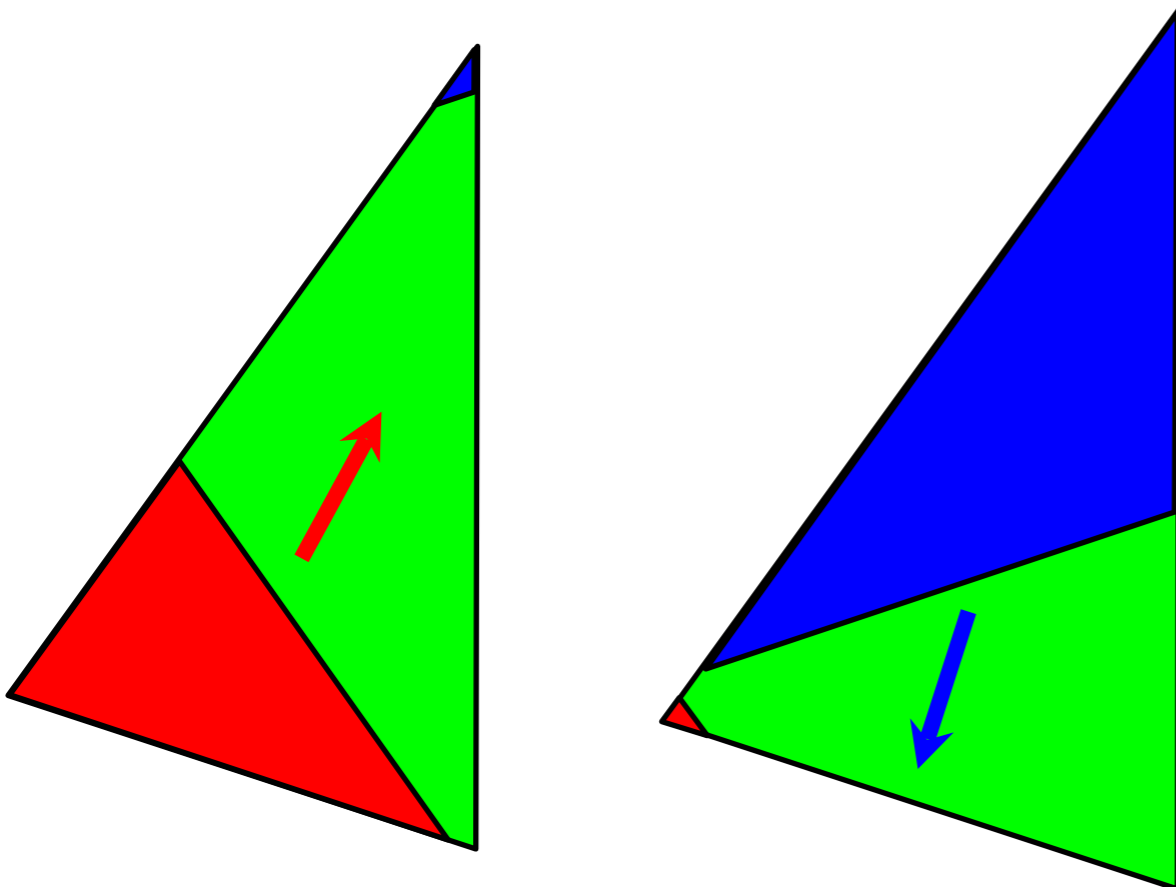
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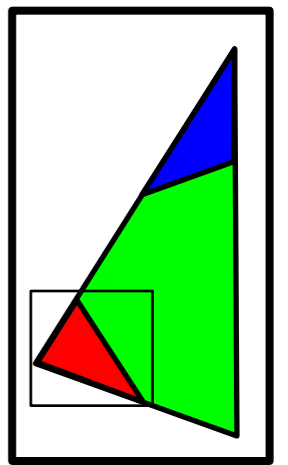
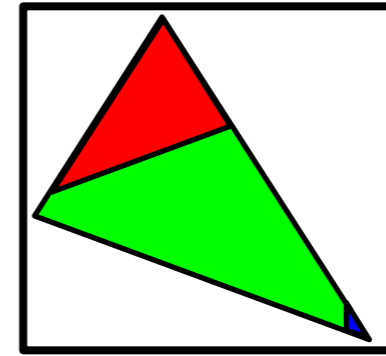


global reflection

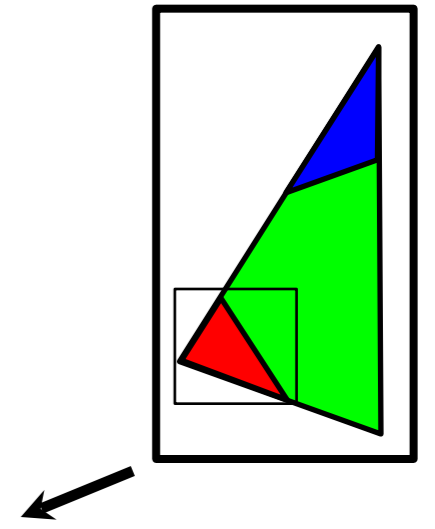
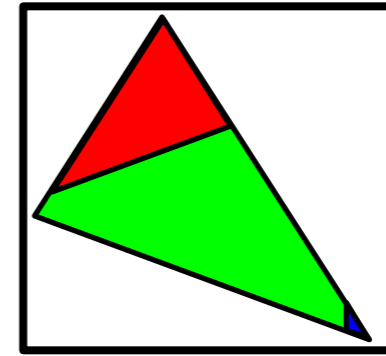
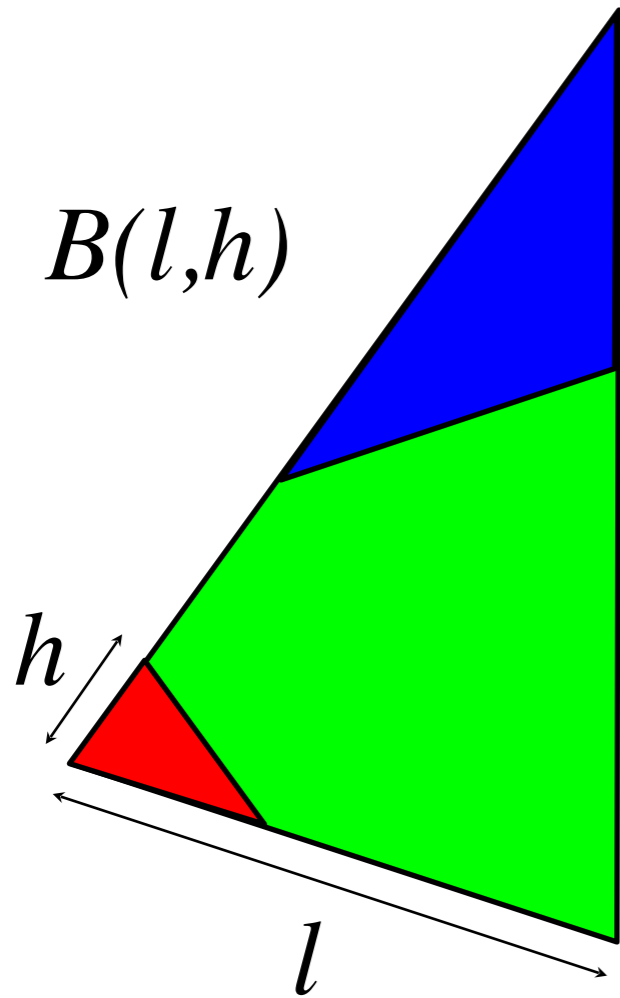
Bifurcations at end-points of parameter interval: one atom disappears.



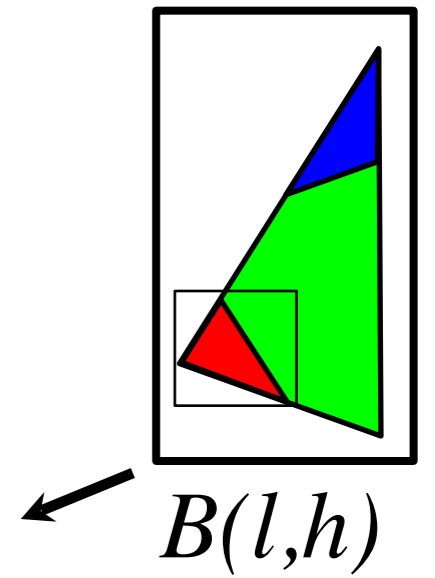
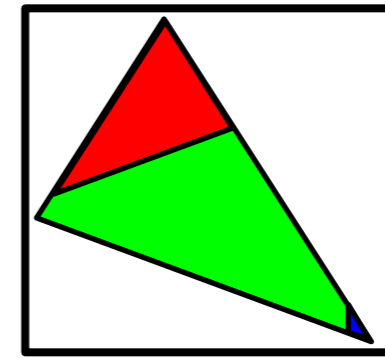
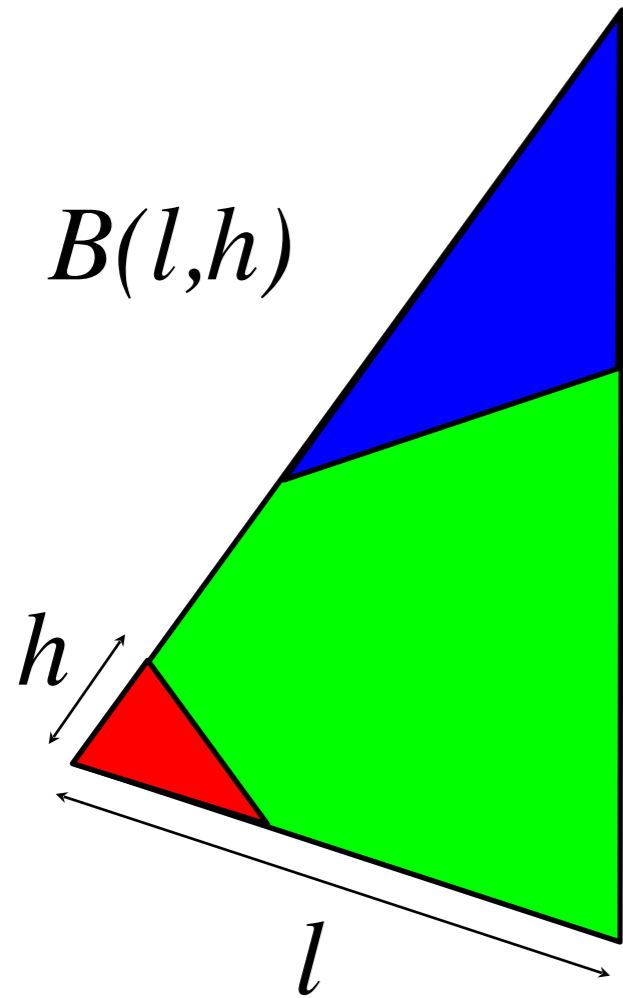
How does the parameter change under induction?



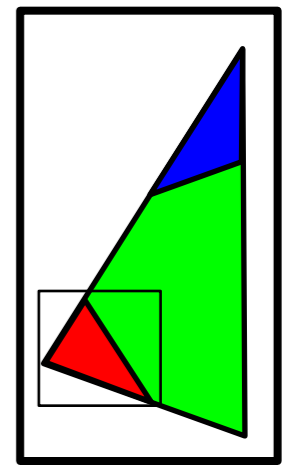
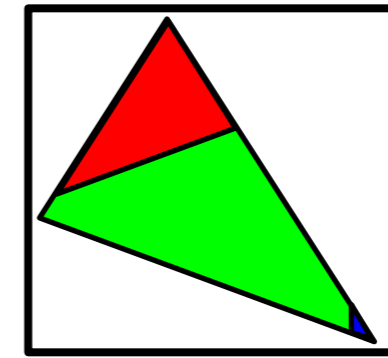
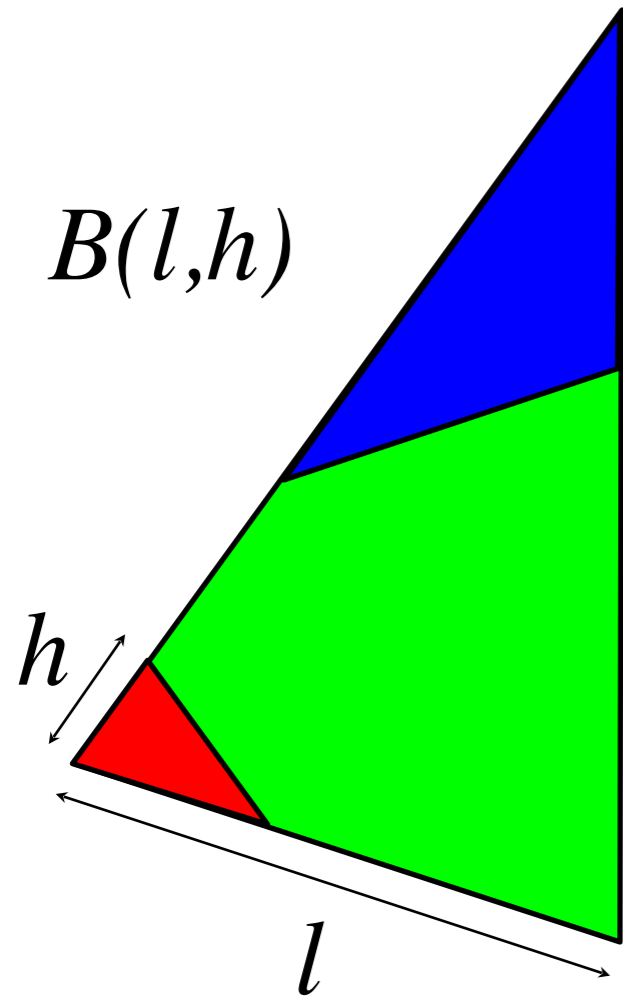
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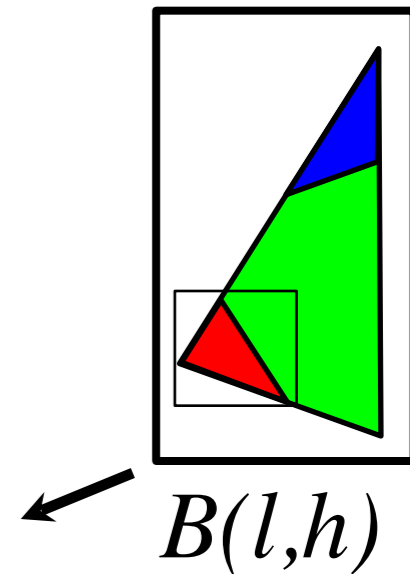
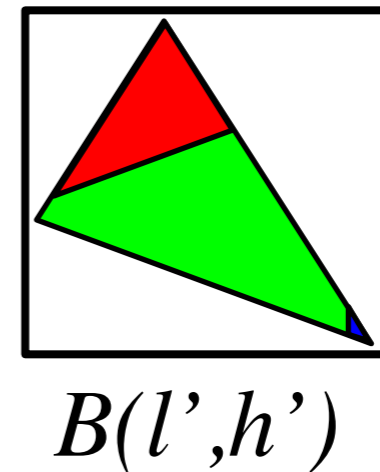
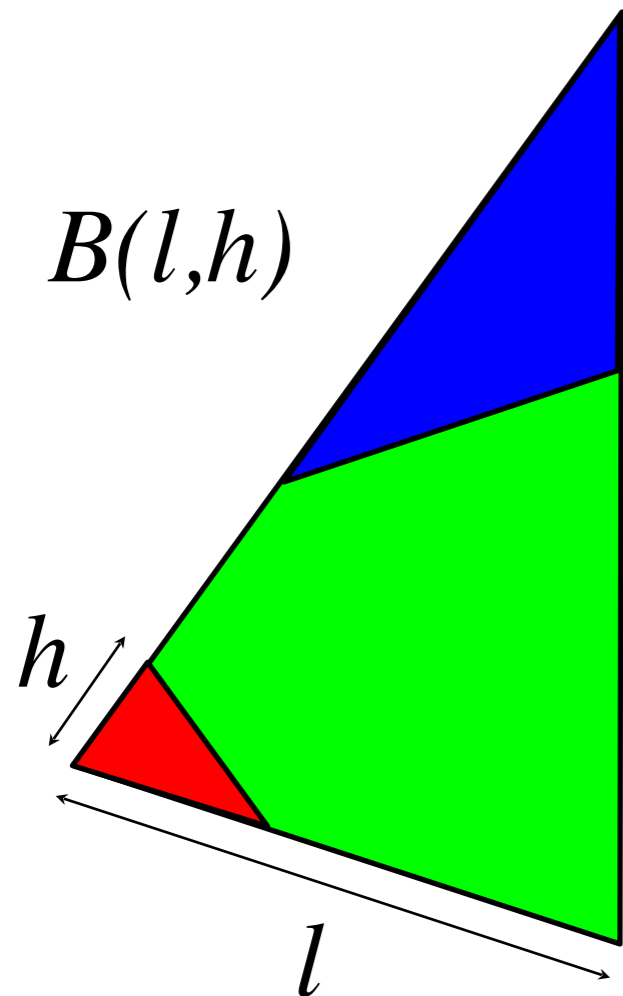


units

$$\omega = \frac{\sqrt{5} + 1}{2} \quad \beta = \frac{\sqrt{5} - 1}{2} = \omega^{-1}$$



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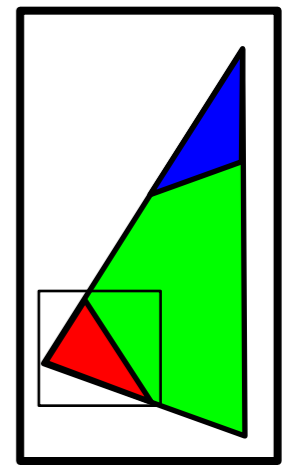
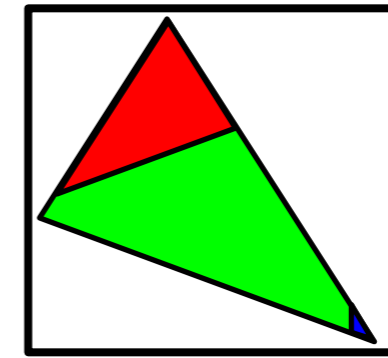
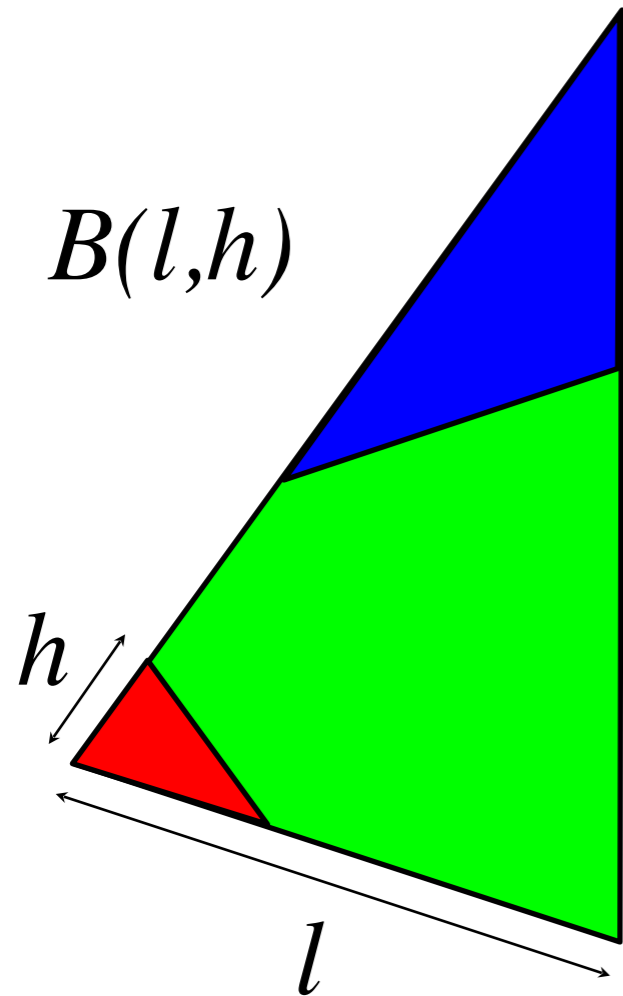
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new parameter:  $t=t(s)$

