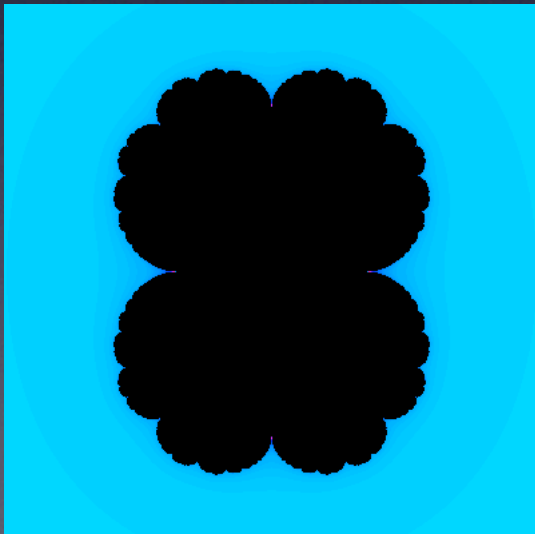


Applications of Parabolic Renormalization

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Local Holomorphic Dynamics

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When a parabolic point is perturbed....

$$f_0(z) = z + z^2$$

Discontinuous change of Julia sets, etc.

Inou's talk

$$f_0(z) = z + a_2 z^2 + \dots \quad a_2 \neq 0$$

$$E_{f_0} = \Phi_{attr} \circ \Phi_{rep}^{-1} \quad \text{horn map} = \text{Ecalle-Voronin invariant}$$

Want to consider iteration

$$\mathbb{C}/\mathbb{Z} \ni z, \quad E_{f_0}(z), \quad E_{f_0}^2(z), \quad E_{f_0}^3(z), \dots$$

$$\mathcal{R}_0 f_0 = \text{Exp}^\sharp \circ E_{f_0} \circ (\text{Exp}^\sharp)^{-1} \quad \text{Parabolic renormalization}$$

$$\text{Exp}^\sharp(z) = e^{2\pi iz} : \mathbb{C}/\mathbb{Z} \xrightarrow{\cong} \mathbb{C}^*$$

Want to iterate the process

$$f_0 \rightsquigarrow \mathcal{R}_0 f_0 \rightsquigarrow \mathcal{R}_0^2 f_0 \rightsquigarrow \mathcal{R}_0^3 f_0 \rightsquigarrow \dots$$

$$\text{Main Theorem I: } \exists \mathcal{F}_1 \quad \mathcal{R}_0(\mathcal{F}_1) \subset \mathcal{F}_1$$

Continued fraction and Renormalization

(for rotation-like dynamics)

$$\alpha = \frac{1}{a_1 + \alpha_1} = \frac{1}{a_1 + \frac{1}{a_2 + \alpha_2}} = \dots = \frac{1}{a_1 + \frac{1}{a_2 + \frac{1}{\ddots}}}$$

$$\alpha_0 = \alpha \in (0, 1) \setminus \mathbb{Q}, \quad a_n = \left[\frac{1}{\alpha_{n-1}} \right] \in \mathbb{N}, \quad \alpha_n = \frac{1}{\alpha_{n-1}} - a_n \in (0, 1)$$

$$\alpha \longleftrightarrow T_\alpha : x \mapsto x + \alpha \text{ on } \mathbb{R}/\mathbb{Z}$$

$$\text{convergents } \frac{p_n}{q_n} \rightarrow \alpha \longleftrightarrow T_\alpha^{q_n} \rightarrow id \text{ "closest returns"}$$

first return map $\doteq T_\alpha^{q_n} \doteq (T_\alpha^{q_{n-1}})^{a_n}$



rescale



Renormalization for rotation-like dynamics

The same construction for non-linear maps

$$f_0 = f$$

$f_{n+1} =$ (first return map of f_n to a fundamental domain)
up to rescaling

$$\alpha_n = \text{rotation number of } f_n \qquad f_{n+1} \doteq f_n^{\alpha_n}$$

Yoccoz renormalization for Siegel-Bruno Theorem

$f(z) = e^{2\pi i \alpha} z + \dots$, $\sum \log \frac{q_{n+1}}{q_n} < \infty \implies f$ is conjugate to $z \mapsto e^{2\pi i \alpha} z$

