

Some Extensions and Analysis of Flux and Stress Theory

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Structures of the Mechanics of Complex Bodies

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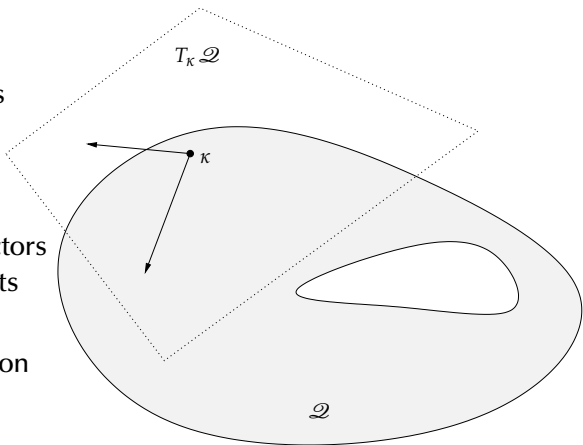
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The Global Point of View

C^n -Functionals

Review of Basic Kinematics and Statics on Manifolds

- The mechanical system is characterized by its configuration space—a manifold \mathcal{Q} .
- *Velocities* are tangent vectors to the manifold—elements of $T\mathcal{Q}$.
- A *Force* at the configuration κ is a linear mapping $F: T_{\kappa}\mathcal{Q} \rightarrow \mathbb{R}$.



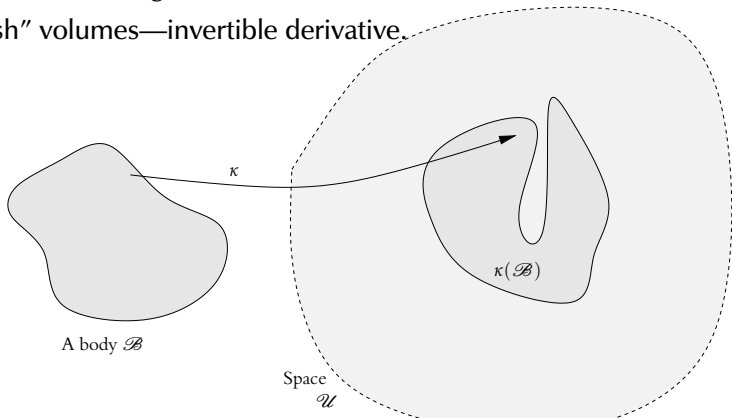
Can we apply this framework to Continuum Mechanics?

Problems Associated with the Configuration Space in Continuum Mechanics

- What is a configuration?
- Does the configuration space have a structure of a manifold?
- The configuration space for continuum mechanics is infinite dimensional.

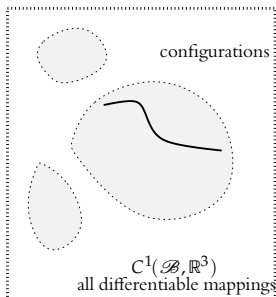
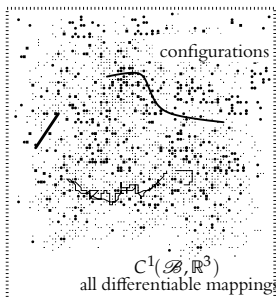
Configurations of Bodies in Space

- A mapping of the body into space;
- material impenetrability—one-to-one;
- continuous deformation gradient (derivative);
- do not “crash” volumes—invertible derivative.



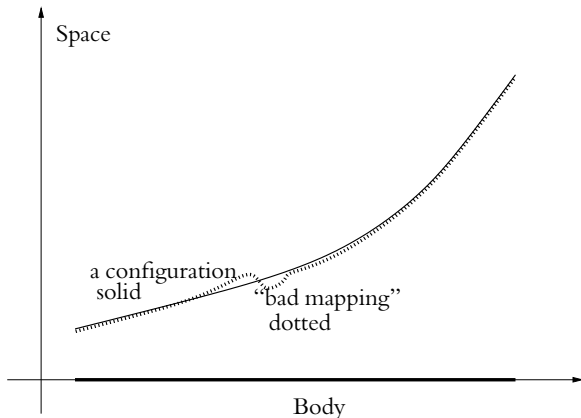
Manifold Structure for Euclidean Geometry

- If the body is a subset of \mathbb{R}^3 and space is modeled by \mathbb{R}^3 , the collection of differentiable mappings $C^1(\mathcal{B}, \mathbb{R}^3)$ is a vector space
- However, the subset of “good” configurations is not a vector space, e.g., $\kappa - \kappa = 0$ —not one-to-one.
- We want to make sure that the subset of configurations \mathcal{Q} is an open subset of $C^1(\mathcal{B}, \mathbb{R}^3)$, so it is a trivial manifold.