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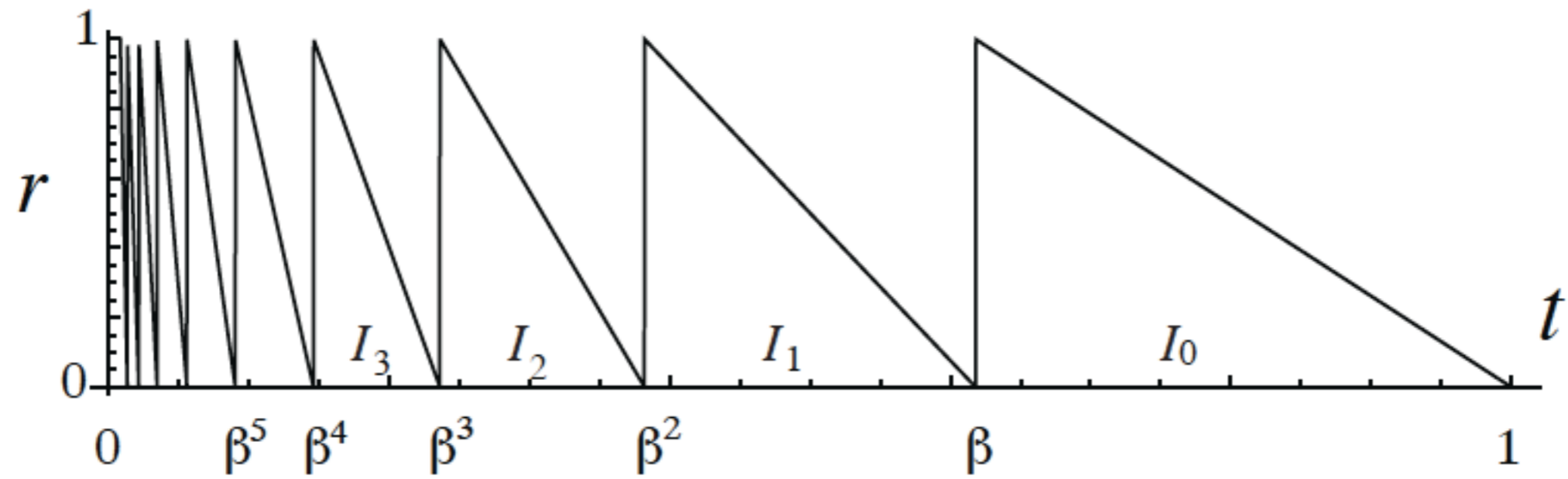
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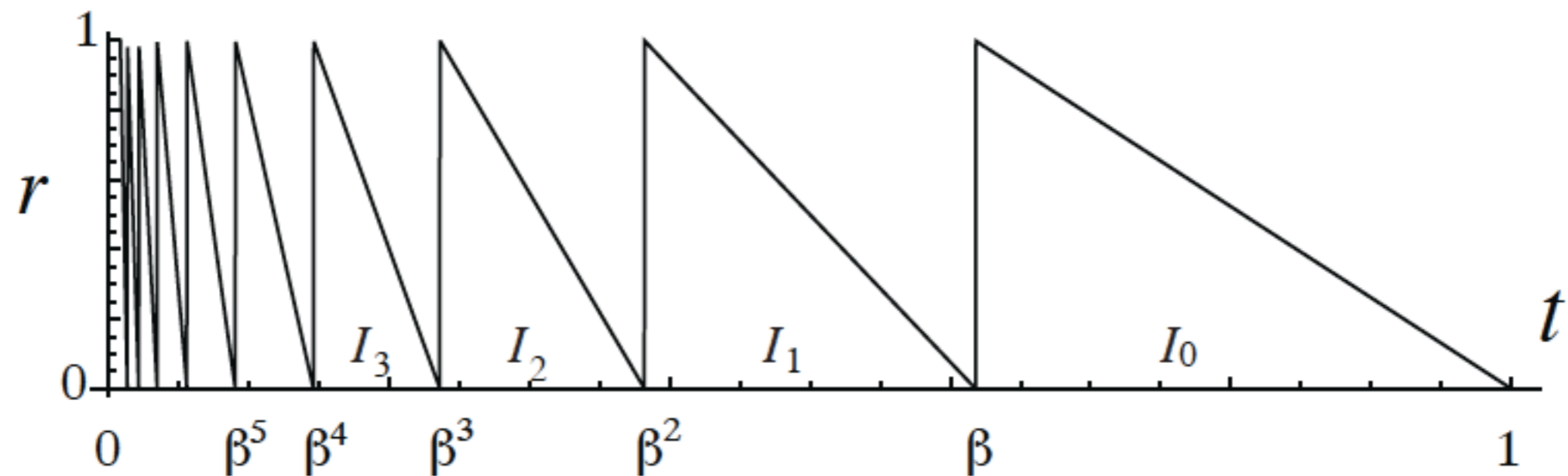
parameter renormalization function

$$t' = r(t)$$

# The parameter renormalization function

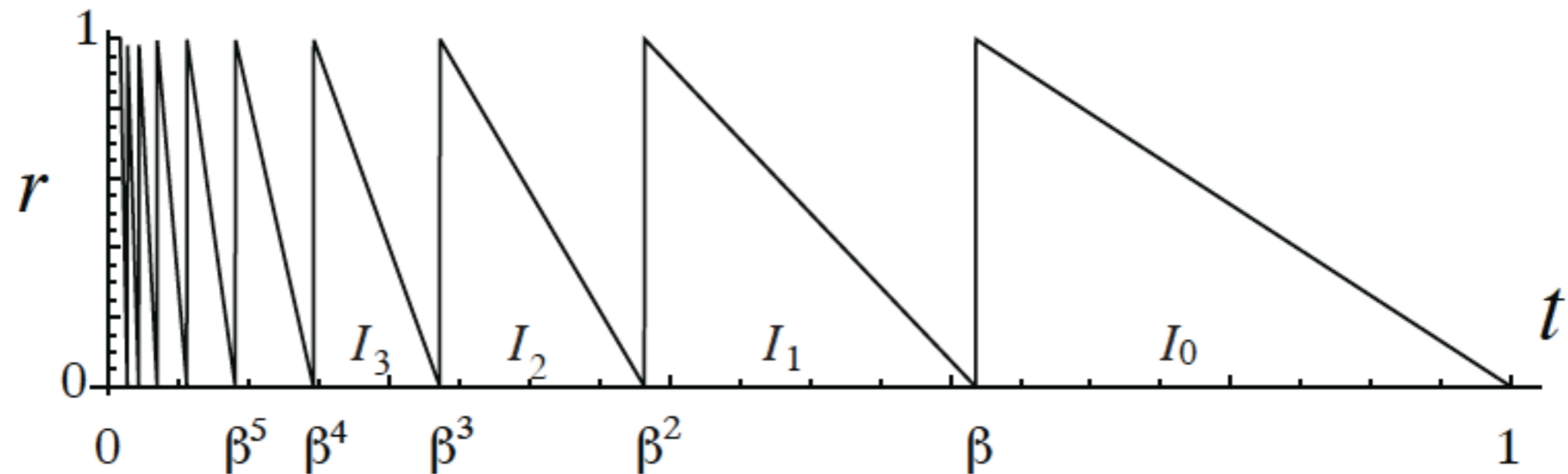


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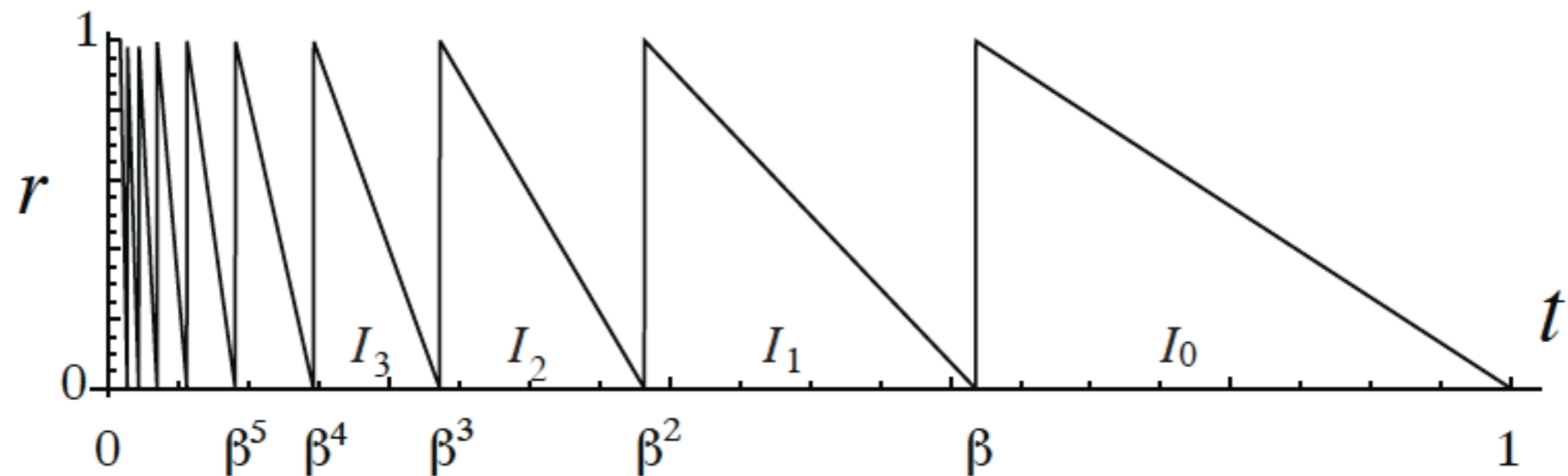
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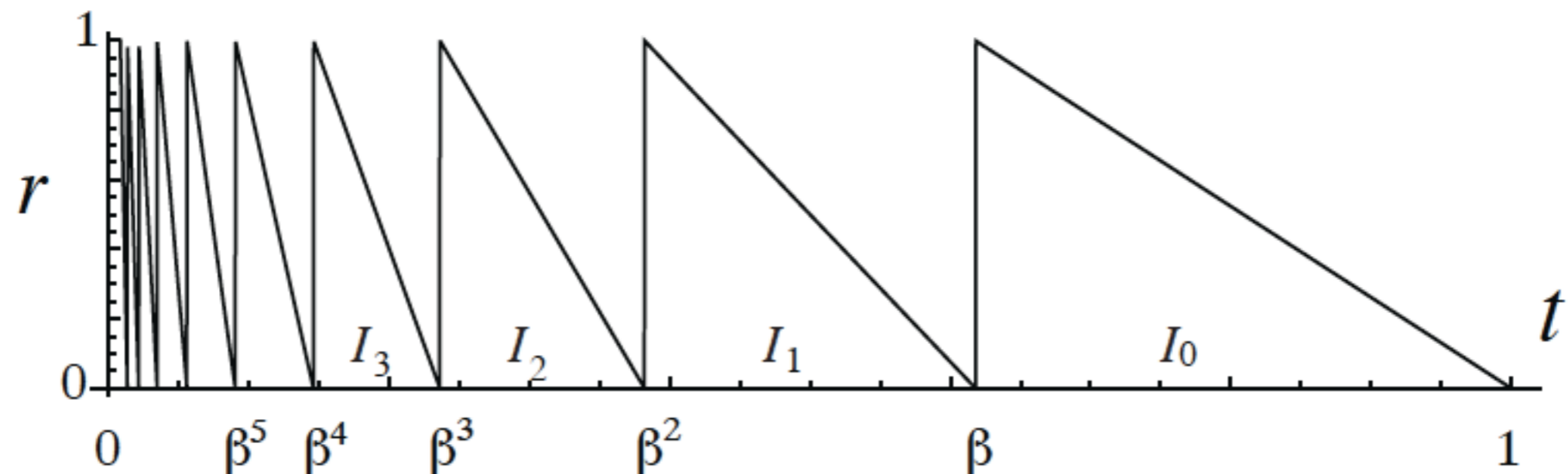
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# Generalised Luroth series



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There is a natural symbolic dynamics, recording the sub-intervals visited by an orbit:

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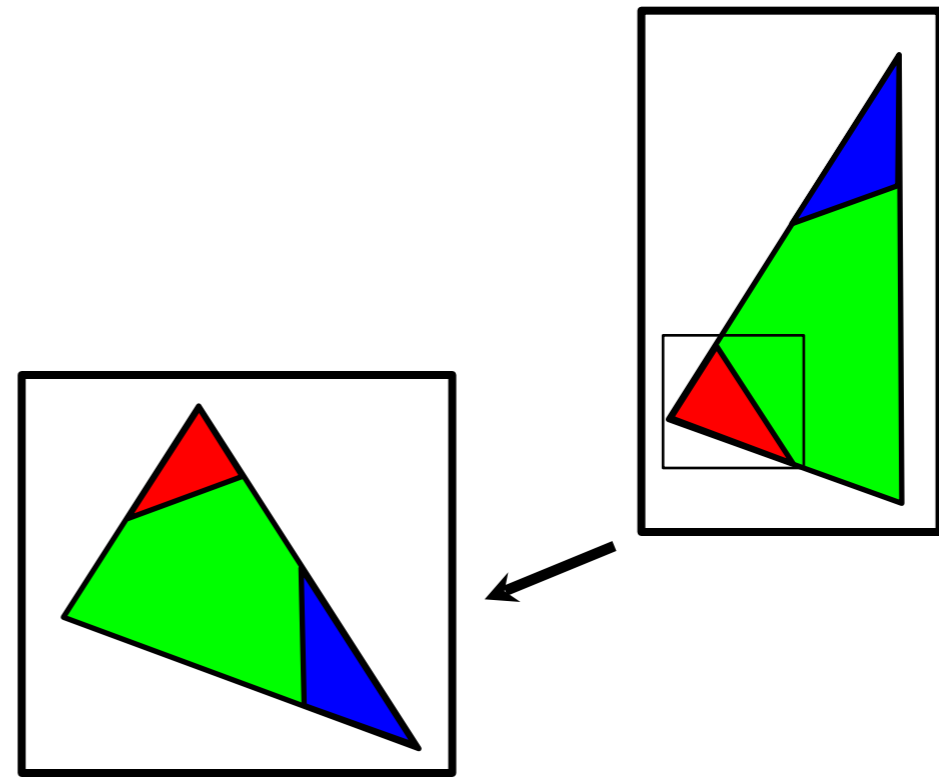
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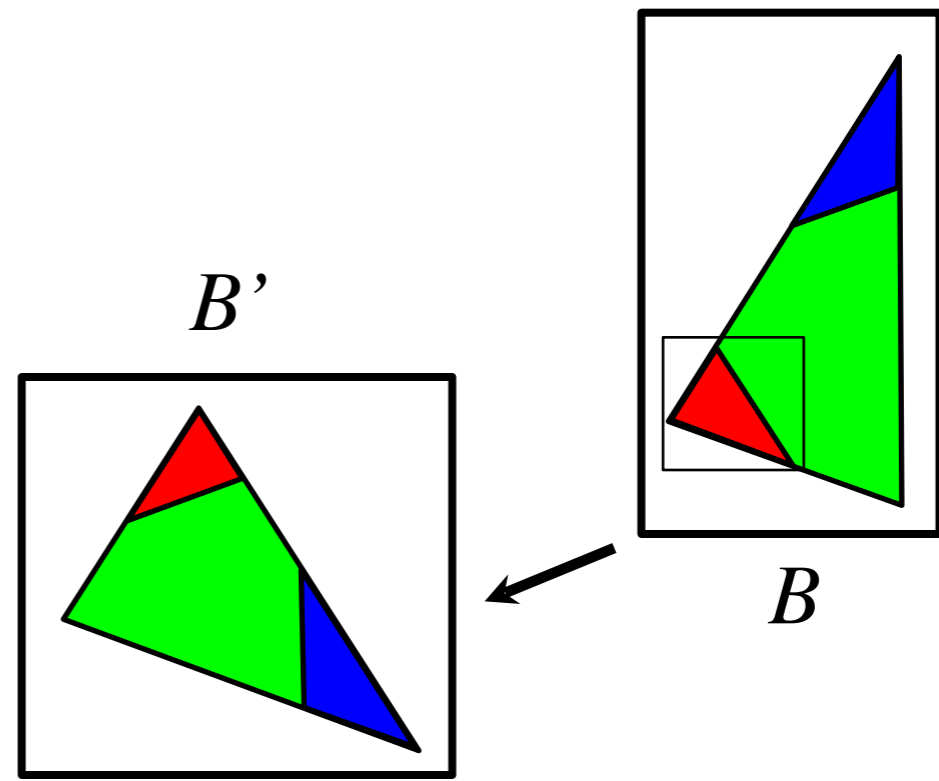
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A contraction argument shows that the field condition is sufficient for periodicity.

# Self-similarity

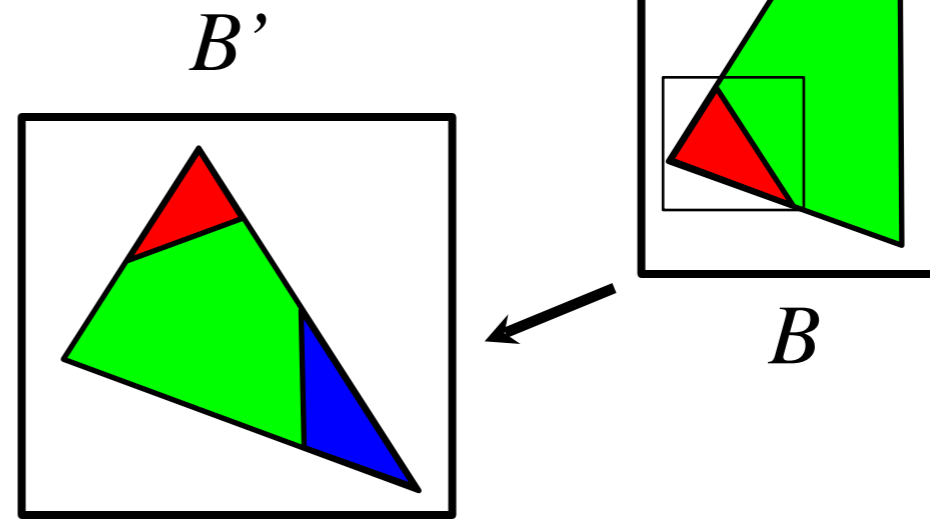


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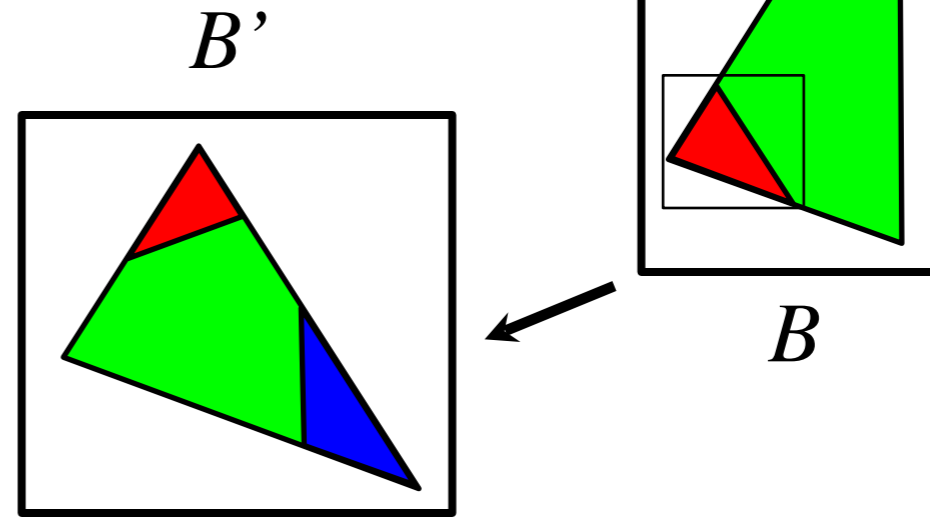
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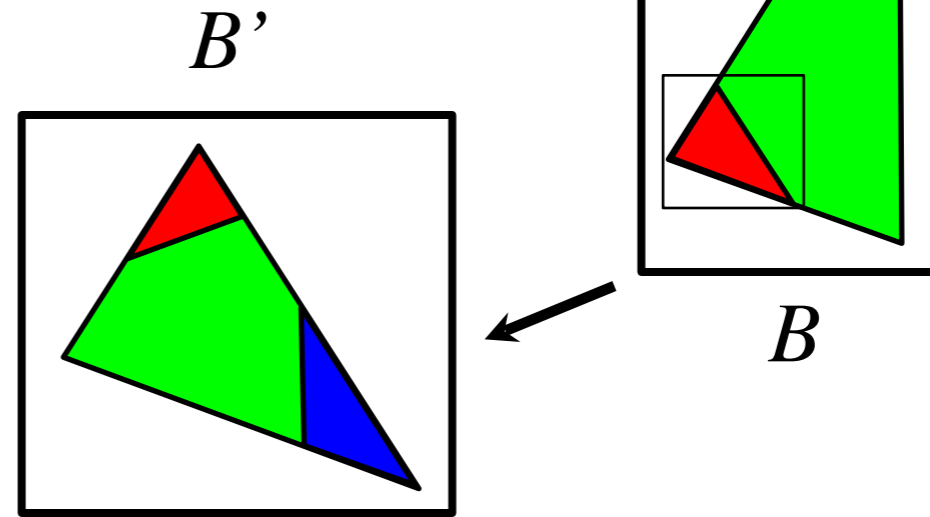