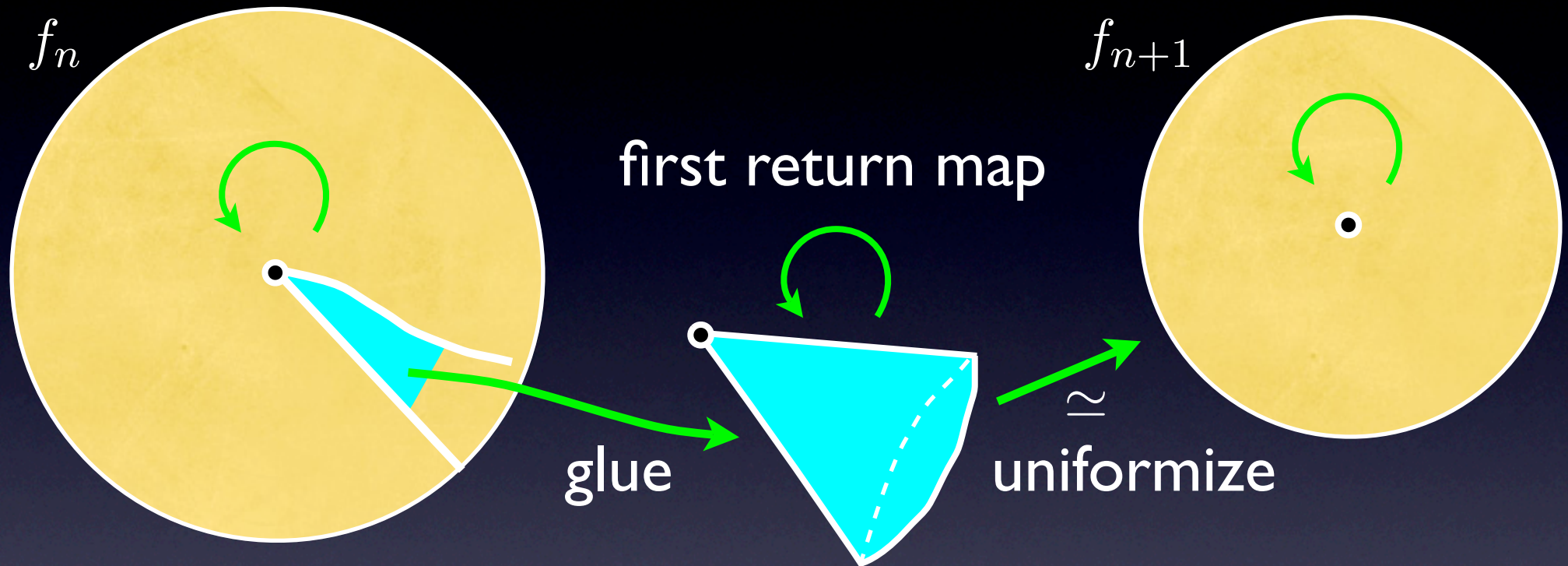
$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$

= first return map of f_n
to a fundamental domain
up to *uniformization*

$$\alpha_{n+1} \equiv -\frac{1}{\alpha_n} \pmod{\mathbb{Z}}$$

Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$



Optimality of Bruno condition (Yoccoz)

inverse construction for germs (*unrenormalization*)

can add extra fixed / periodic points

another argument for quadratic polynomials

Near-parabolic renormalization (cylinder renormalization)

take crescent-shaped fundamental region

quotient = cylinder $\mathbb{C}/\mathbb{Z} \simeq \mathbb{C}^*$

Advantages:

- quotient cylinder is canonical

- can include critical point

- when α is small the fundamental region does not shrink

- limit can be described by horn map

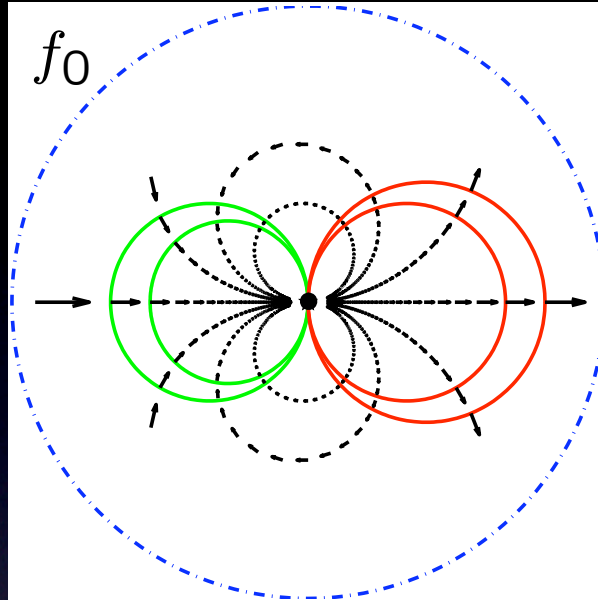
Disadvantage:

- only applies to α with large continued fraction coefficients $a_n \geq N$

Horn map and Parabolic Renormalization

$$f_0(z) = z + a_2 z^2 + \dots$$

$$a_2 \neq 0$$



Horn map

$$E_{f_0} = \Phi_{attr} \circ \Phi_{rep}^{-1}$$

Parabolic Renormalization

$$\mathcal{R}_0 f_0 = \text{Exp}^\# \circ E_{f_0} \circ (\text{Exp}^\#)^{-1}$$

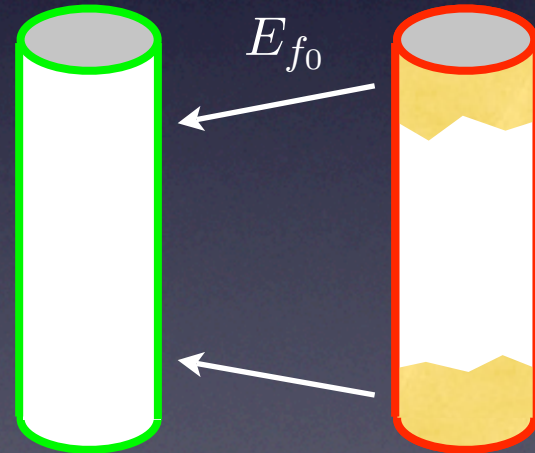
$$\text{Exp}^\#(z) = e^{2\pi i z} : \mathbb{C}/\mathbb{Z} \xrightarrow{\cong} \mathbb{C}^*$$

$$\mathcal{R}_0 f_0(z) = z + \dots$$

by normalization

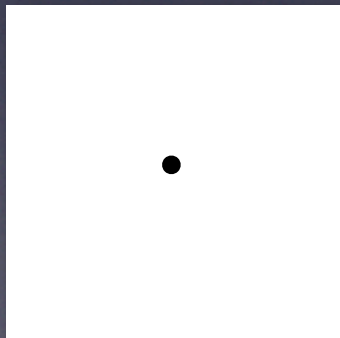
$$E_{f_0}(z) = z + o(1) \quad (\text{Im } z \rightarrow +\infty)$$