The C^0 -Distance Between Functions

- The C^0 -distance between functions measures the maximum difference between functions.
- A configuration is arbitrarily close to a “bad” mapping.

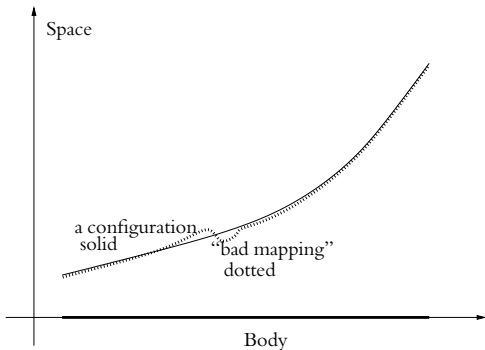


The C^1 -Distance Between Functions

- The C^1 distance between functions measures the maximum difference between functions and their derivative

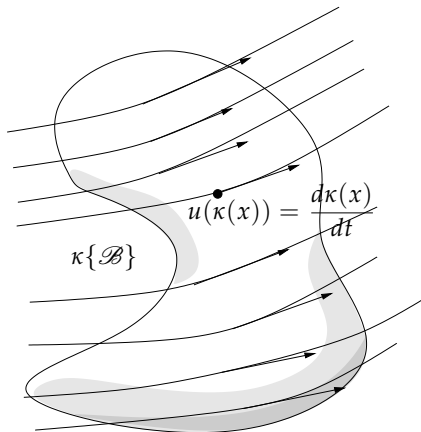
$$\|u - v\|_{C^1} = \sup\{|u(x) - v(x)|, |Du(x) - Dv(x)|\}.$$

- A configuration is always a finite distance away from a “bad” mapping.



Conclusions for \mathbb{R}^3

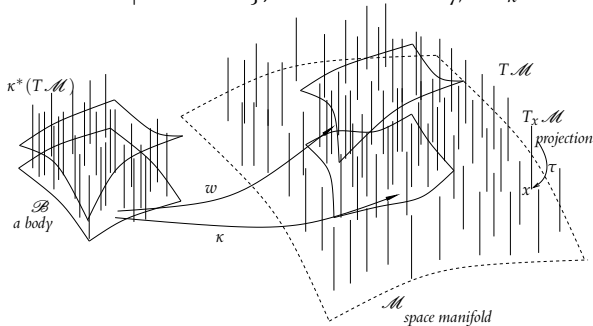
- If we use the C^1 -norm, the configuration space of a continuous body in space is an open subset of $C^1(\mathcal{B}, \mathbb{R}^3)$ -the vector space of all differentiable mapping.
- \mathcal{Q} is a trivial infinite dimensional manifold and its tangent space at any point may be identified with $C^1(\mathcal{B}, \mathbb{R}^3)$.
- A tangent vector is a velocity field.



For Manifolds

- Both the body \mathcal{B} and space \mathcal{U} are differentiable manifolds.
- The configuration space is the collection $\mathcal{Q} = \text{Emb}(\mathcal{B}, \mathcal{U})$ of the embeddings of the body in space. This is an open submanifold of the infinite dimensional manifold $C^1(\mathcal{B}, \mathcal{U})$.
- The tangent space $T_\kappa \mathcal{Q}$ may be characterized as