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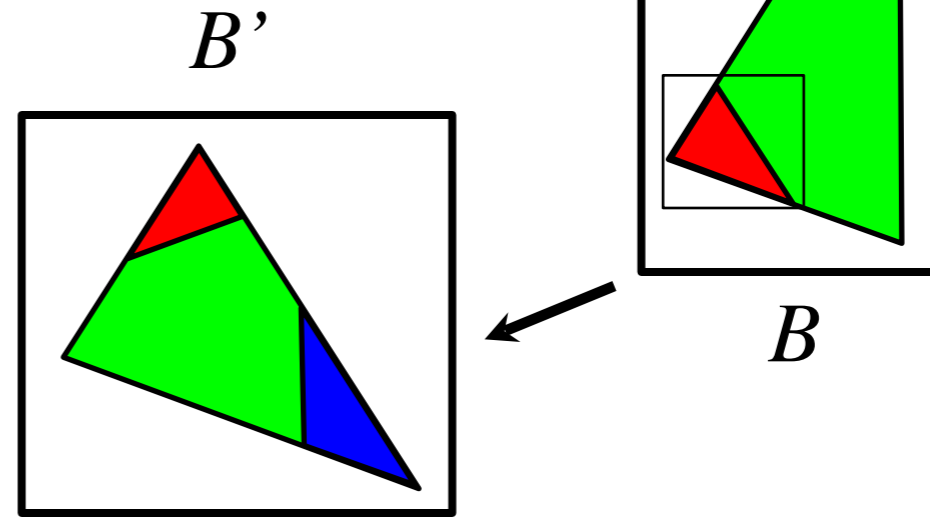


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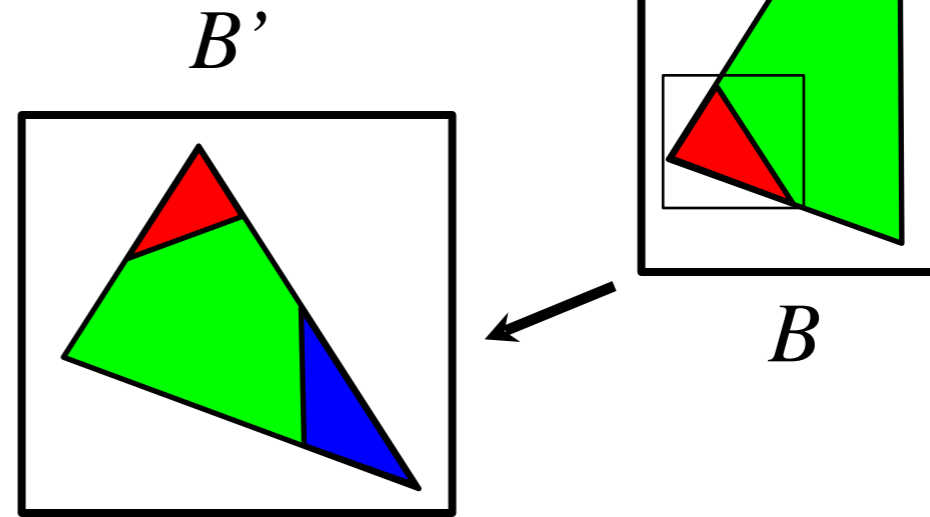


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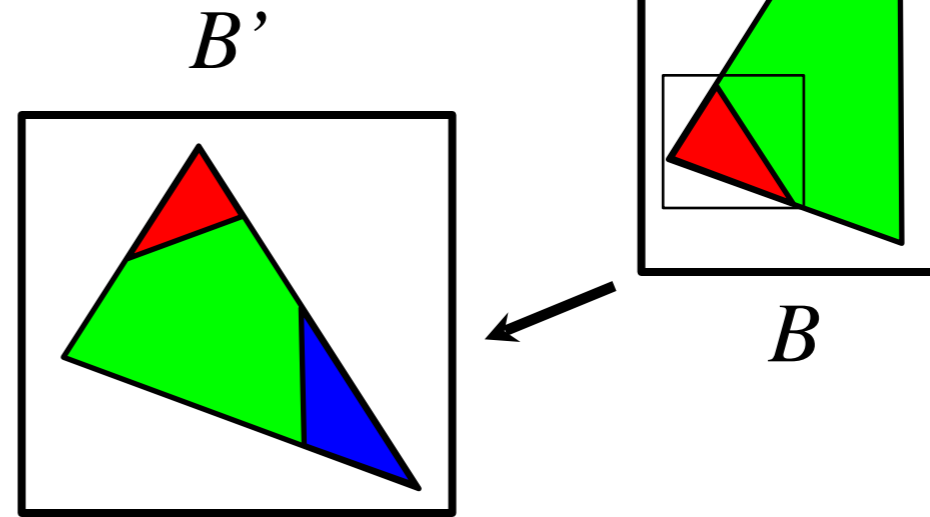


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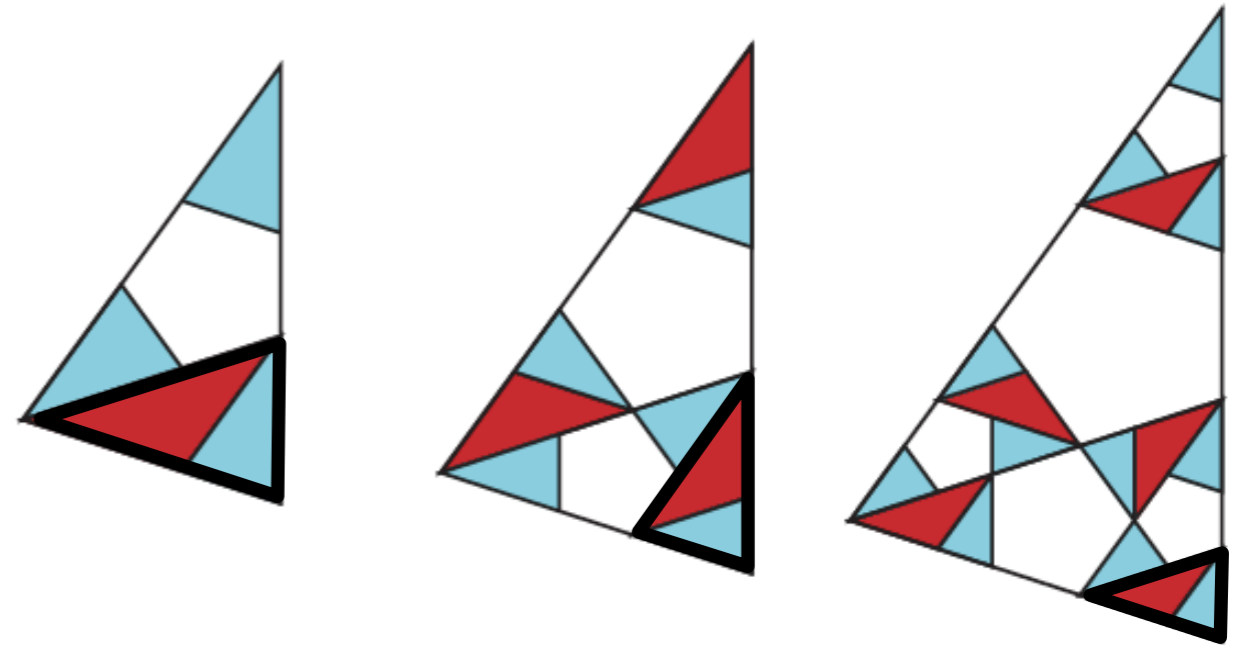
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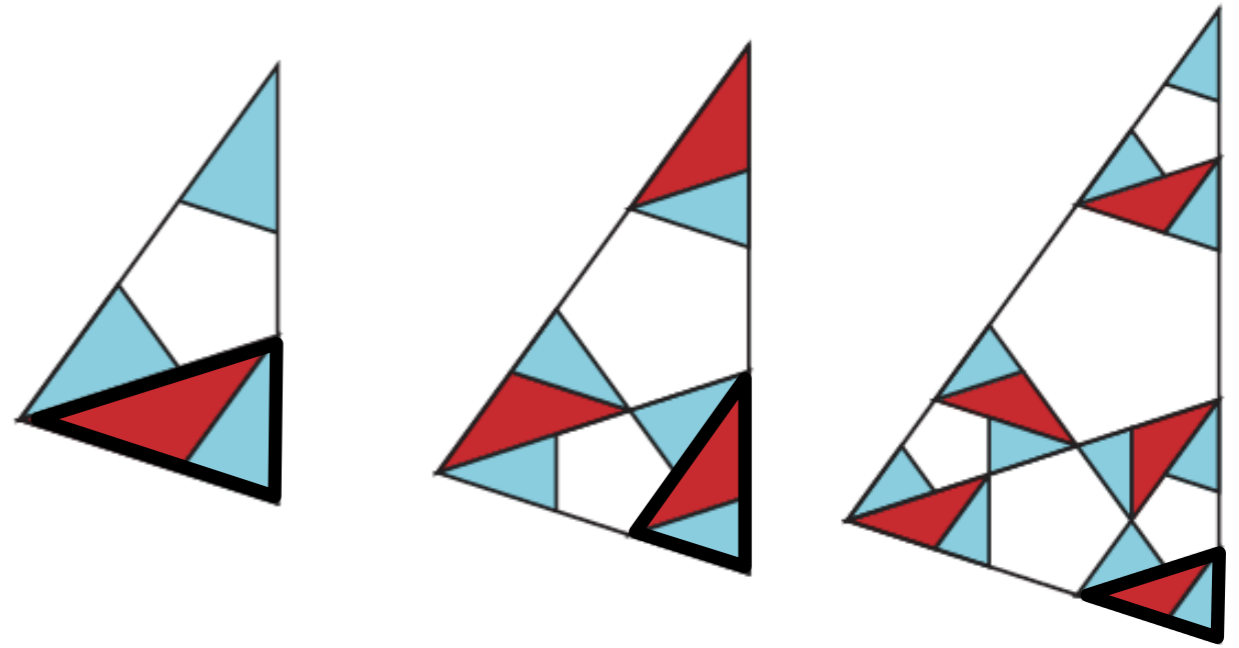
- We write  $B \sim B'$  to denote congruence with respect to the following transformations:
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- This equivalence extends naturally to the PWIs on  $B$  and  $B'$  (written  $\mathbf{B}, \mathbf{B}'$ ), by matching the corresponding atoms and their images:  $\mathbf{B} \sim \mathbf{B}'$ .

# Tiling and renormalizability



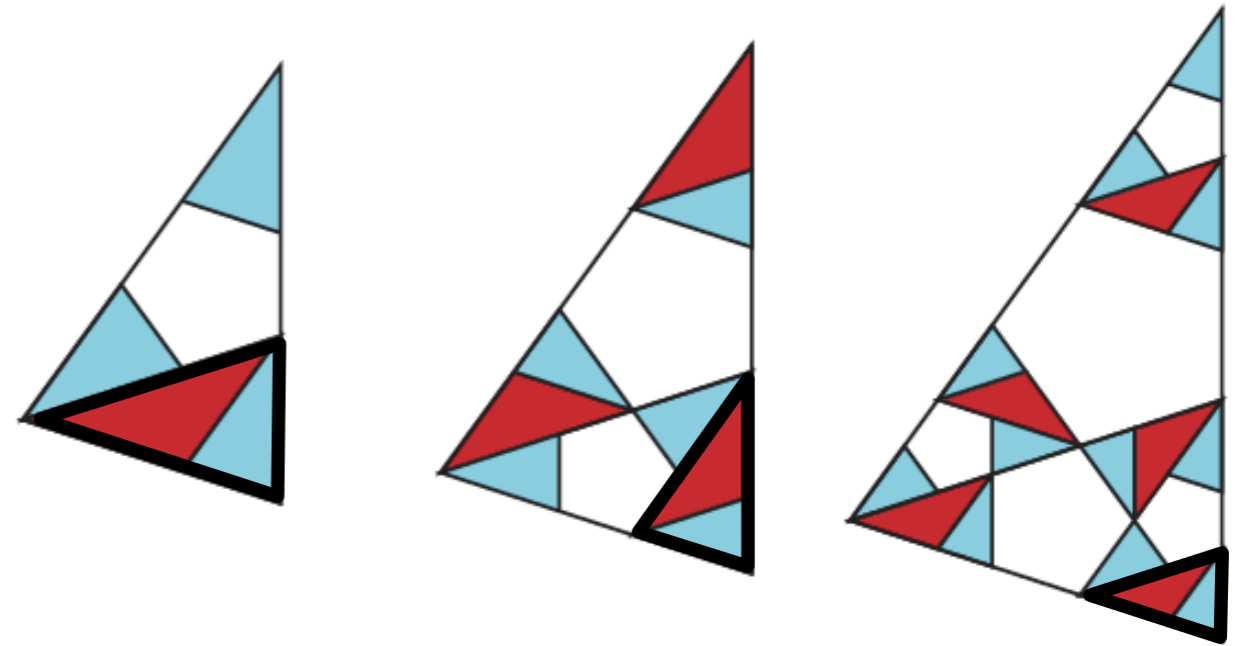
# Tiling and renormalizability

Refined coverings of the exceptional set, via recursive tiling.



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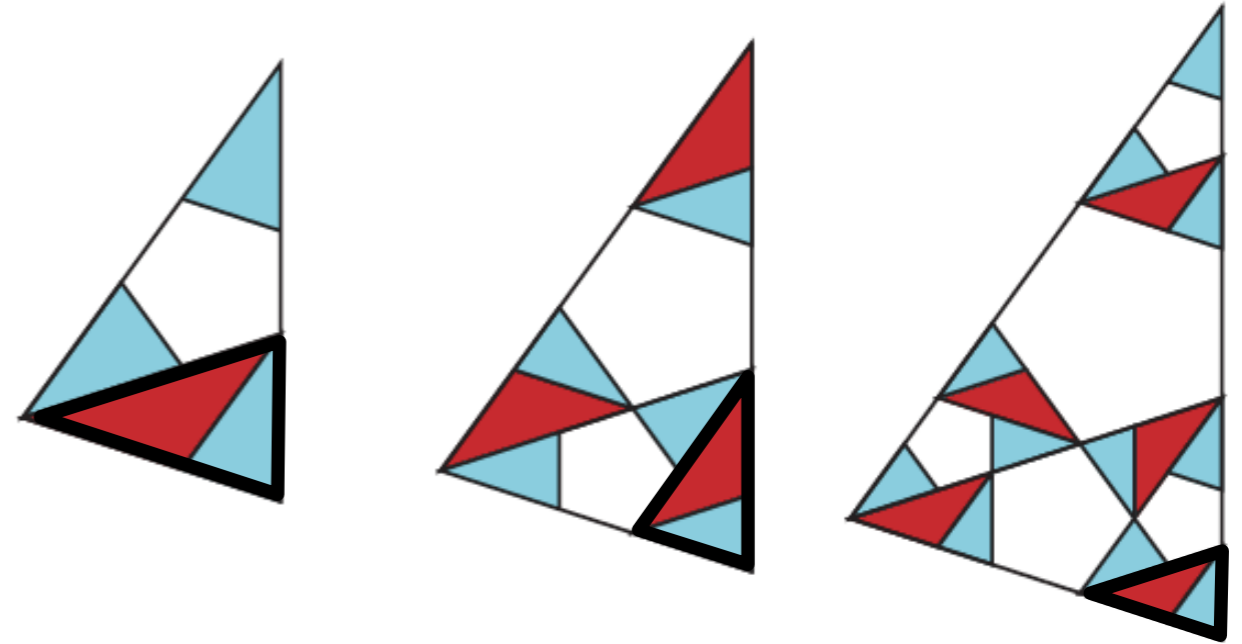
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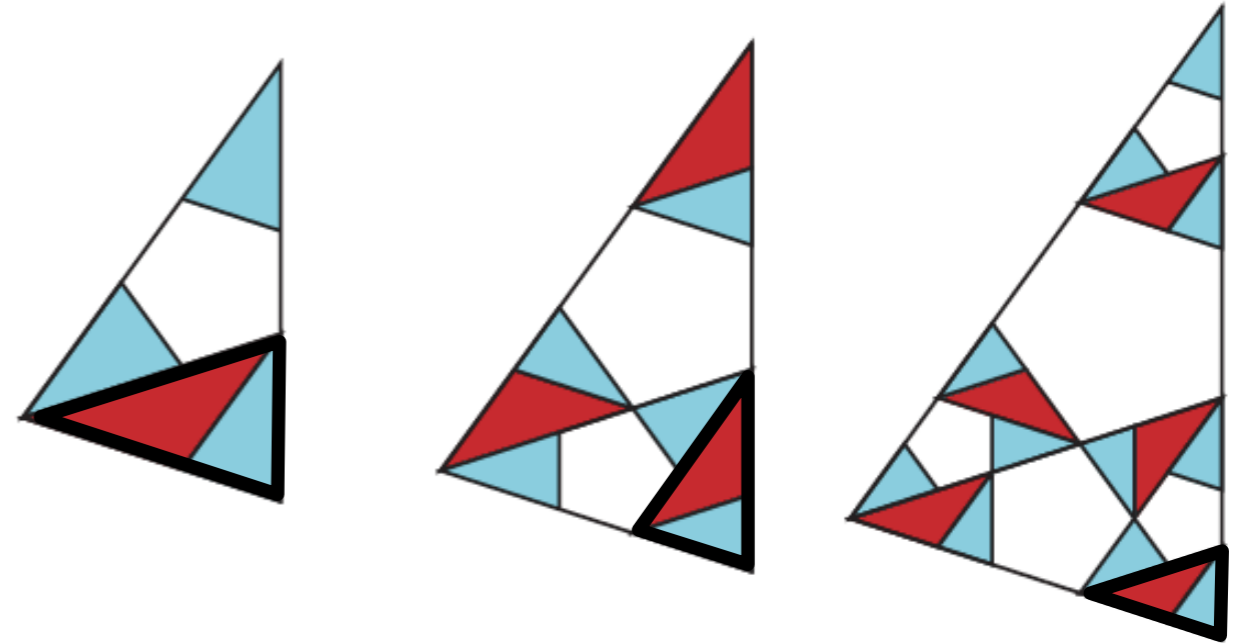