Φ_{attr} Φ_{rep}



$\text{Exp}^\#$

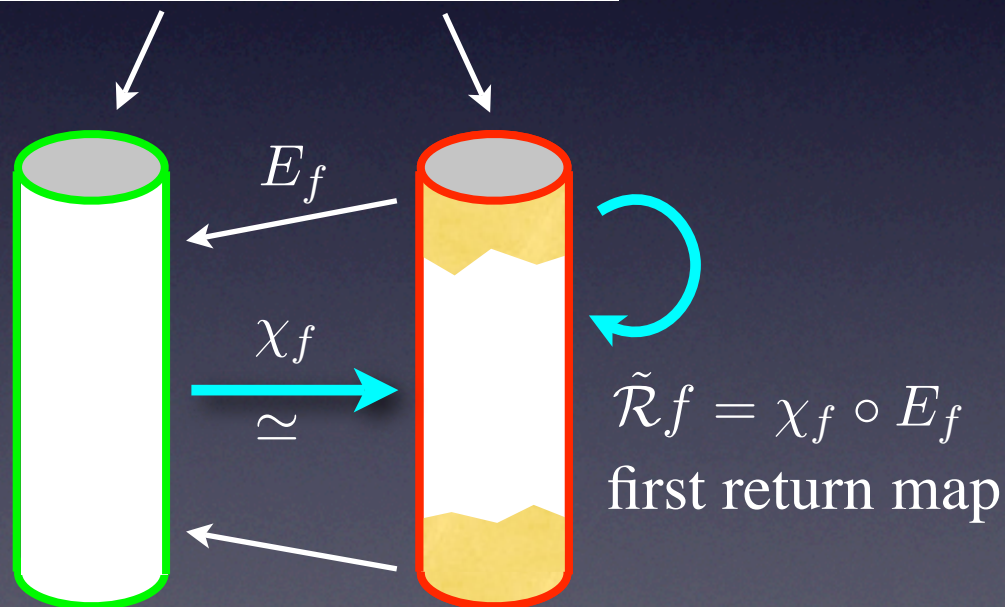
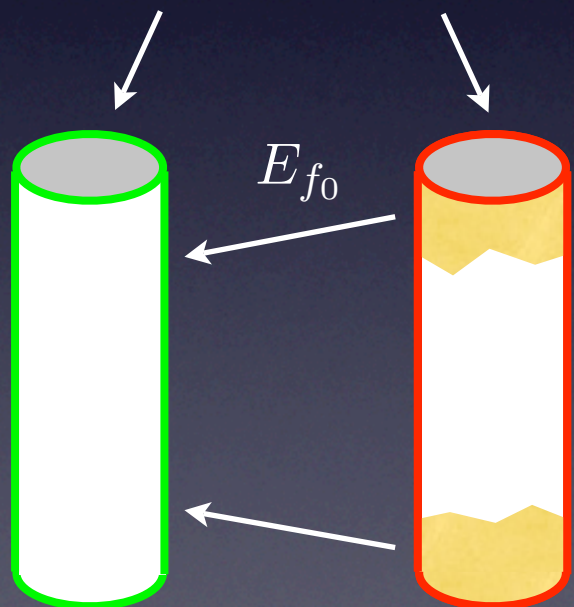
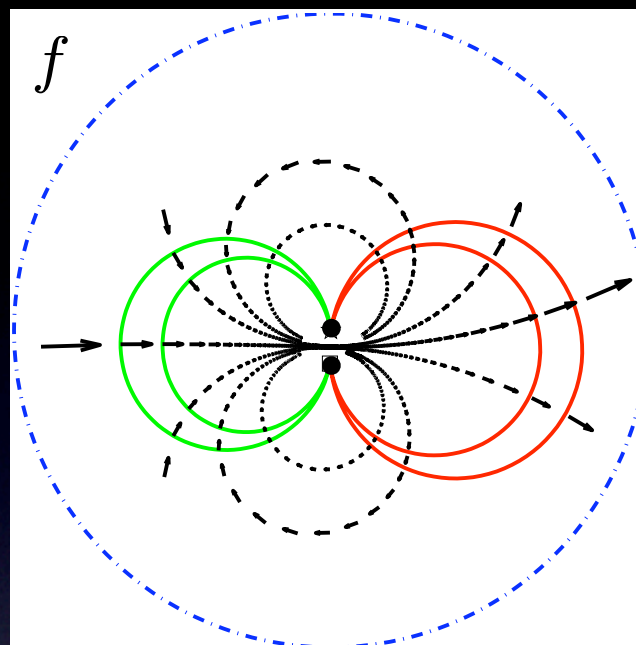
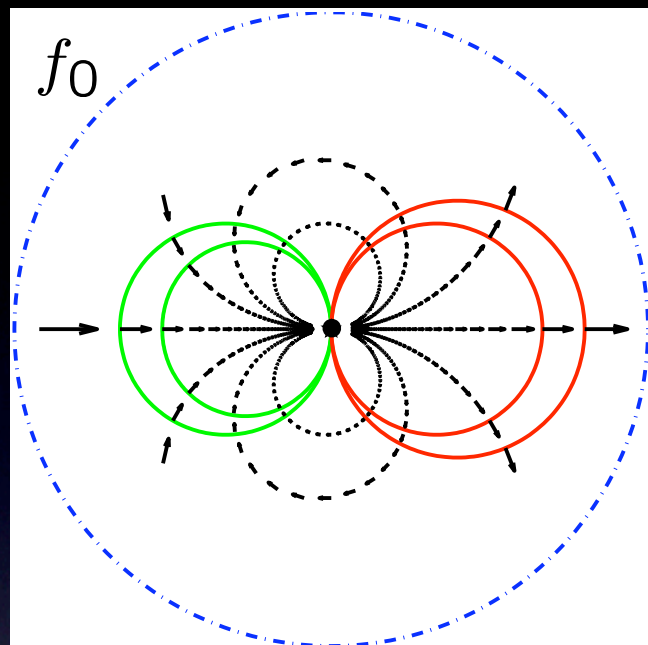
$\text{Exp}^\#$