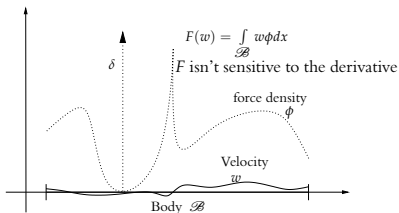
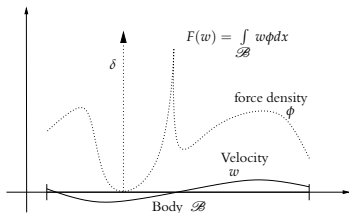
$$T_\kappa \mathcal{Q} = \{w: \mathcal{B} \rightarrow T\mathcal{Q} \mid \tau \circ w = \kappa\}, \quad \text{or alternatively,} \quad T_\kappa \mathcal{Q} = C^1(\kappa^* T\mathcal{U}).$$



Representation of C^0 -Functionals by Integrals

- Assume you measure the size of a function using the C^0 -distance, $\|w\| = \sup\{|w(x)|\}$.
- A linear functional $F: w \mapsto F(w)$ is continuous with respect to this norm if $F(w) \rightarrow 0$ when $\max |w(x)| \rightarrow 0$.
- *Riesz representation theorem*: A continuous linear functional F with respect to the C^0 -norm may be represented by a unique measure μ in the form

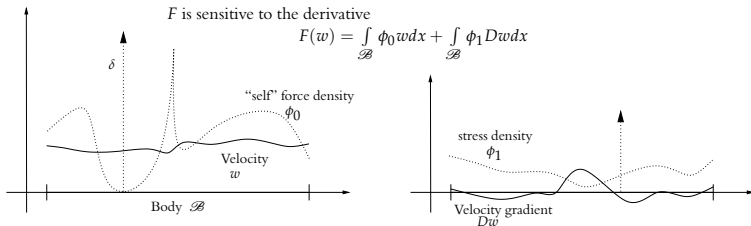
$$F(w) = \int_{\mathcal{B}} w d\mu.$$



Representation of C^1 -Functionals by Integrals

- Now, you measure the size of a function using the C^1 -distance, $\|w\| = \sup\{|w(x)|, |Dw(x)|\}$.
- A linear functional $F: w \mapsto F(w)$ is continuous with respect to this norm if $F(w) \rightarrow 0$ when both $\max |w(x)| \rightarrow 0$ and $\max |Dw(x)| \rightarrow 0$.
- Representation theorem: A continuous linear functional F with respect to the C^1 -norm may be represented by measures σ_0, σ_1 in the form

$$F(w) = \int_{\mathcal{B}} w d\sigma_0 + \int_{\mathcal{B}} Dw d\sigma_1.$$

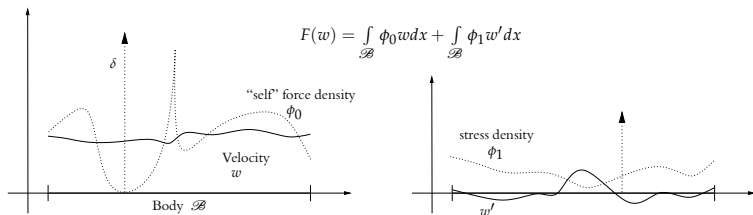


Non-Uniqueness of C^1 -Representation by Integrals

- We had an expression in the form

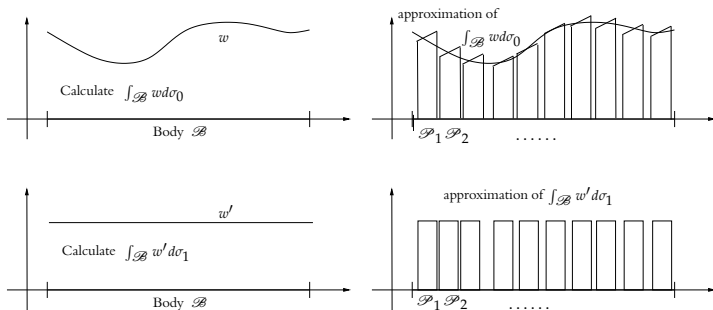
$$F(w) = \int_{\mathcal{B}} w d\sigma_0 + \int_{\mathcal{B}} w' d\sigma_1.$$

- If we were allowed to vary w and w' independently, we could determine σ_0 and σ_1 uniquely.
- This cannot be done because of the condition $w' = Dw$.



Unique Representation of a Force System

- Assume we have a force system, i.e., a force $F_{\mathcal{P}}$ for every subbody \mathcal{P} of \mathcal{B} .
- We can approximate pairs of non-compatible functions w and w' , i.e., $w' \neq Dw$, by piecewise compatible functions.