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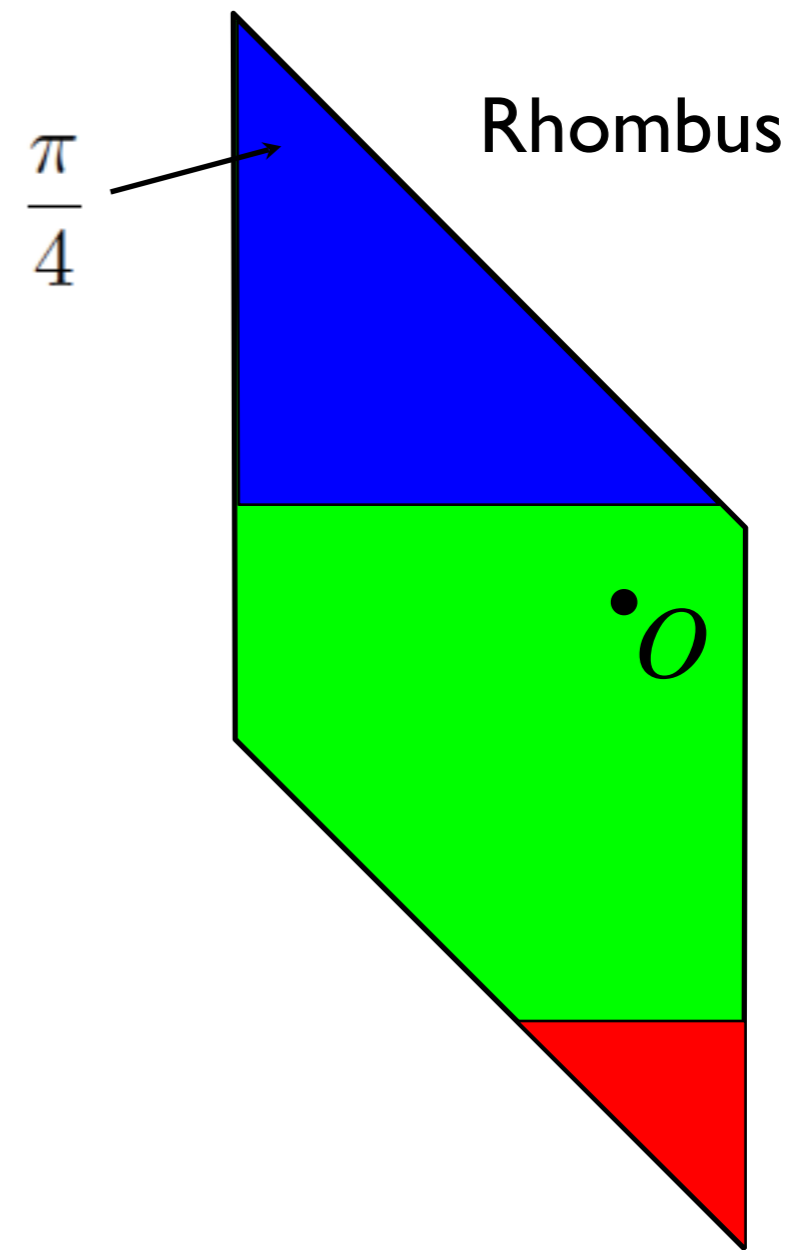


- We say that  $\mathbf{B}'$  tiles  $\mathbf{B}$  if the first return orbit of the atoms of  $\mathbf{B}'$  covers  $B$  apart from a finite number of periodic tiles.
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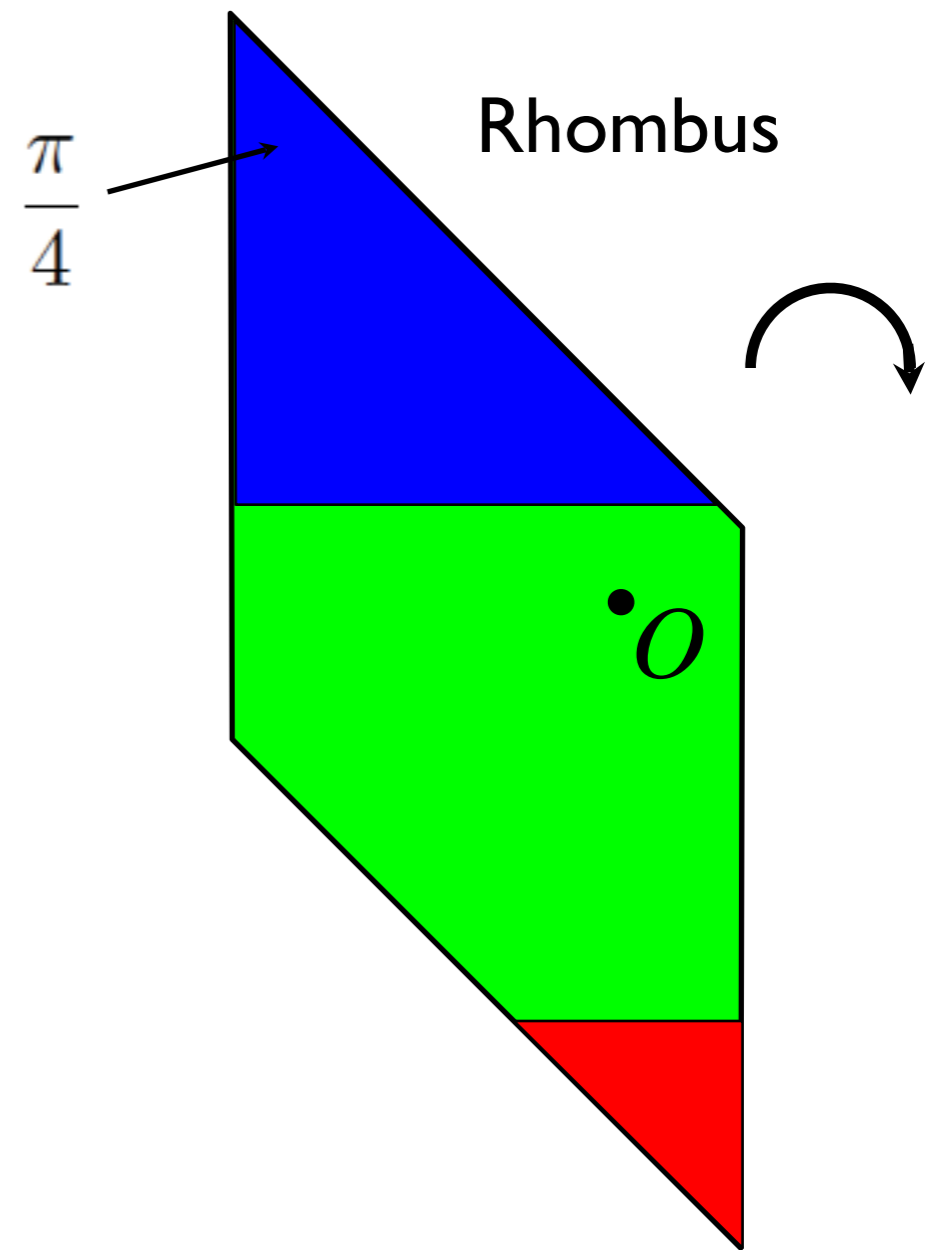
**Theorem.** The one-parameter pentagonal model is renormalizable if and only if the parameter belongs to the field  $\mathbb{Q}(\sqrt{5})$ .

The octagonal model (field  $\mathbb{Q}(\sqrt{2})$ )

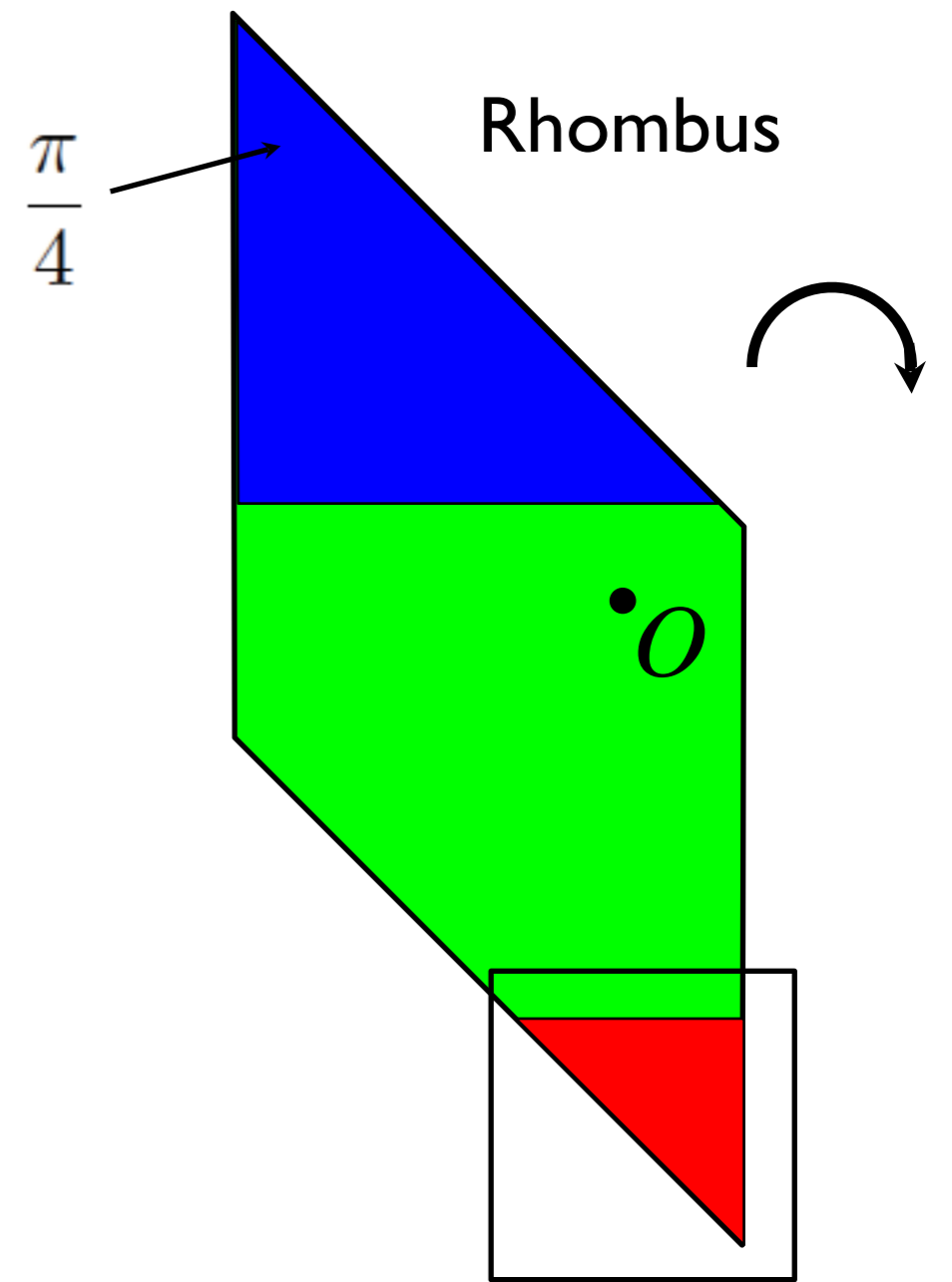
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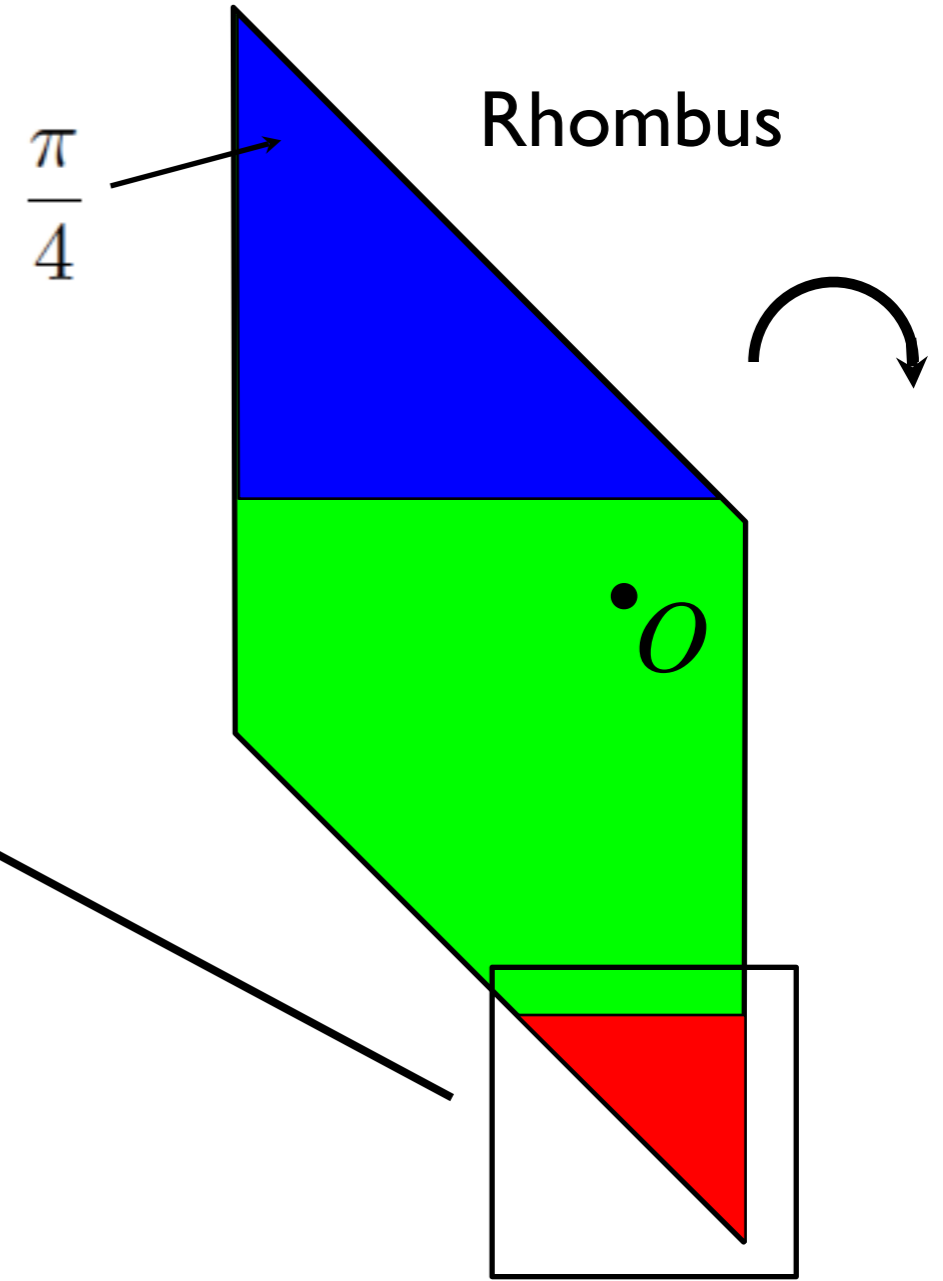
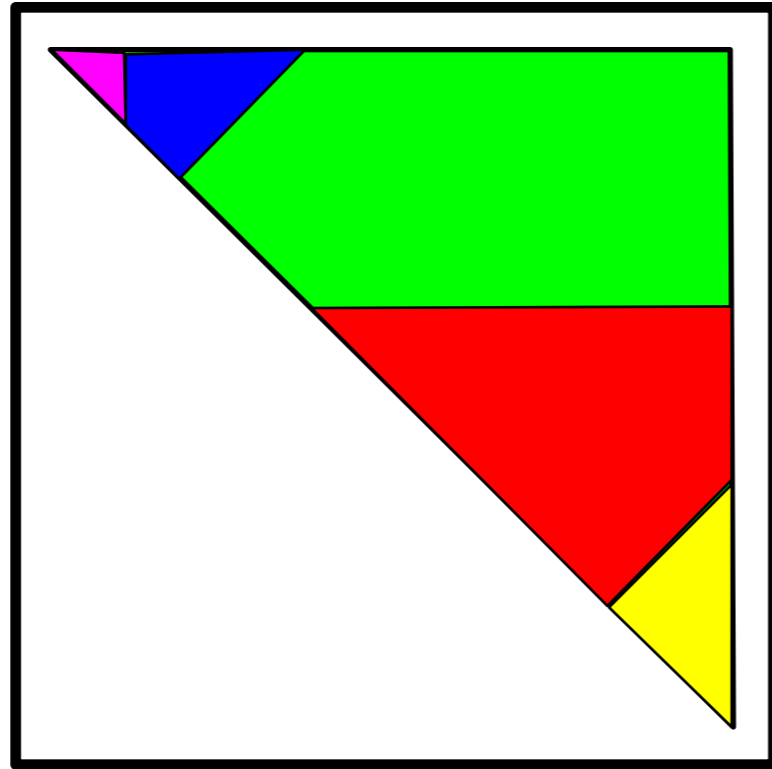


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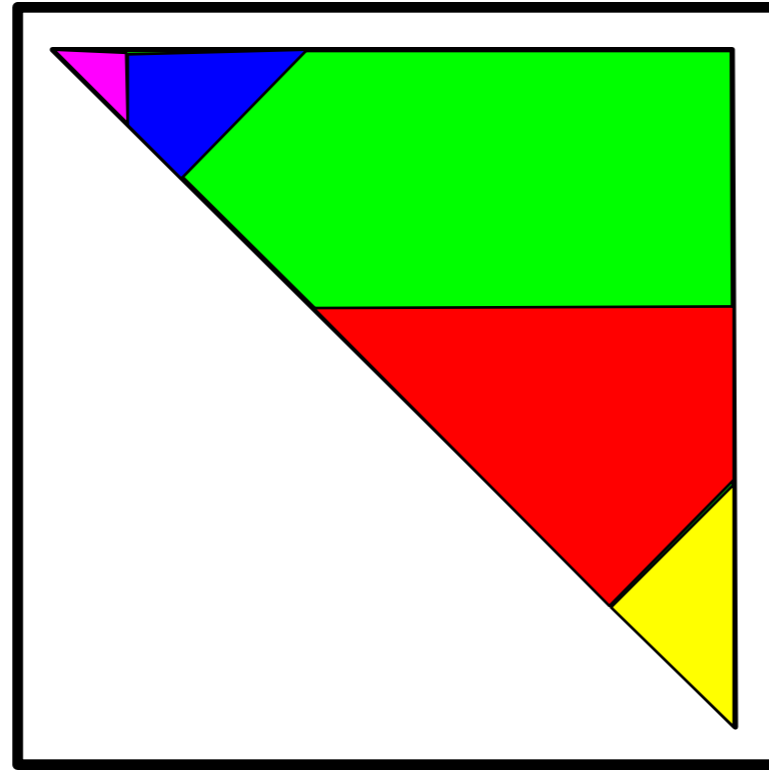
# The octagonal model (field $\mathbb{Q}(\sqrt{2})$ )