$\mathcal{R}_0 f_0$

Perturbation

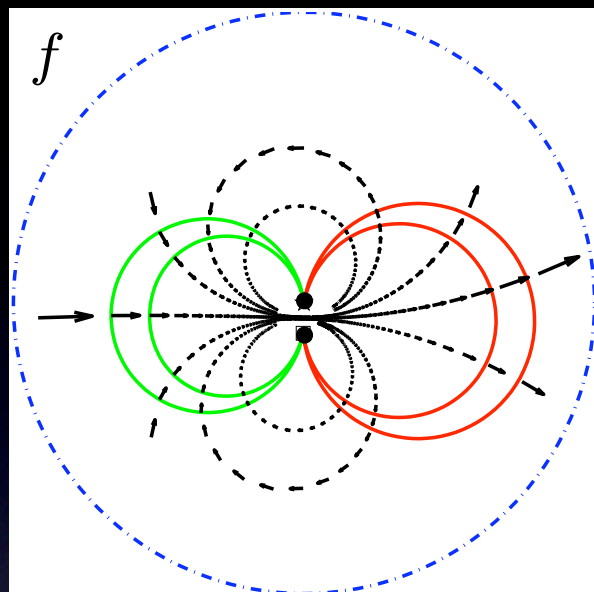
$$f'(0) = e^{2\pi i\alpha}, \quad \alpha \text{ small} \quad |\arg \alpha| < \frac{\pi}{4}$$



E_f depends continuously on f
(after a suitable normalization)

$$\chi_f(z) = z - \frac{1}{\alpha}$$

Near-parabolic Renormalization (cylinder renorm.)



$$\begin{aligned} \mathcal{R}f &= \text{Exp}^\# \circ \tilde{\mathcal{R}}f \circ (\text{Exp}^\#)^{-1} \\ &= \text{Exp}^\# \circ \chi_f \circ E_f \circ (\text{Exp}^\#)^{-1} \\ &= e^{2\pi i\beta} z + O(z^2) \end{aligned}$$

where $\beta = -\frac{1}{\alpha} \pmod{\mathbb{Z}}$

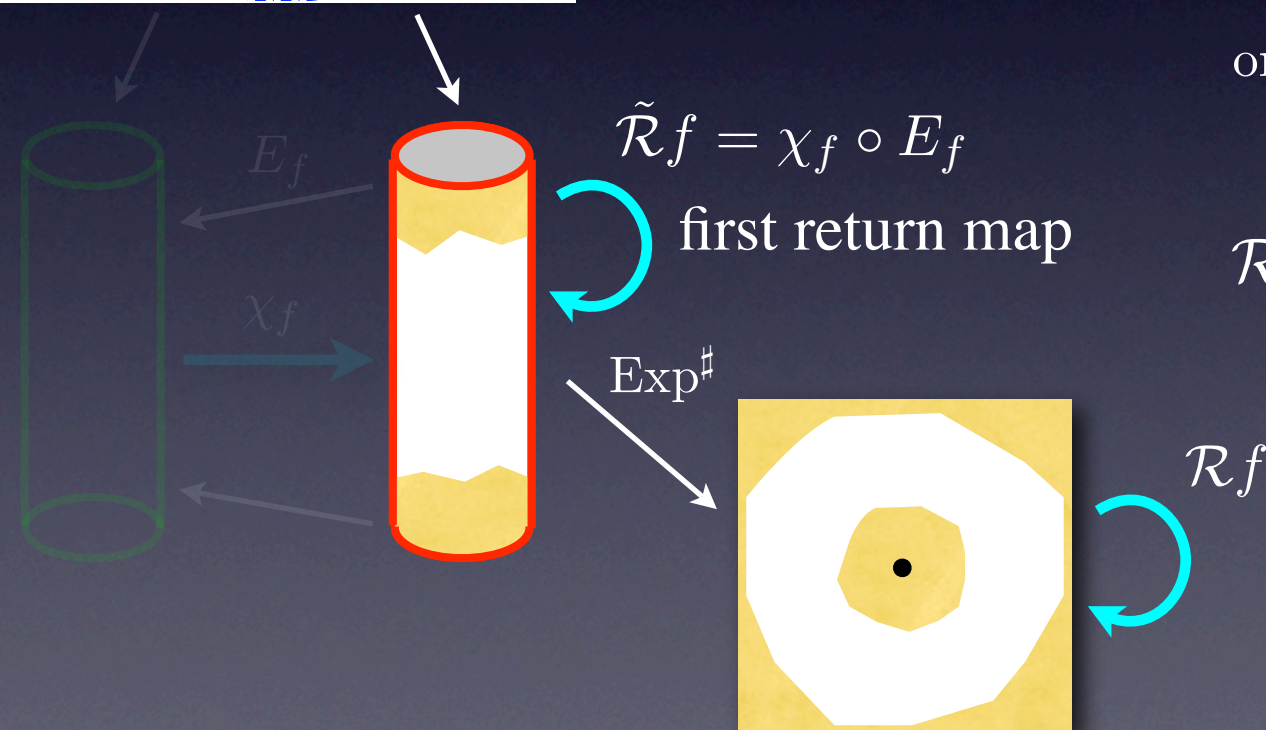
or $\alpha = \frac{1}{m - \beta} \pmod{\mathbb{Z}}$ ($m \in \mathbb{N}$)