- This way the two measures are determined uniquely.
- One needs consistency conditions for the force system.

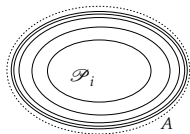
Generalized Cauchy Consistency Conditions

- *Additivity:*

$$F_{\mathcal{P}_1 \cup \mathcal{P}_2}(w|_{\mathcal{P}_1 \cup \mathcal{P}_2}) = F_{\mathcal{P}_1}(w|_{\mathcal{P}_1}) + F_{\mathcal{P}_2}(w|_{\mathcal{P}_2}).$$



- *Continuity:* If $\mathcal{P}_i \rightarrow A$, then $F_{\mathcal{P}_i}(w|_{\mathcal{P}_i})$ converges and the limit depends on A only.



- *Uniform Boundedness:* There is a $K > 0$ such that for every subbody \mathcal{P} and every w ,

$$|F_{\mathcal{P}}(w|_{\mathcal{P}})| \leq K \|w_{\mathcal{P}}\|.$$

Main Tool in Proof: Approximation of measurable sets by bodies with smooth boundaries.

Generalizations

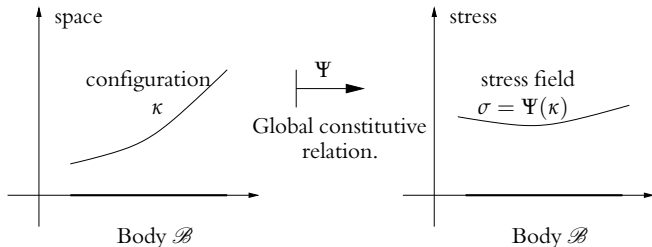
- All the above may be formulated and proved for differentiable manifolds.
- This formulation applies to continuum mechanics of order $k > 1$ (stress tensors of order k). One should simply use the C^k -norm instead of the C^1 -norm.
- The generalized Cauchy conditions also apply to continuum mechanics of order $k > 1$. This is the only formulation of Cauchy conditions for higher order continuum mechanics.

Locality and Continuity in Constitutive Theory

Global Constitutive Relations

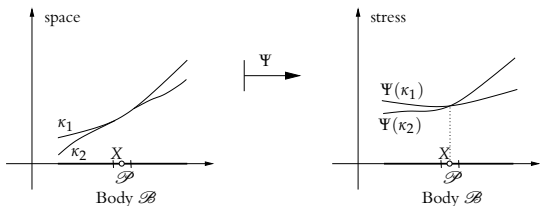
(Elasticity for Simplicity)

- \mathcal{Q} , the configuration space of a body \mathcal{B} .
- $C^0(\mathcal{B}, L(\mathbb{R}^3, \mathbb{R}^3))$, the collection of all stress fields over the body.
- $\Psi: \mathcal{Q} \rightarrow C^0(\mathcal{B}, L(\mathbb{R}^3, \mathbb{R}^3))$, a global constitutive relation.

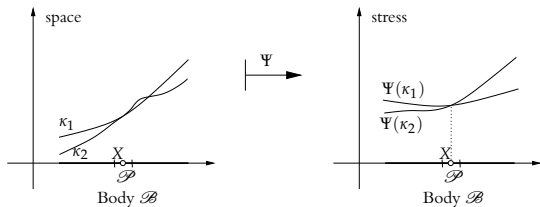


Locality and Materials of Grade- n

Germ Locality: If two configurations κ_1 and κ_2 are equal on a subbody containing X , then the resulting stress fields are equal at X .



Material of Grade- n or n -Jet Locality: If the first n derivatives of κ_1 and κ_2 are equal at X , then, $\Psi(\kappa_1)(X) = \Psi(\kappa_2)(X)$.
(Elastic = grade 1.)



n -Jet Locality and Continuity

Basic Theorem: If a constitutive relation $\Psi: \mathcal{Q} \rightarrow C^0(\mathcal{B}, L(\mathbb{R}^3, \mathbb{R}^3))$ is local and continuous with respect to the C^n -norm, then, it is n -jet local. In particular, if Ψ is continuous with respect to the C^1 -topology, the material is elastic.