5-atom base triangle



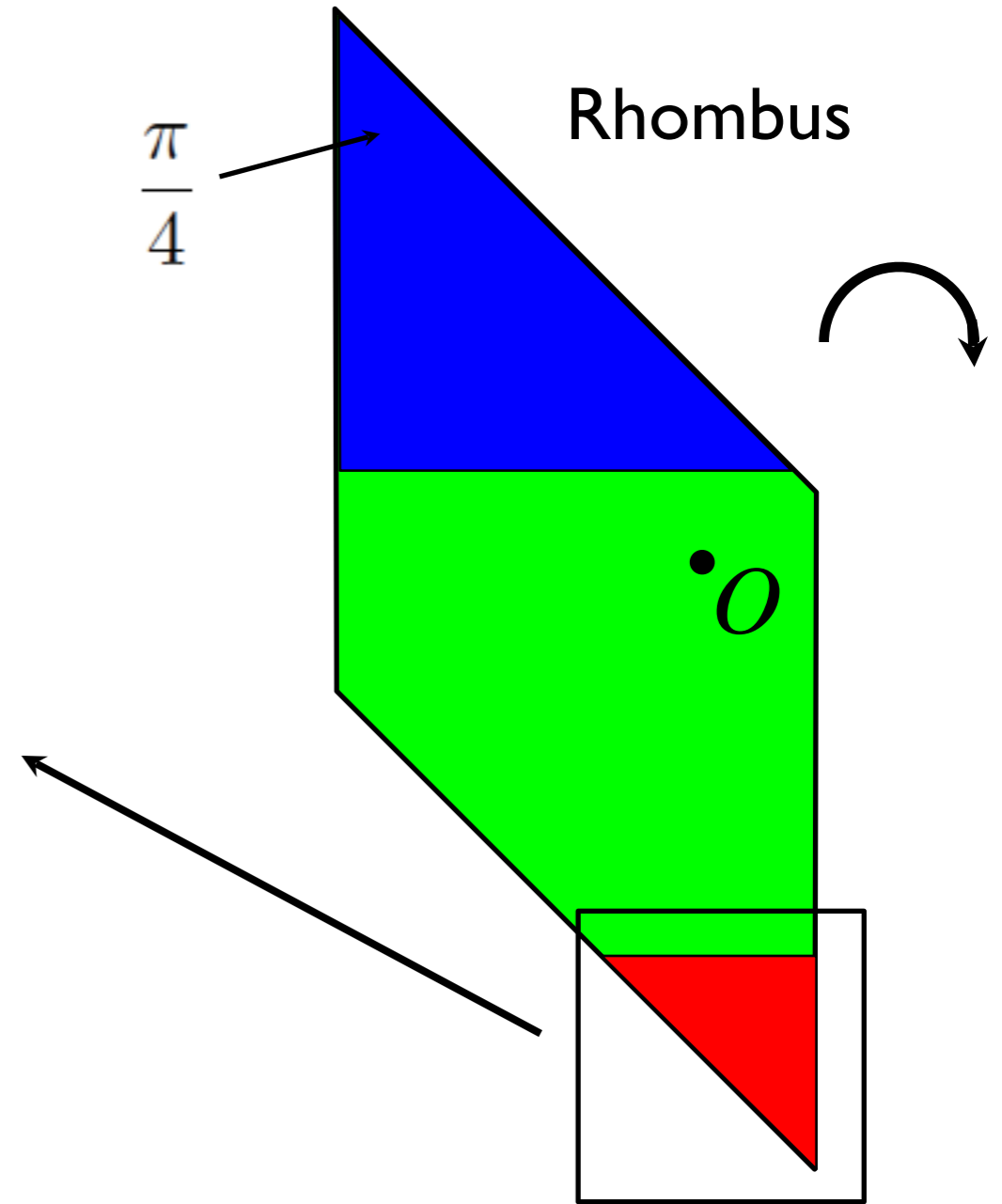
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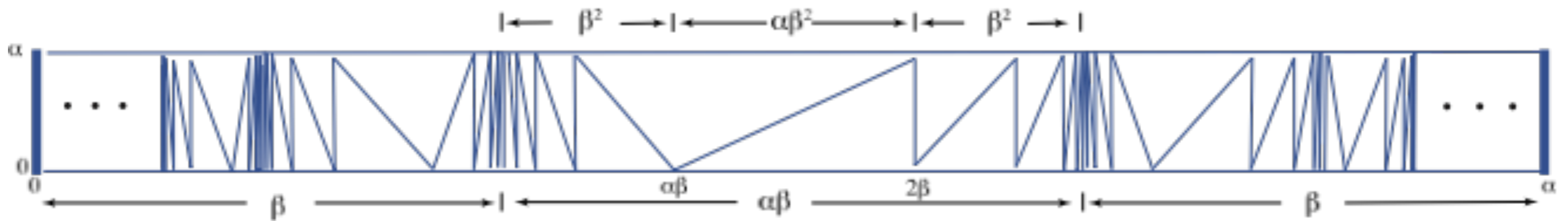


$$\frac{\pi}{4}$$

Rhombus

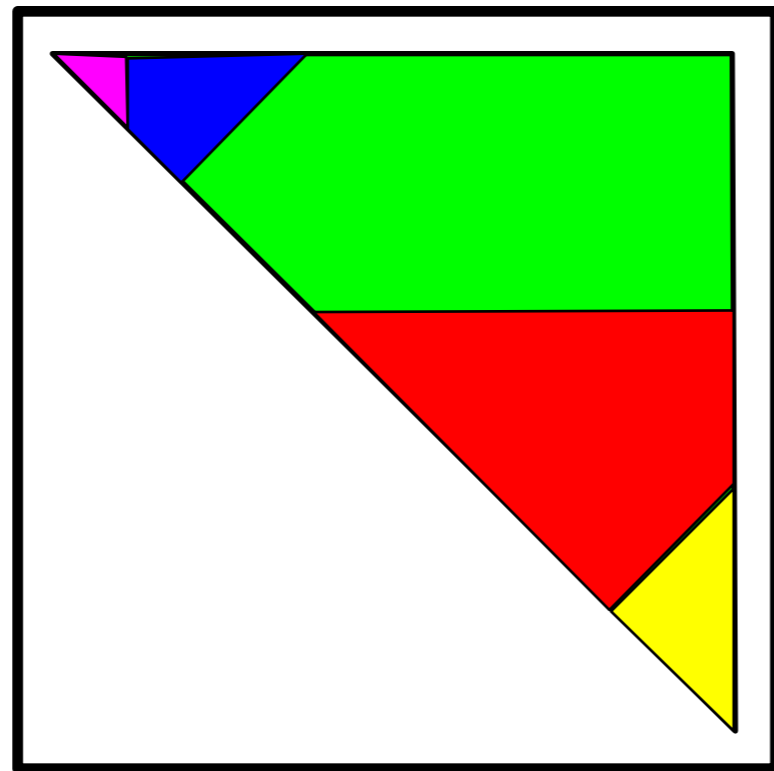


## The renormalization function



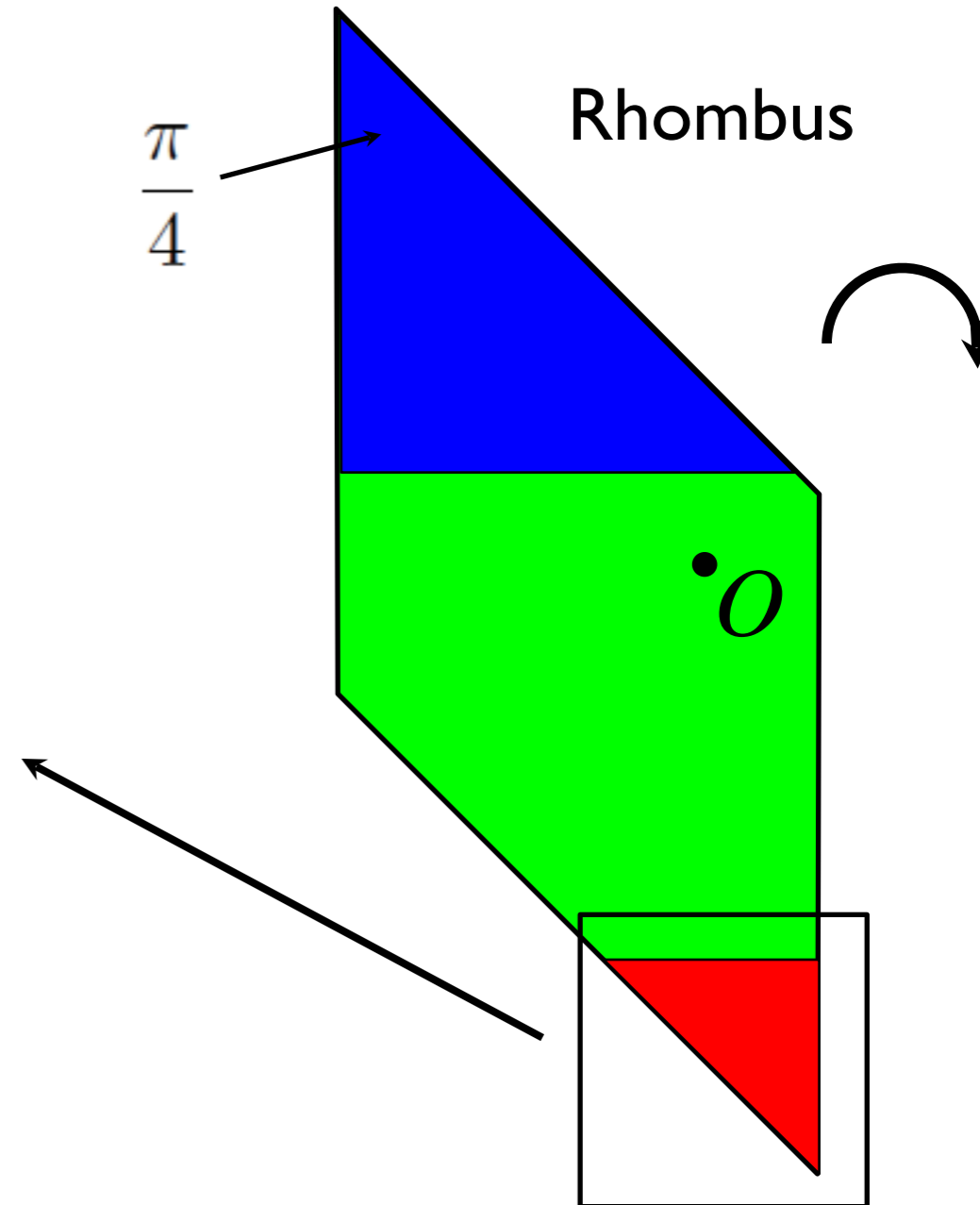
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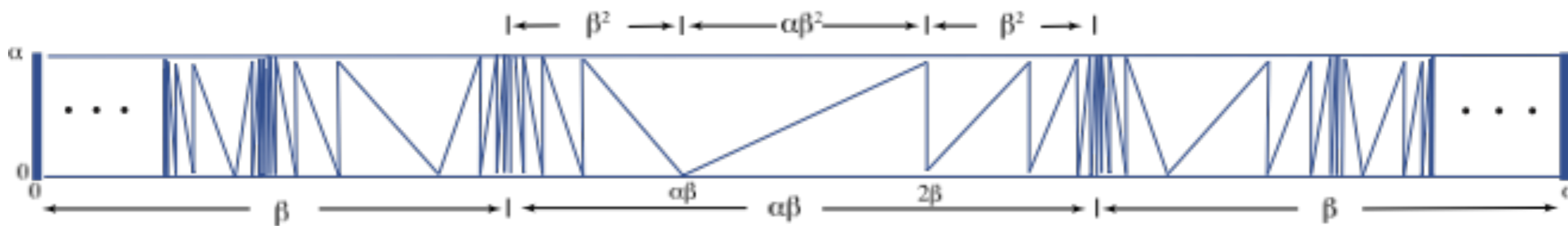


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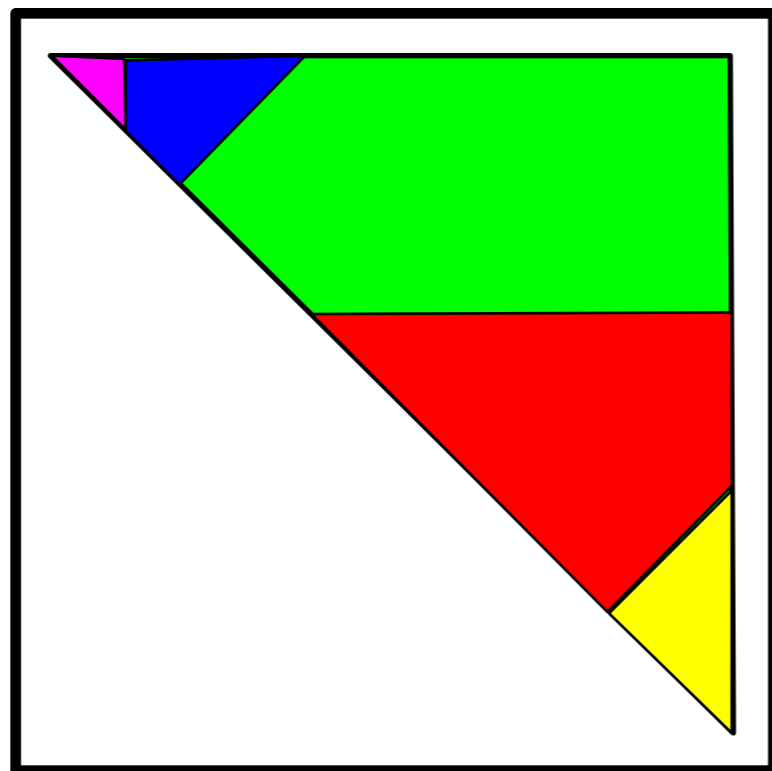


$$r(s) = f^2(s)$$



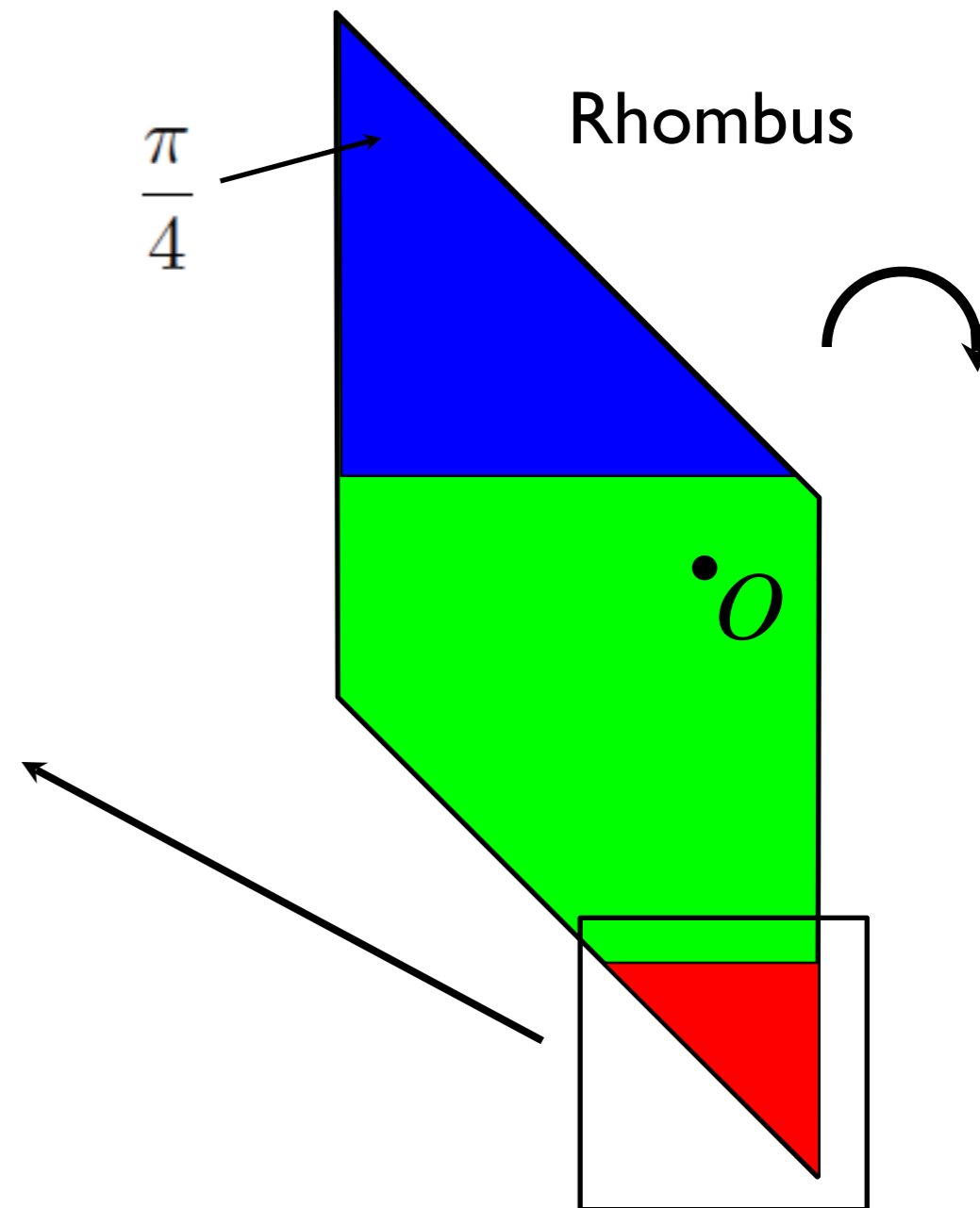
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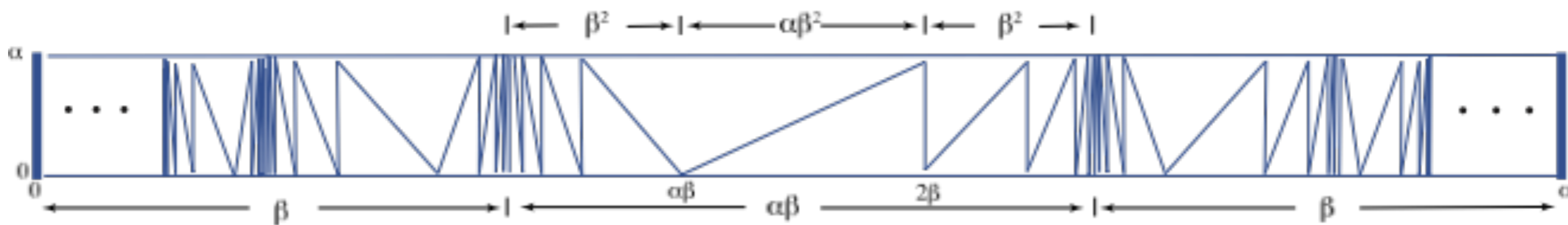


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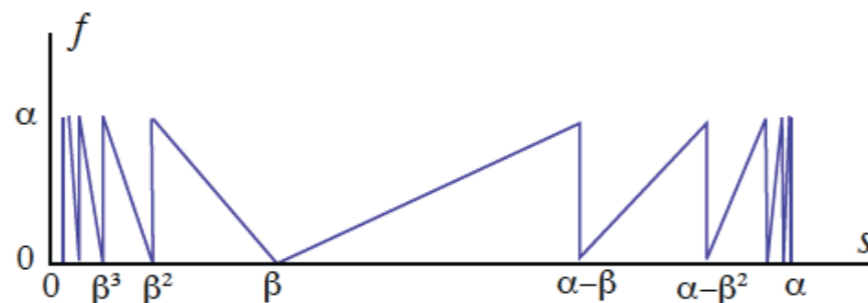
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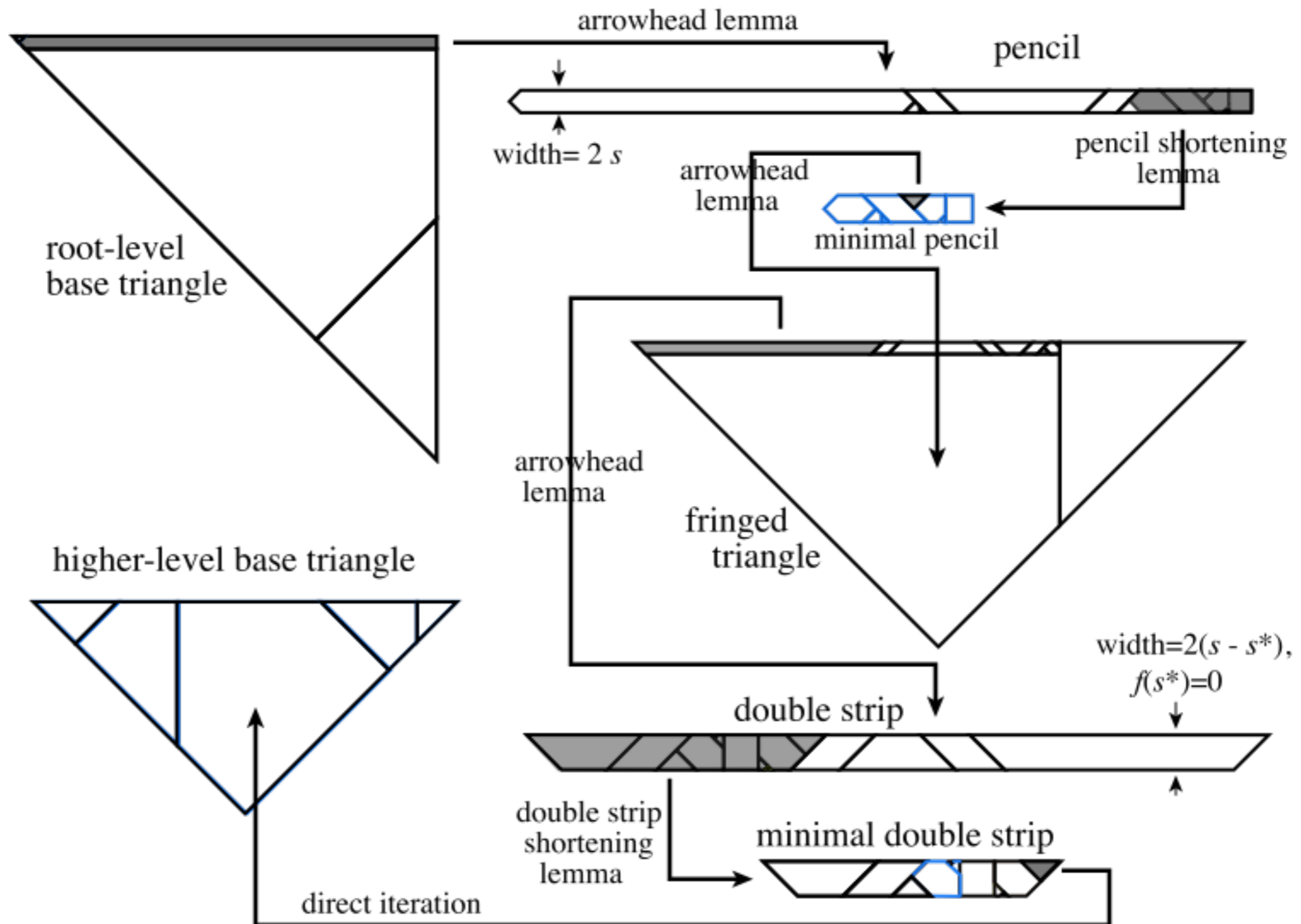