$$\mathcal{R}f \doteq e^{2\pi i\beta} \mathcal{R}_0 f_0$$

write $f = e^{2\pi i\alpha} f_0$

$$\alpha \mapsto -\frac{1}{\alpha} \pmod{\mathbb{Z}}$$

$$f_0 \mapsto \mathcal{R}_0 f_0$$



Suppose the near-parabolic renormalization can be iterated:

$$f_0 \xrightarrow{\mathcal{R}} f_1 \xrightarrow{\mathcal{R}} f_2 \xrightarrow{\mathcal{R}} f_3 \xrightarrow{\mathcal{R}} \dots$$

high iterates of f_0 corresponds to low iterates of f_1

highly recurrent behavior of f_0 can be analyzed through f_i 's

fine structure of orbits or invariant sets are magnified

\mathcal{R} = a dynamical system in the space of certain type of dynamical systems

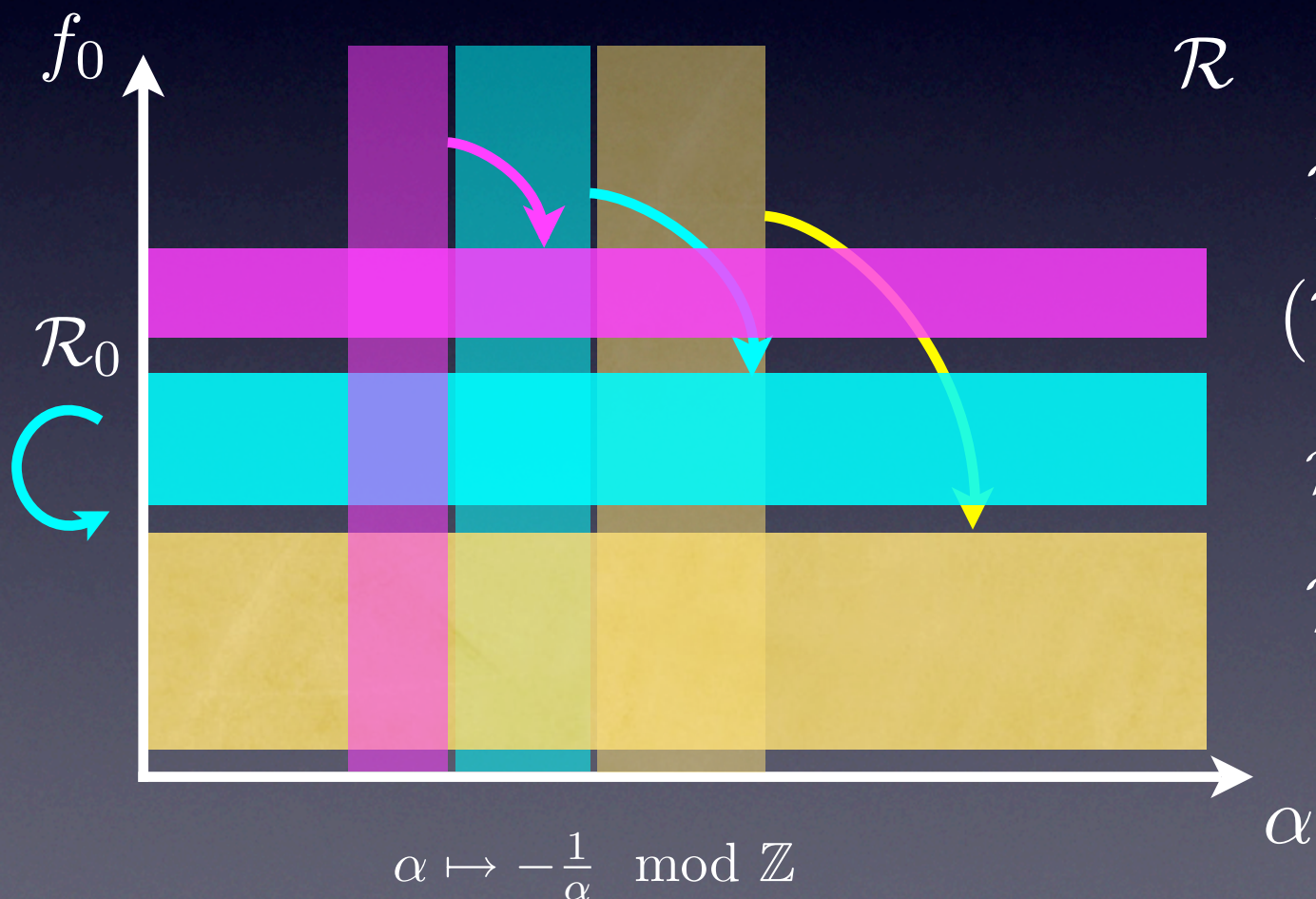
But, ... Can you really iterate infinitely many times?