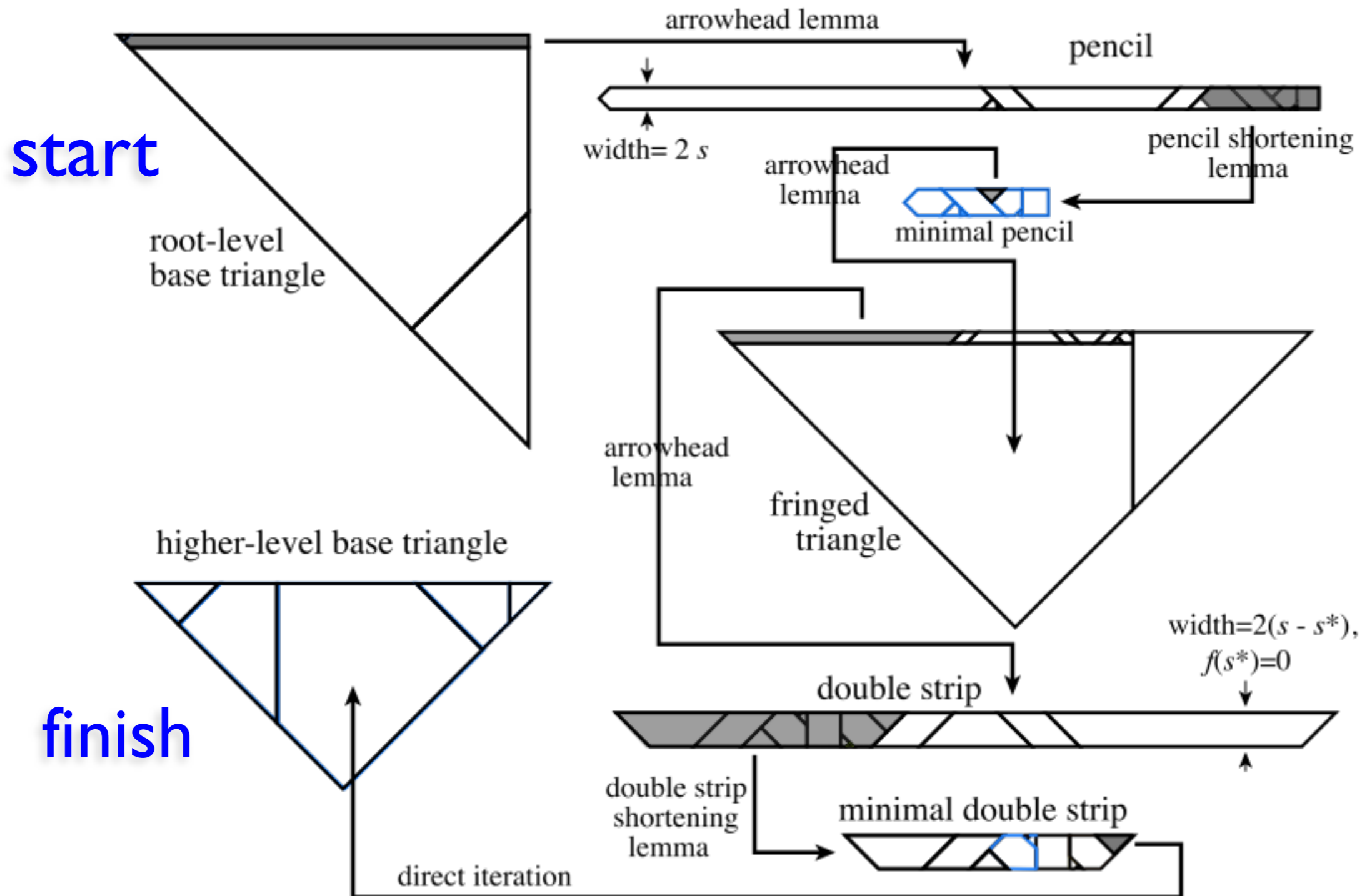
■ Very elaborate induction process

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More than one parameter: re-combination of atoms

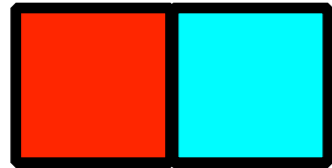
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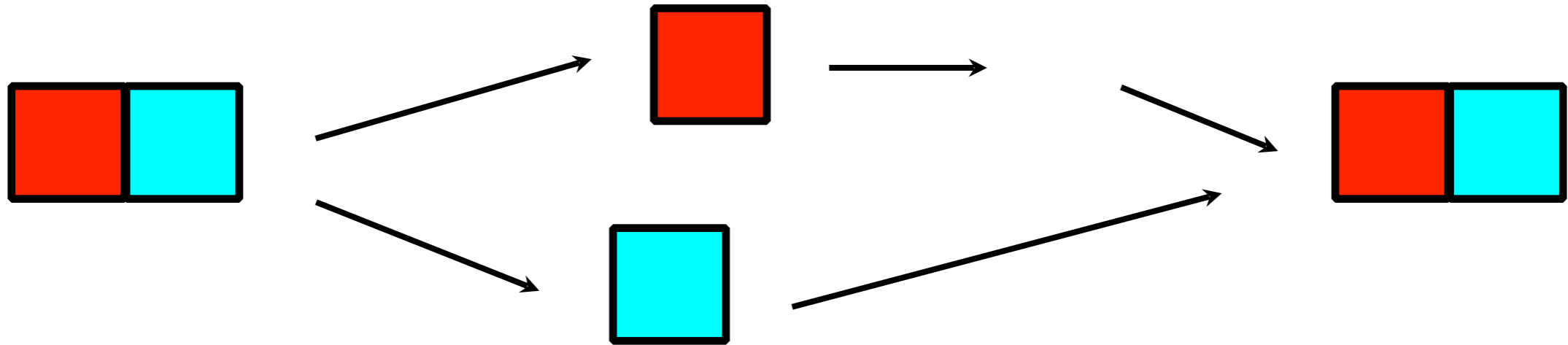
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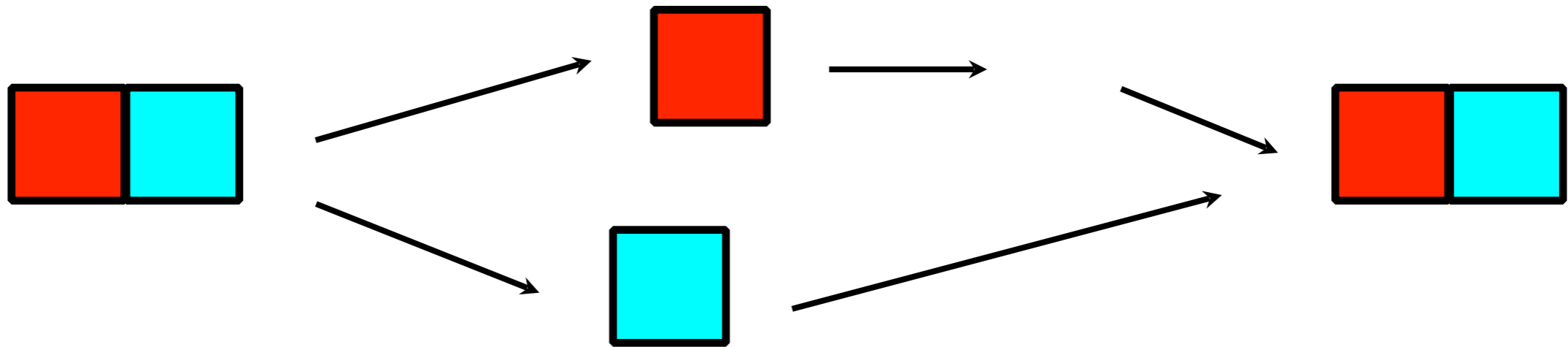
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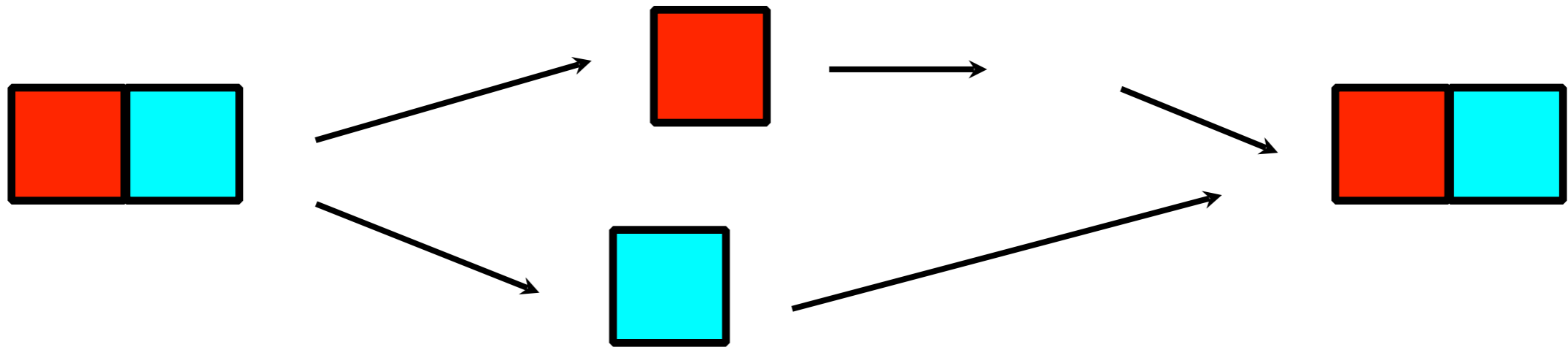
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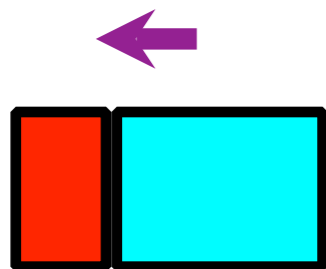
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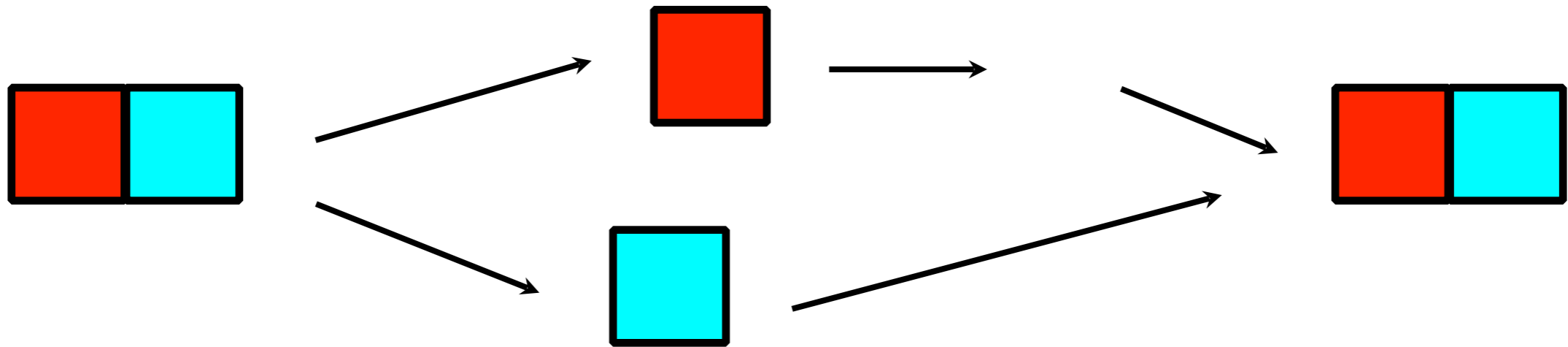


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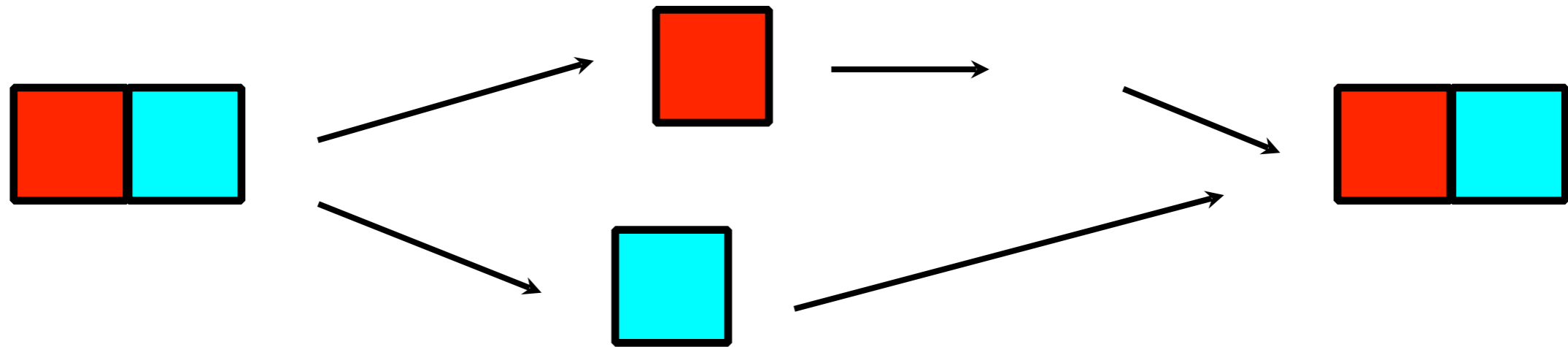


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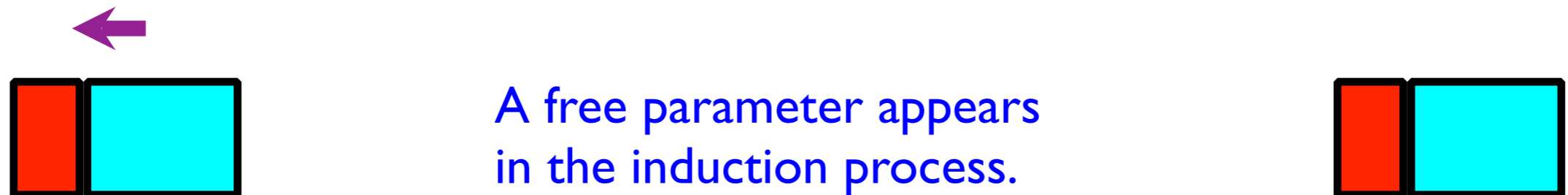


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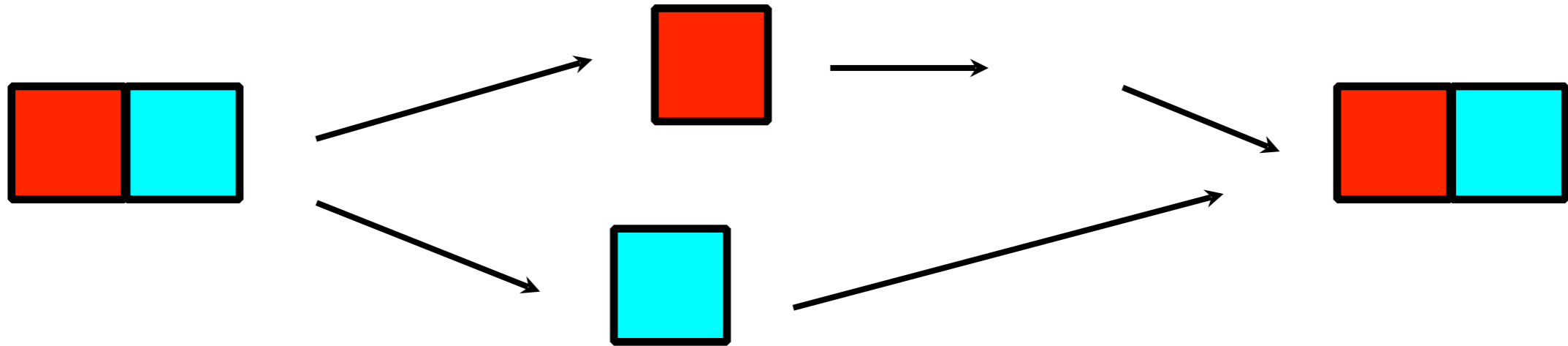


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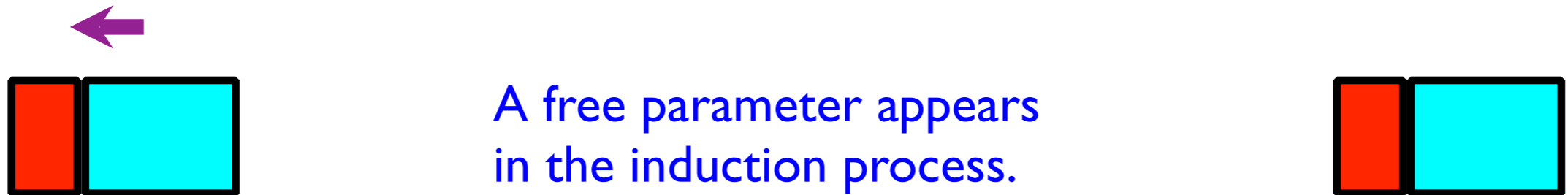


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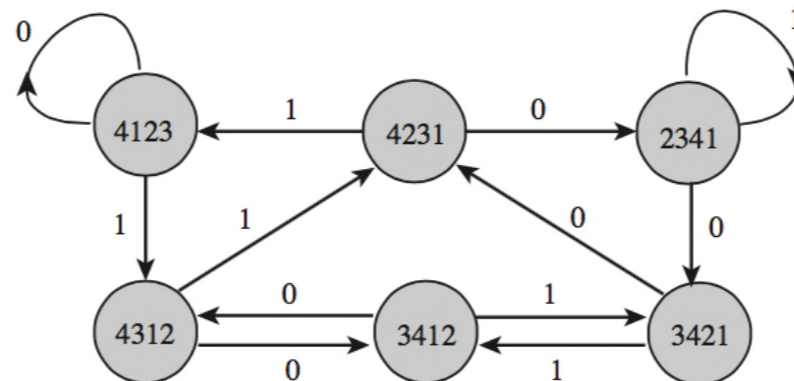