Renormalization: The Picture

$f(z) = e^{2\pi i\alpha} z + O(z^2) = e^{2\pi i\alpha} f_0(z)$ where $f_0(z) = z + O(z^2)$ 1-parabolic

$f \leftrightarrow (\alpha, f_0)$

Write $\mathcal{R}f(z) = e^{-2\pi i\frac{1}{\alpha}} \mathcal{R}_\alpha f_0(z)$ then $\mathcal{R} : (\alpha, f_0) \mapsto (-\frac{1}{\alpha}, \mathcal{R}_\alpha f_0)$



\mathcal{R}

\mathcal{R} hyperbolic?

(\mathcal{R}_α contracting?)

$\mathcal{R}_\alpha f_0 \rightarrow \mathcal{R}_0 f_0$ ($\alpha \rightarrow 0$)

\mathcal{R}_0 contracting?

YES for α small

Main Theorems

Theorem 1 *Let $P(z) = z(1+z)^2$. There exist bounded simply connected open sets V and V' with $0 \in V \subset \bar{V} \subset V' \subset \mathbb{C}$ such that the class*

$$\mathcal{F}_1 = \left\{ f = P \circ \varphi^{-1} : \varphi(V) \rightarrow \mathbb{C} \mid \begin{array}{l} \varphi : V \rightarrow \mathbb{C} \text{ is univalent} \\ \varphi(0) = 0, \varphi'(0) = 1 \end{array} \right\}$$

satisfies the following: univalent = holomorphic and injective

(0) *every $f \in \mathcal{F}_1$ is non-degenerate;*

(i) *$\mathcal{F}_0 \setminus \{\text{quadratic polynomial}\}$ can be naturally embedded into \mathcal{F}_1 (in particular, $\mathcal{R}_0^n(z+z^2) \in \mathcal{F}_1$ $n = 1, 2, \dots$);*

(ii) *The renormalization \mathcal{R}_0 is well defined on \mathcal{F}_1 so that $\mathcal{R}_0(\mathcal{F}_1) \subset \mathcal{F}_1$;*

(iii) *If we write $\mathcal{R}_0 f = P \circ \psi^{-1}$, then ψ can be extended univalently to V' ;*

(iv) *$f \mapsto \mathcal{R}_0 f$ is “holomorphic.”*

Theorem 2 *The above statements hold for \mathcal{R}_α for α small. Hence there exists an N such that the above holds for*

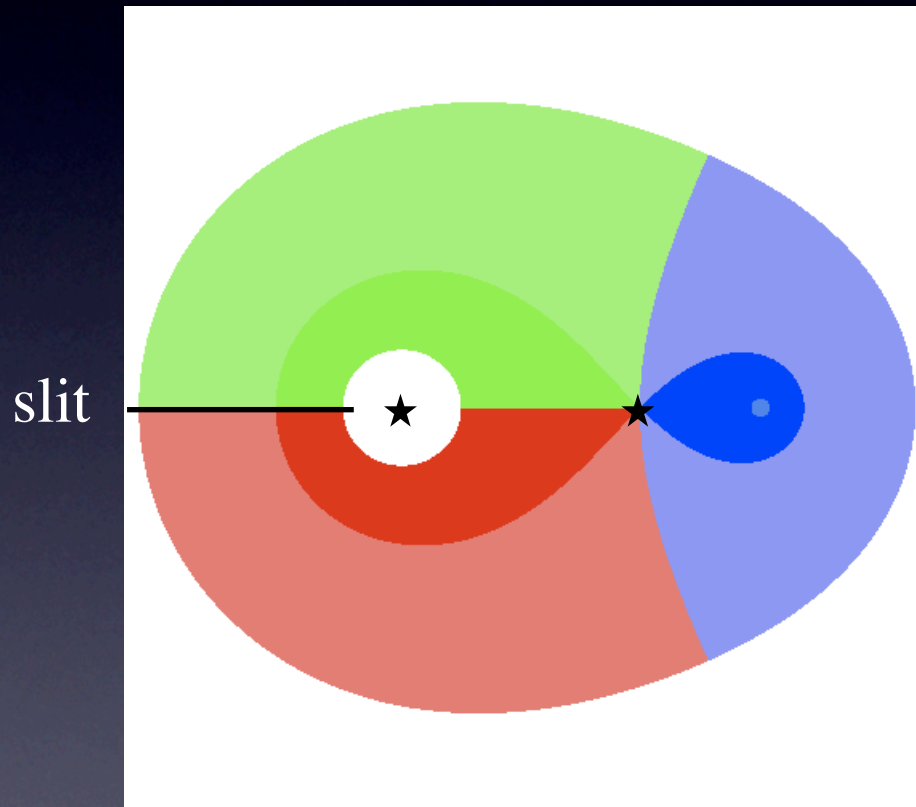
$$\alpha = \frac{1}{m + \beta} \quad \text{with } m \in \mathbb{N}, \beta \in \mathbb{C} \text{ and } |\beta| \leq 1.$$

$$P(z) = z(1+z)^2 \text{ and } V, V'$$

$$P(0) = 0, \quad P'(0) = 1$$

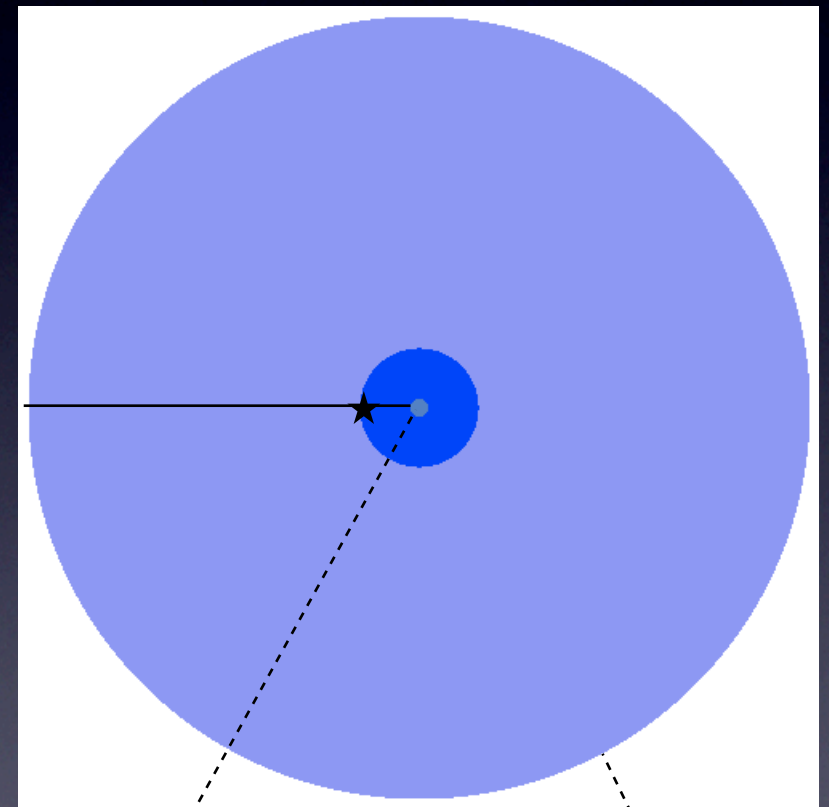
critical points: $-\frac{1}{3}$ and -1 critical values: $P(-\frac{1}{3}) = -\frac{4}{27}$ and $P(-1) = 0$

$$\eta = 2$$



V'

P
→



$$\frac{4}{27} e^{-2\pi\eta}$$

$$\frac{4}{27} e^{2\pi\eta}$$

V slightly smaller domain than V'

Applications

Theorem 2'. *There exists a large N such that the following holds:
If $f(z) = e^{2\pi i\alpha}h(z)$ with $h \in \mathcal{F}_1$ (or $h(z) = z + z^2$) and*

$$\alpha = \pm \frac{1}{a_1 \pm \frac{1}{a_2 \pm \frac{1}{\ddots}}} \quad \text{where } a_i \in \mathbb{N} \text{ and } a_i \geq N,$$

then the sequence of near-parabolic renormalizations

$$f = f_0 \xrightarrow{\mathcal{R}} f_1 \xrightarrow{\mathcal{R}} f_2 \xrightarrow{\mathcal{R}} f_3 \xrightarrow{\mathcal{R}} \dots$$

is defined so that $f_n = e^{2\pi i\alpha_n}h_n(z)$, $h_n \in \mathcal{F}_1$ (possibly with complex conjugation).

Corollary. *Under the assumption of Theorem 2', the critical orbit stays in the domain of f and can be iterated infinitely many times. Moreover if f is (a part of) a rational map, then the critical orbit is not dense.*

Julia set with positive area

Theorem (Buff-Chéritat). *There exists an irrational number α such that the Julia set of the quadratic polynomial $P_\alpha(z) = e^{2\pi i\alpha}z + z^2$ has positive Lebesgue measure.*

Adrien Douady's plan

$$P_{\alpha(0)} \xrightarrow{\text{perturb}} P_{\alpha(1)} \xrightarrow{\text{perturb}} P_{\alpha(2)} \xrightarrow{\text{perturb}} P_{\alpha(3)} \xrightarrow{\text{perturb}} \dots \longrightarrow P_\alpha$$

Each $\alpha(n)$ is of bounded type, hence $P_{\alpha(n)}$ has a Siegel disk.

$$\alpha(0) = [0, N, N, N, N, \dots]$$

$\alpha(n) \rightsquigarrow \alpha(n+1)$: modify only a_k (larger), for some large k .

Want:

Do not lose too much area in each step $K(P_{\alpha(n)}) \rightsquigarrow K(P_{\alpha(n+1)})$.

$$\alpha = \lim_{n \rightarrow \infty} \alpha(n) \text{ non Bruno, } |J(P_\alpha)| = |K(P_\alpha)| \geq \limsup_{n \rightarrow \infty} |K(P_{\alpha(n)})| > 0$$