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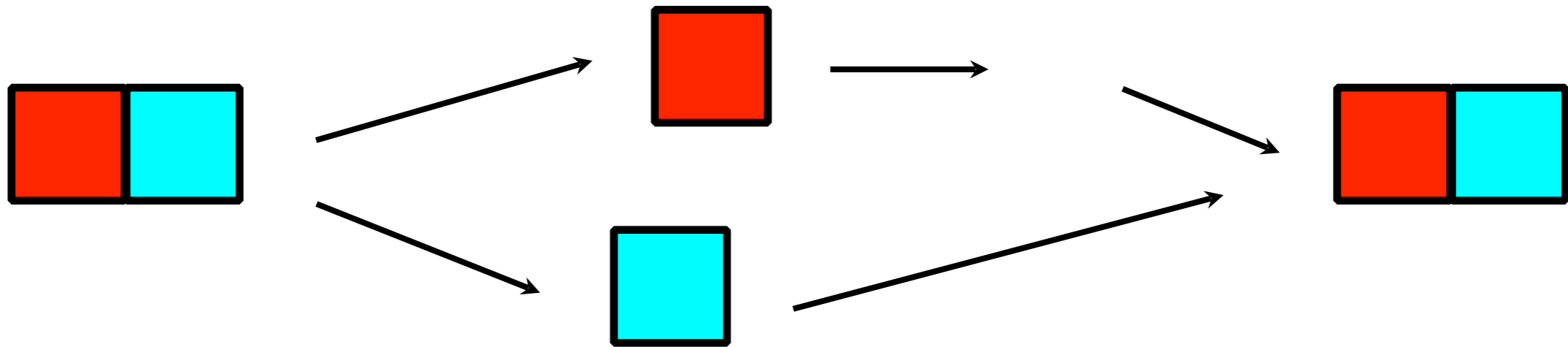
IETs:

Rauzy class

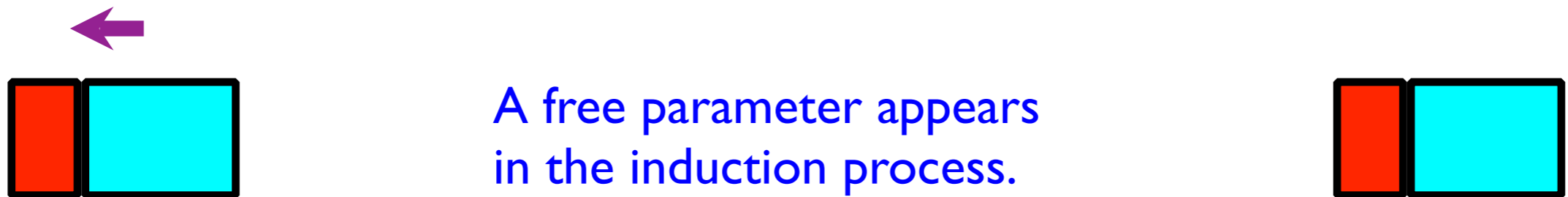


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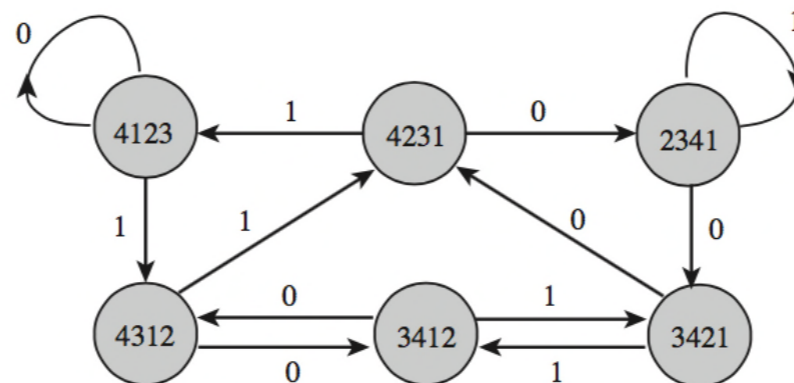


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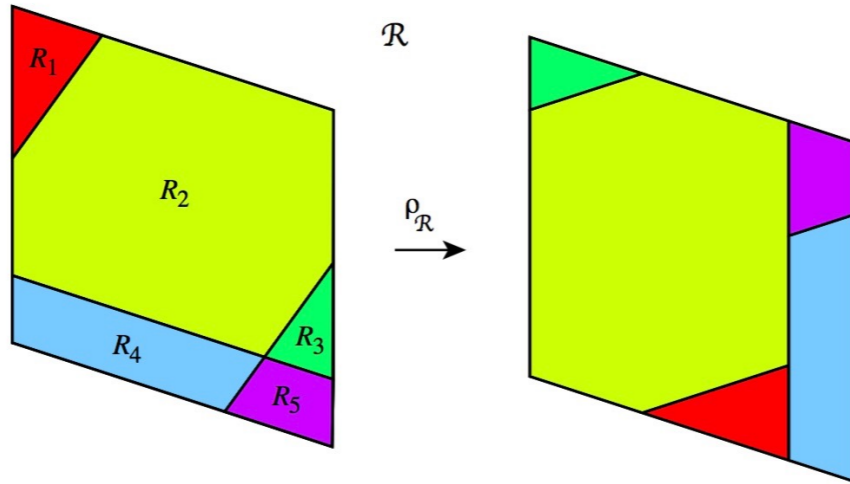


All permutations have a pair of consecutive indices

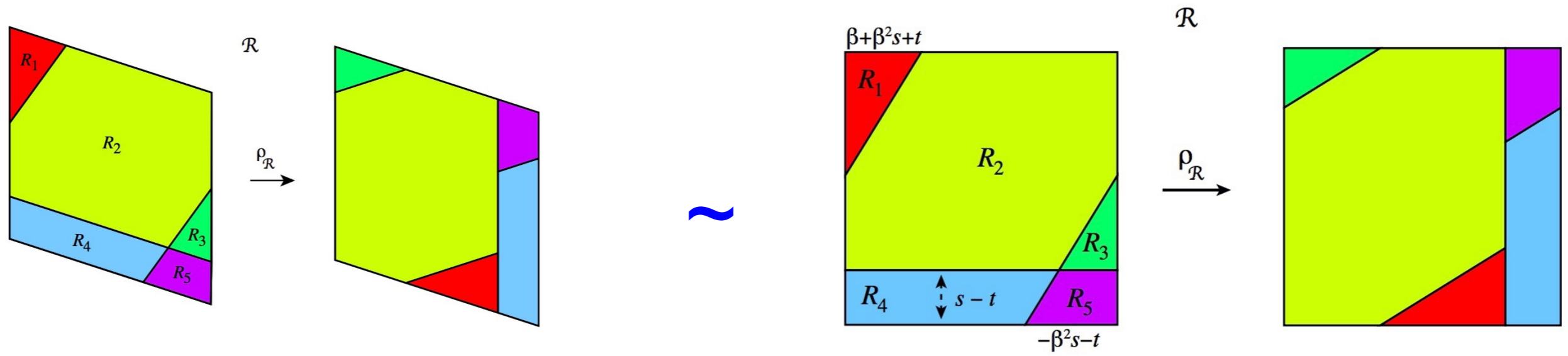


Two-parameter PETs: degenerate renormalisation (field  $\mathbb{K} = \mathbb{Q}(\sqrt{5})$ )

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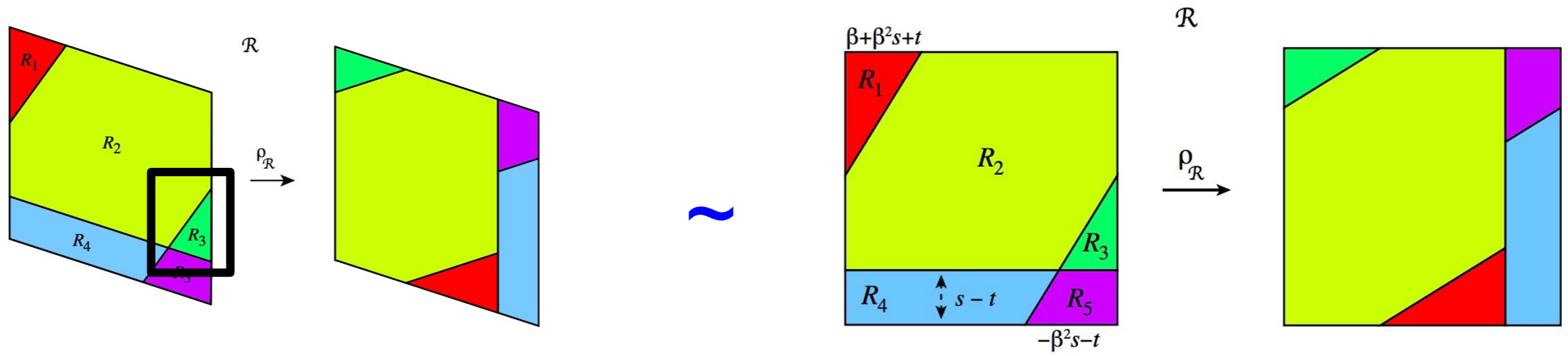


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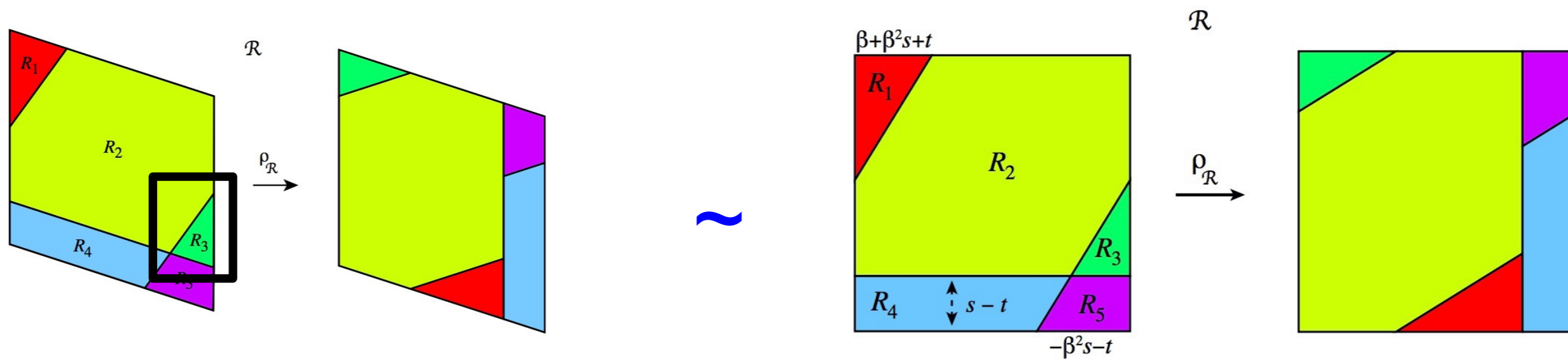
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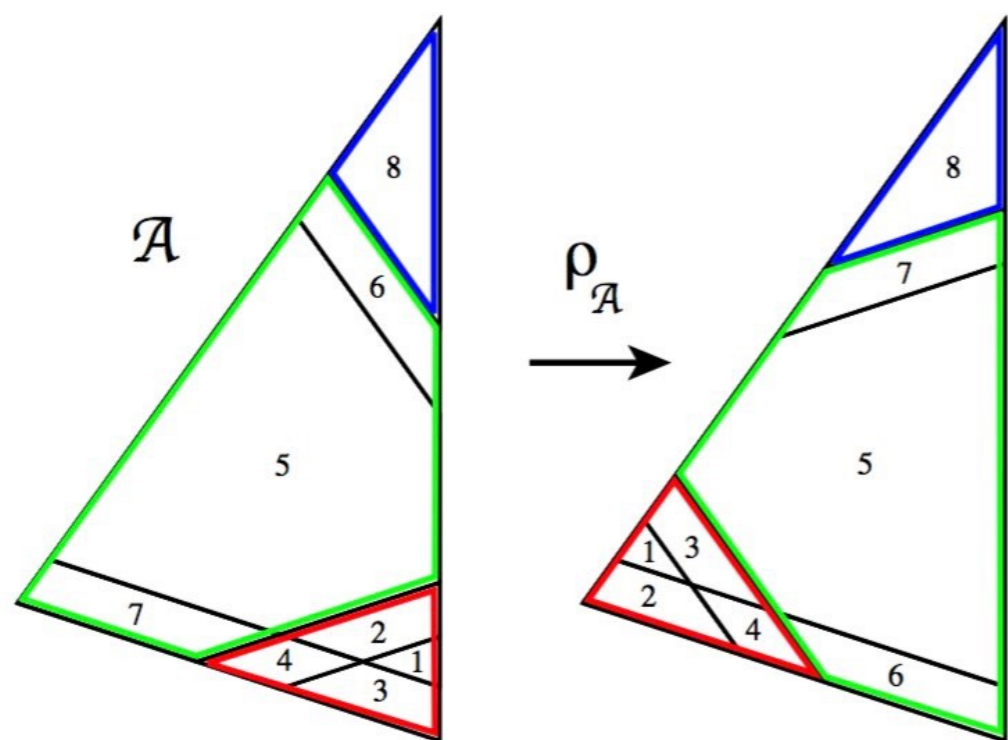


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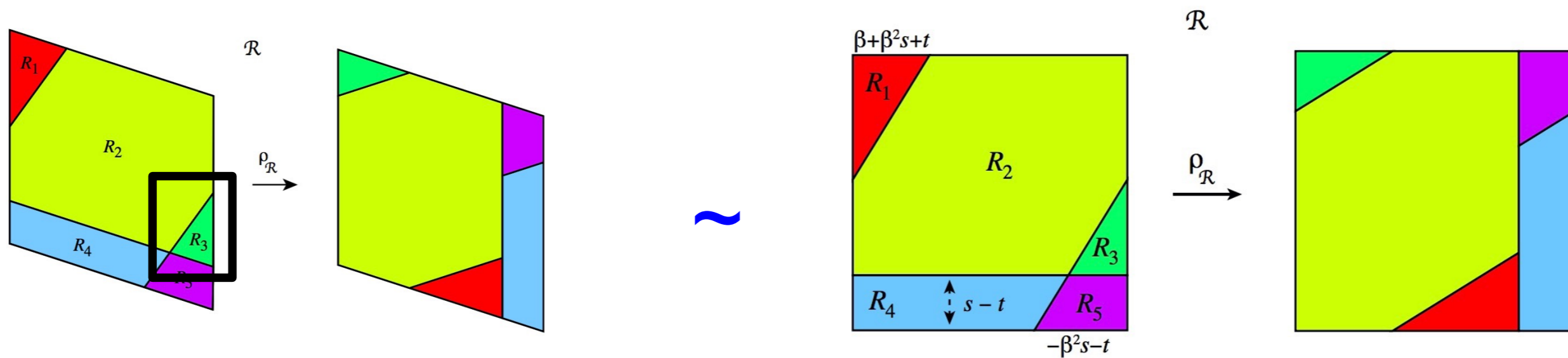


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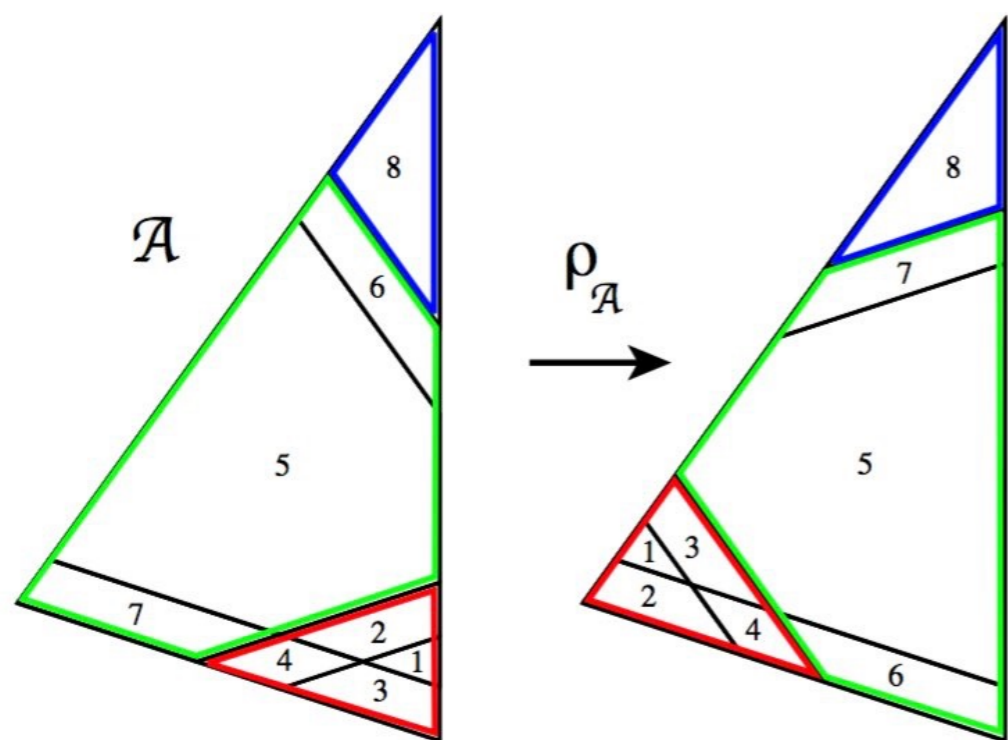


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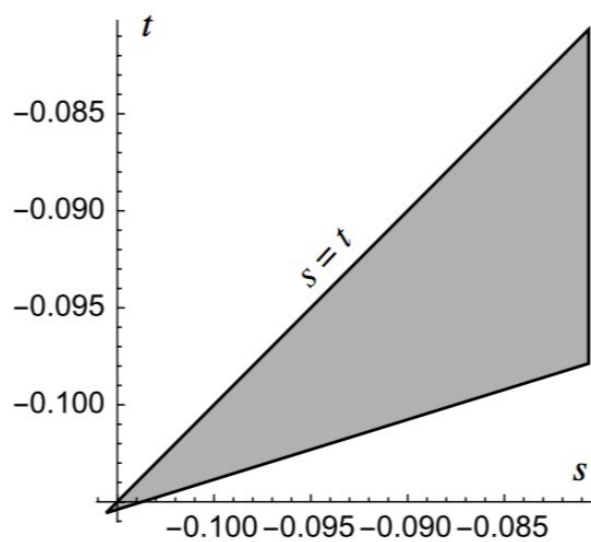
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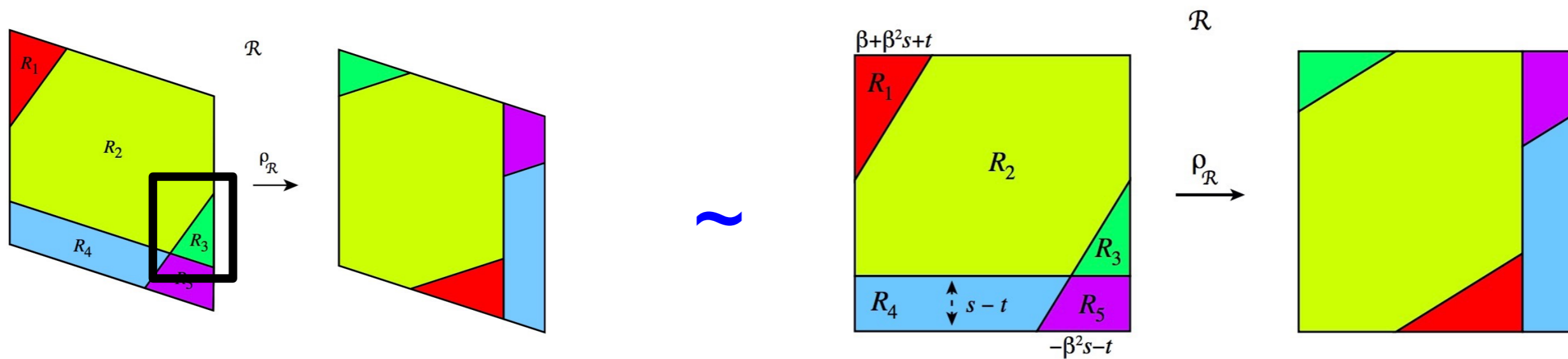