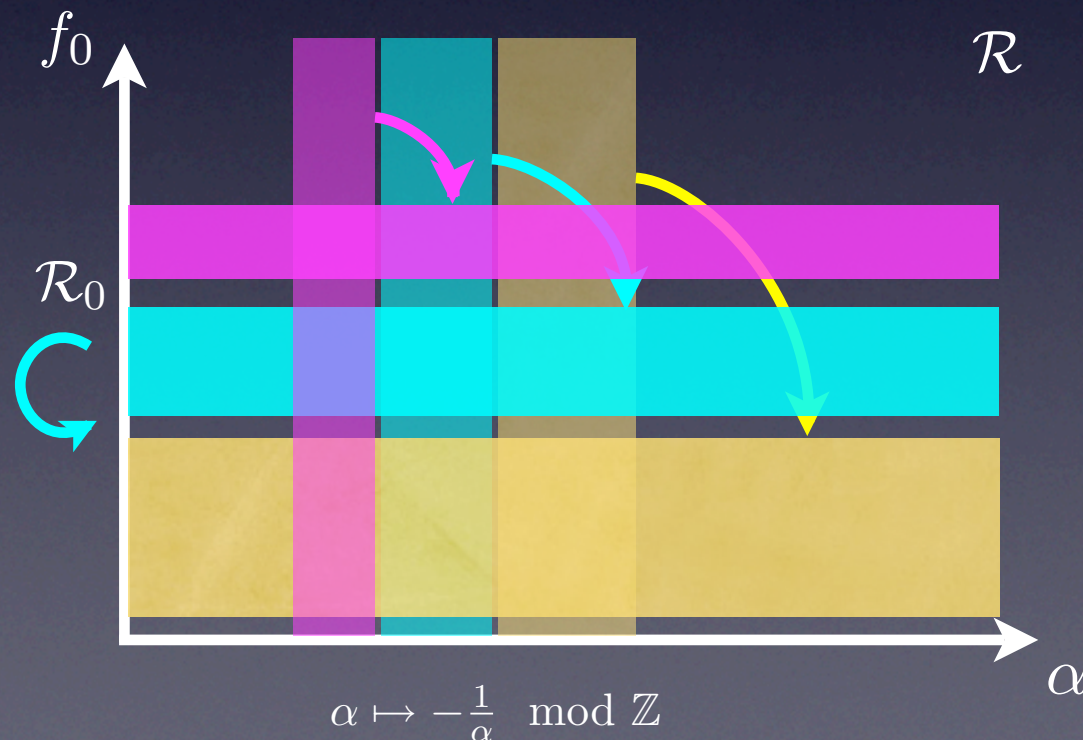
Need to control the boundary of Siegel disks, which is the closure of the critical orbit. “Semi-continuity” by our Theorem 2’ which yields a uniform control on the sequence of renormalizations, even when a_k are taken to be very big.

Also uses McMullen’s renormalization result on Siegel disks of bounded type rotation number.

Contraction and Hyperbolicity

Theorem 3 *Modifying the definition slightly (requiring that φ has a quasi-conformal extension to \mathbb{C}), \mathcal{F}_1 is in one to one correspondence with the Teichmüller space $\text{Teich}(W)$ of $W = \mathbb{C} \setminus \bar{V} (\simeq \mathbb{D}^*)$. The induced map $\mathcal{R}_0^{\text{Teich}}$ is a uniform contraction with respect to the Teichmüller distance. (The Lipschitz constant $\leq \exp(-2\pi \text{mod}(V' \setminus \bar{V}))$.)*

Theorem 4 *The above statements hold for the fiber map \mathcal{R}_α for α small. Hence the total renormalization \mathcal{R} is hyperbolic in this region.*



$$f = e^{2\pi i \alpha} f_0$$

$$\mathcal{R} : (\alpha, f_0) \mapsto \left(-\frac{1}{\alpha}, \mathcal{R}_\alpha f_0\right)$$

$$\mathcal{R}_\alpha f_0 \rightarrow \mathcal{R}_0 f_0 \quad (\alpha \rightarrow 0)$$

Proof of Theorems 3 and 4

$\mathcal{F}_1 \ni f = P \circ \varphi^{-1} \rightsquigarrow [\tilde{\varphi}|_W] \in \text{Teich}(W) = \{\psi : W \rightarrow \mathbb{C} \text{ qc}\} / \sim$
where $\tilde{\varphi}$ is a quasiconformal extension of φ to \mathbb{C} .

Teichmüller space is like the unit disk with Poincaré metric.
holomorphic self map does not expand the distance,
and is often a contraction.

Royden-Gardiner Theorem (Teichmüller distance = Kobayashi distance)
cotangent space = {integrable holomorphic quadratic differentials}
modulus-area inequality for holom. quad. differentials
isoperimetric inequality for holom. quad. differentials
modified Carleman's inequality

Application of contraction

Theorem. *Suppose f and f' satisfy the assumption of Theorem 2', with the same rotation number α . Then they have small periodic cycles ζ_n and ζ'_n around 0 with period q_n . Let $\lambda(\zeta_n)$, $\lambda(\zeta'_n)$ be their multipliers. The differences*

$$|\lambda(\zeta_n) - \lambda(\zeta'_n)| \quad \text{and} \quad \left| \frac{1}{1 - \lambda(\zeta_n)} - \frac{1}{1 - \lambda(\zeta'_n)} \right|$$

tends to 0 exponentially fast as $n \rightarrow \infty$ with a uniform rate.

Corollary. *Not arbitrary germ with irrationally indifferent fixed point can be realized in such a map (for example, quadratic polynomials, or a perturbation of rational map with parabolic fixed point with only one critical point in its basin).*

cf. Yoccoz unrenormalization construction for the inverse of Siegel-Bruno Theorem.

Grazie!

A draft is available at

<http://www.math.kyoto-u.ac.jp/~mitsu/pararenorm/>