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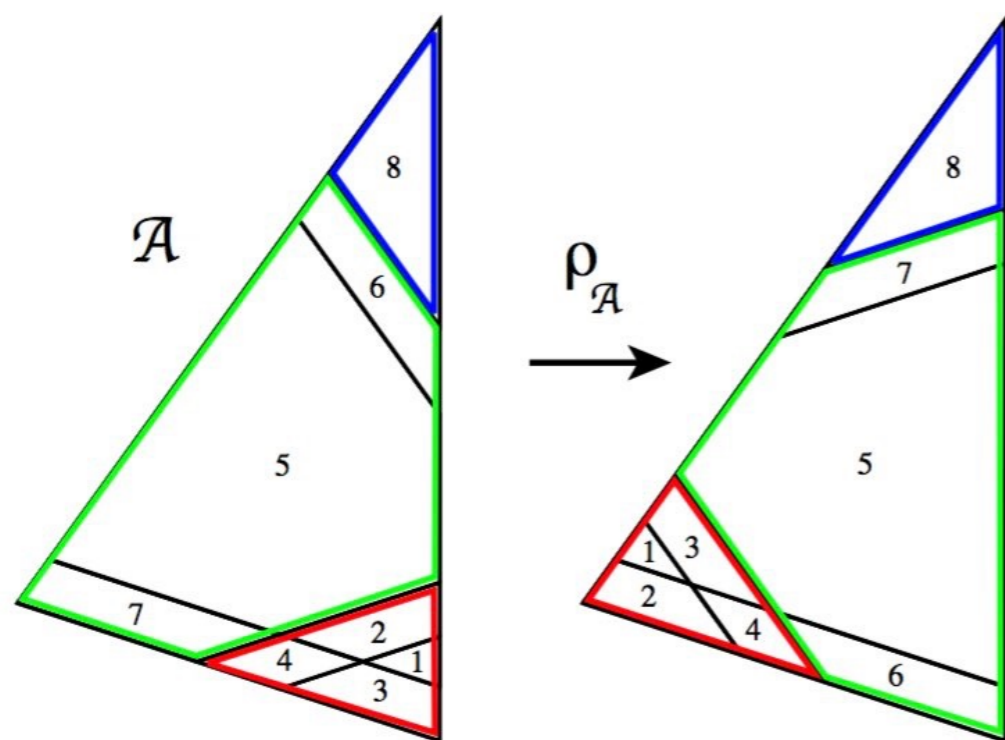


bifurcation-free parametric domain

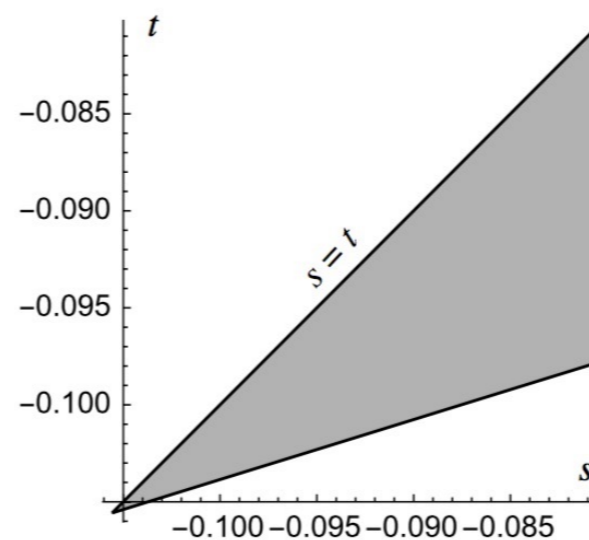
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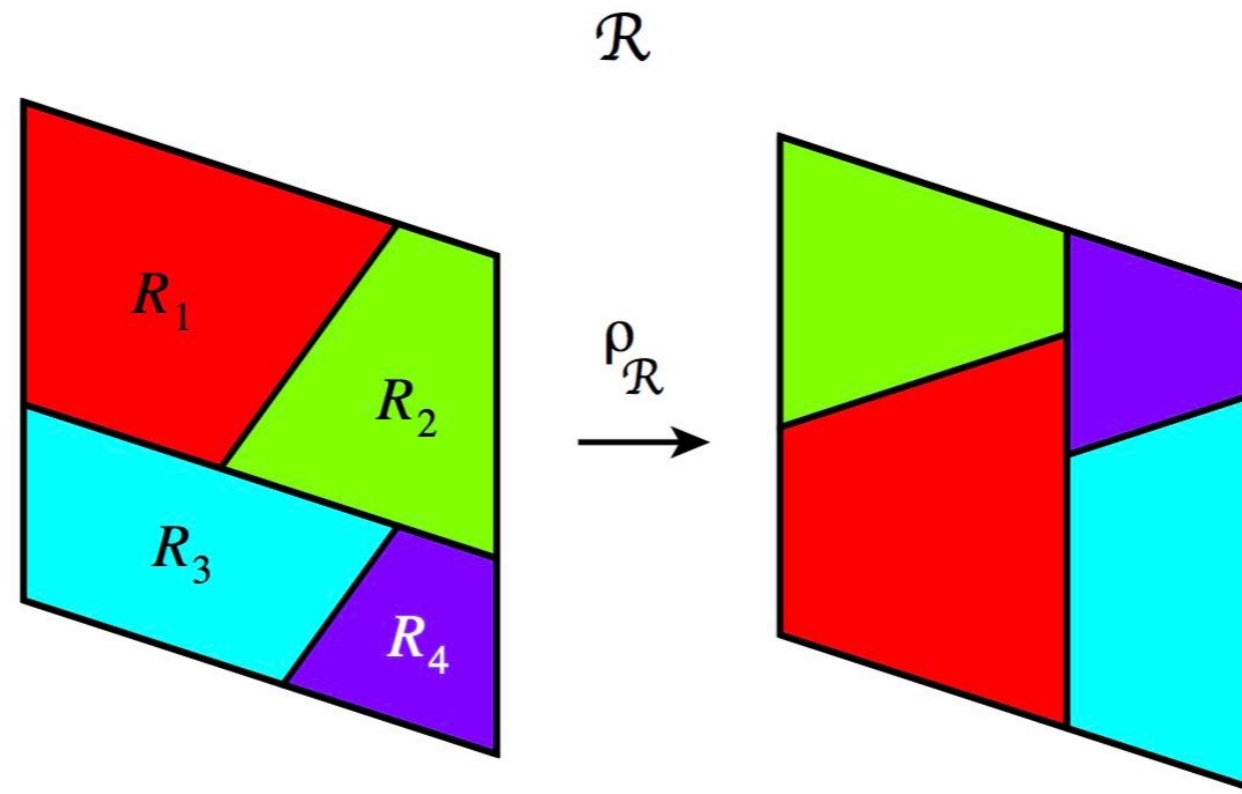
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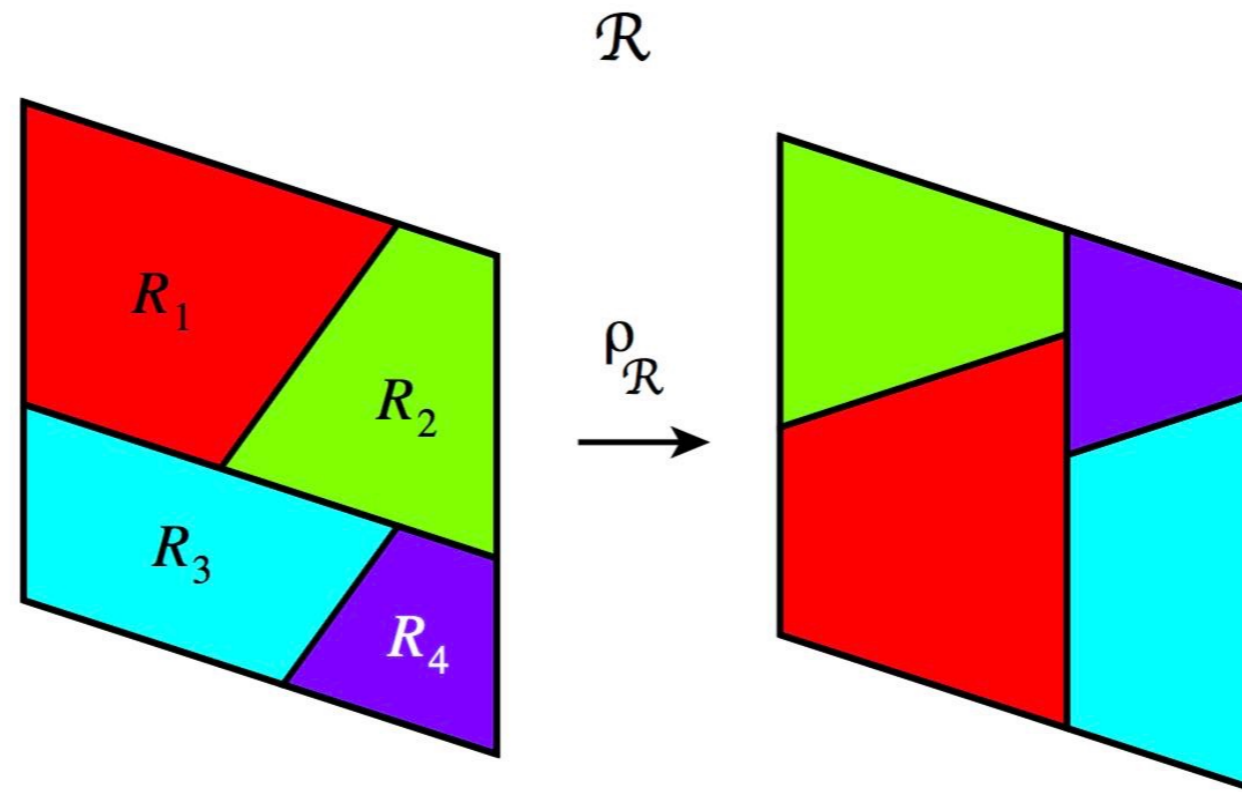
For renormalizability,  $s$  must be constrained to  $\mathbb{K}$ , while  $t$  is unconstrained.

# A weakly non-degenerate renormalisation, with two parameters



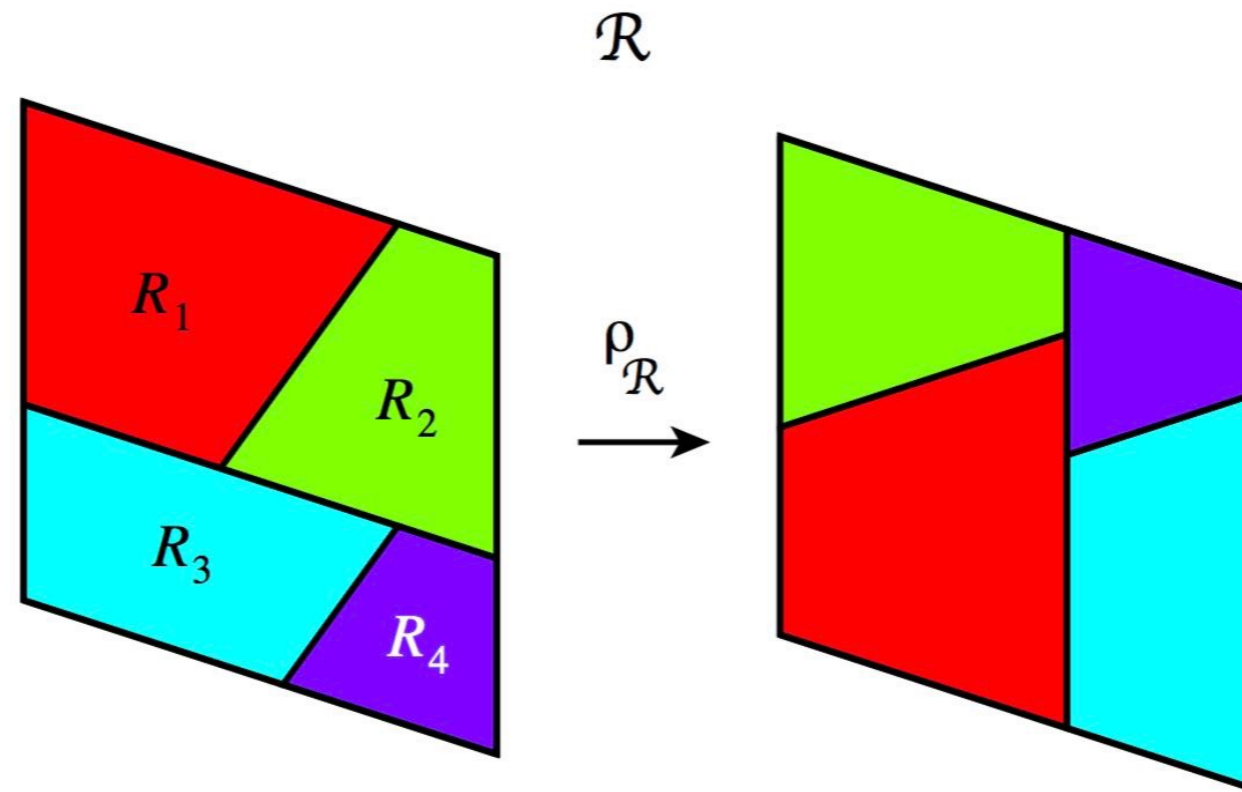


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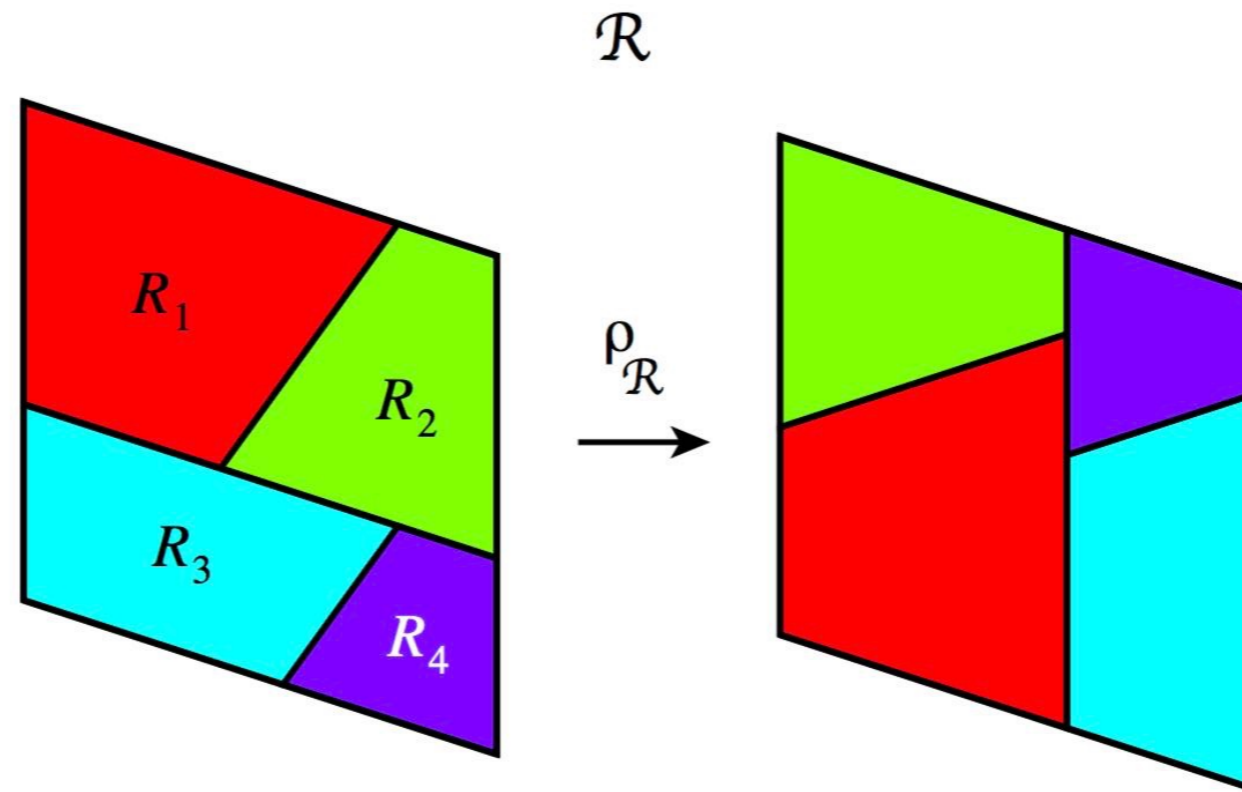
- The renormalization domain splits into three components.

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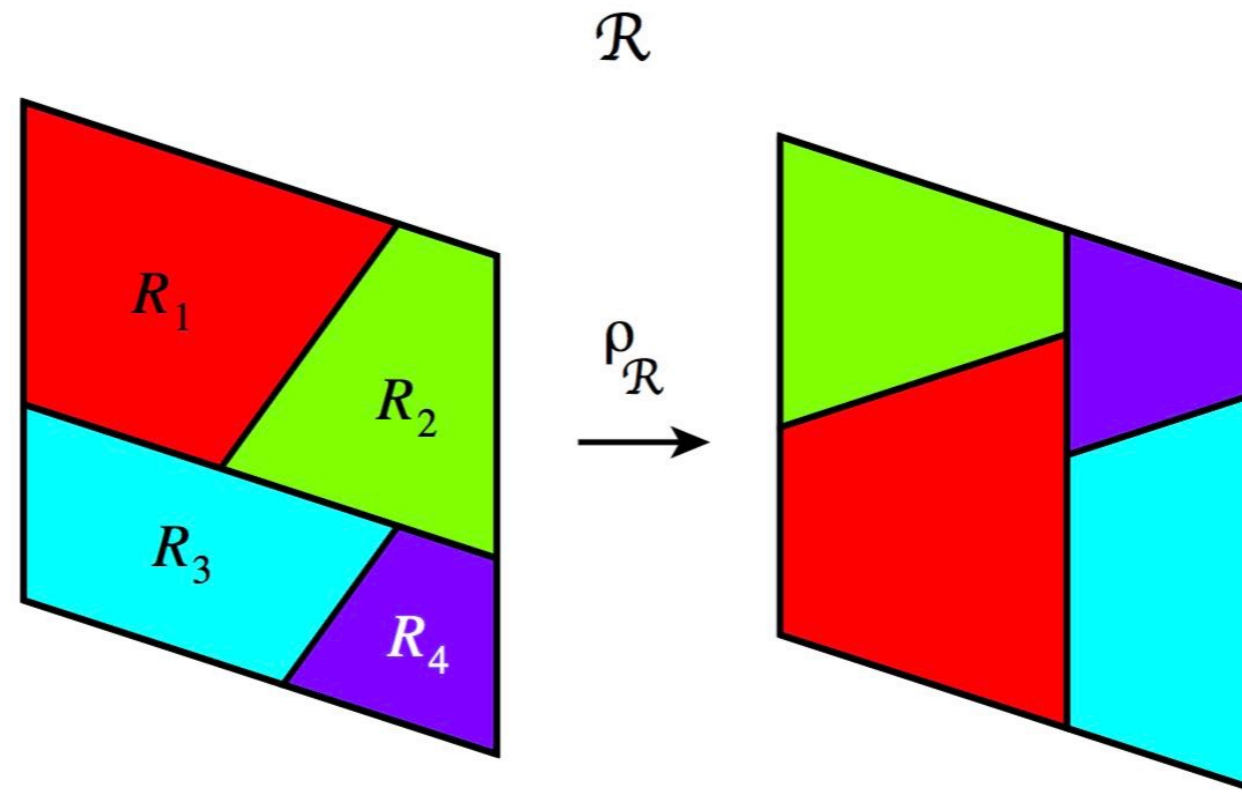
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- The system is renormalisable iff both parameters belong to  $\mathbb{K}$ .

Thank you for your attention

